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A STUDY OF LONGITUDINAL OSCILLATIONS
OF PROPELLANT TANKS AND WAVE PROPAGATIONS IN FEED LINES
PART II - Longitudinal Oscillation of a Liquid-Filled
Elastic Cylindrical Tank with a Flexible
Inverted Conical Bulkhead



3 April 1967

NAS8-11490

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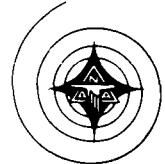
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FOREWORD

This report was prepared by the Space and Information Systems Division of North American Aviation, Inc., Downey, California, for the George C. Marshall Space Flight Center, National Aeronautics and Space Administration, Huntsville, Alabama, under the Supplemental Agreement Modification No. 2 of Contract NAS8-11490, "Study of Longitudinal Oscillations of Propellant Tanks and Wave Propagations in Feed Lines," dated April 6, 1966. Dr. George F. McDonough (Principal) and Mr. Robert S. Ryan (Alternate) of Aero-Astroynamics Laboratory, MSFC, are Contracting Officer Representatives. The work is published in two separate parts:

PART I - Propagating Pressure Waves in a Fluid Filled Cylindrical Shell

PART II - Longitudinal Oscillation of A Liquid-Filled Elastic Cylindrical Tank with A Flexible Inverted Conical Bulkhead

The project was carried out by the Vehicle Dynamics Branch, Structures and Materials Department of Research, Engineering, and Testing Division, S&ID. Dr. F. C. Hung was the Program Manager for North American Aviation, Inc. The study was conducted by Dr. Clement L. Tai (Principal Investigator), Dr. Shoichi Uchiyama, and Mr. John S. Kanno. The computer program was developed by Mr. Shigeo Miyashiro.



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ABSTRACT

A procedure has been formulated to determine the natural frequencies and the corresponding mode shapes of a liquid-filled elastic cylindrical tank with a flexible inverted conical bulkhead. It is assumed that the liquid is incompressible and inviscid and its motion is irrotational. Under such assumptions, the velocity potential is obtained from the solution of Laplace's equation in both circular cylindrical and spherical polar coordinates. This velocity potential, together with Bernoulli's equation, permits the evaluation of the fluctuating liquid pressure at the liquid-shell interfaces.

The interface pressure is treated as a forcing function in the shell equations, and the shell displacement components are then determined analytically. An eigenvalue problem was constructed by the least square technique through the boundary conditions and interface conditions. The solution of such an eigenvalue problem are the desired natural frequencies and the corresponding mode shapes of the system.

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NOMENCLATURE

A_n , B_n	Arbitrary constants
a	Radius of cylindrical shell
b	Height of cone
C_n , $C_{\hat{n}}$, $C_{\bar{n}}$	Undetermined constants $n = 0, 1, \dots, N$; $\hat{n} = 2n$; $\bar{n} = 2n + 1$
e_s , e_ϕ	Extensional and angular (hoop) strain.
D	$Eh/(1 - \nu^2)$, shell extension rigidity.
E	Young's modulus
F , f	Defined functions
h_1 , h_2	Thickness of cylindrical shell and cone respectively.
$J_n(\mu, \bar{s})$	Bessel's function of nth order
N_s , N_ϕ , N_z	Stress resultant in the direction of s , ϕ , z respectively
$P_n(\mu)$	Legendre polynomials
p	Pressure
r	Cylindrical radial coordinate
s	Spherical radial coordinate
t	Time
u_1 , v_1 , w_1	Axial, meridional and radial displacement component, respectively of cylindrical shell.
u_2 , w_2	Displacement component of cone in the direction of s and θ
z	Axial coordinate



\bar{r} , \bar{s} , \bar{u} , \bar{v} , \bar{w} , \bar{z} , etc. Nondimensional quantity of r , s , u , v , w , z , etc.

g_{mn} , h_{mn} , ℓ_{mn}	Defined constant
α	Semivortex angle of cone
α_n ,	Positive root of Bessel's function of first order
β_n	Undetermined constant
$\Gamma(x)$	Gamma function
ϵ_1 , ϵ_2 , ...	Functional errors
θ	Meridional spherical coordinate
κ_1 , κ_2	$\rho_1 h_1 a^2 / D_1$ and $\rho_2 h_2 a^2 / D_2$ respectively
λ	Defined constant
μ	$\cos\theta$
ν	Poisson's ratio
ρ_f , ρ_1 , ρ_2	Density of fluid, cylindrical shell, and cone respectively
σ	Defined constant
\bar{t}	Nondimensional quantity of t
Φ_1 , Φ_2	Velocity potential
ϕ_1 , ϕ_2	Velocity potential for steady flow
ω	Natural frequency



I. INTRODUCTION

In the development of launch vehicle and spacecraft design, a knowledge of the dynamic behavior of thin-walled fuel tanks is of great importance. Of prime importance are the natural frequencies and the corresponding mode shapes of vibrations. In a recent investigation, the author (Reference 1) studied the coupled oscillations of a liquid with a free surface in a flexible oblate spheroidal tank. For a better understanding of such coupling of oscillation, the study of the interaction of dynamic behavior of a flexible liquid-filled cylindrical tank with a flexible inverted conical bulkhead is developed.

A number of studies have been made during the past years of the dynamic behavior of cylindrical tanks with various shapes of bulkhead. Bhuta and Koval (Reference 2) studied the problem of sloshing of a liquid in a cylindrical tank with rigid walls and a flexible bottom. Bleich (Reference 3) investigated longitudinal forced vibrations of cylindrical fuel tanks with rigid walls and elastic bottoms of arbitrary shape. Coale and Nagano (Reference 4) dealt with the axisymmetric dynamic behavior of a cylindrical shell with a hemispherical shell bottom. The present study is concerned with the axisymmetric dynamic behavior of an elastic, liquid-filled cylindrical tank with a flexible inverted conical bulkhead.

The present method of analysis is based on two previous papers, one for a hemispherical tank (Reference 5) and one for an oblate spheroidal tank (Reference 1), both of which are fully filled by the liquid. The mathematical model for the analysis of such a system is considered and its geometry is described by a circular cylindrical coordinate system for the cylindrical shell and a spherical polar coordinate system for the inverted conical shell as shown in Figure 1. The motion of the liquid is represented by a velocity potential which satisfies the continuity equation stated by Laplace's equation. The motion of the shell is represented by two displacement components, one in the normal direction and one in the tangential direction. The liquid velocity potential in the cylindrical shell is represented by a series of the product of Bessel functions and trigonometric functions. The liquid velocity potential in the inverted conical shell is represented by a series of Legendre functions. The shell displacement components for both the cylindrical shell and the conical shell are expressed in terms of the liquid velocity potential. Coefficients in these series are selected to satisfy the liquid-shell interface conditions by the least squared error technique. For the numerical calculation of the natural frequencies and the corresponding mode shapes, the procedure is described in detail at the end of this report.



II. EQUATIONS OF MOTION AND BOUNDARY CONDITIONS

The system to be analyzed is a thin, flexible, liquid-filled circular cylindrical shell with a thin, flexible, inverted conical shell. The configuration is shown in Figure 1. For analysis, the system is divided into four elements: (1) a cylinder of liquid with a free surface, (2) an inverted cone of liquid, (3) a circular cylindrical shell, and (4) a circular conical shell. The liquid in elements (1) and (2) is assumed to be incompressible, inviscid, and irrotational. The shell in elements (3) and (4) is treated as a membrane and hence the bending effects are ignored. For a thin shell whose thickness-to-radius ratio is very small, such as those in the present case, this assumption is justified. The inertial forces of elements (3) and (4) are included in the equations of motion. Only a small-amplitude, longitudinal, axisymmetric motion is considered. As shown in Figure 1, the system is restrained against axial motion at the lower end of the cylindrical shell but is unrestrained otherwise. The upper end of the cylindrical shell is unrestrained longitudinally.

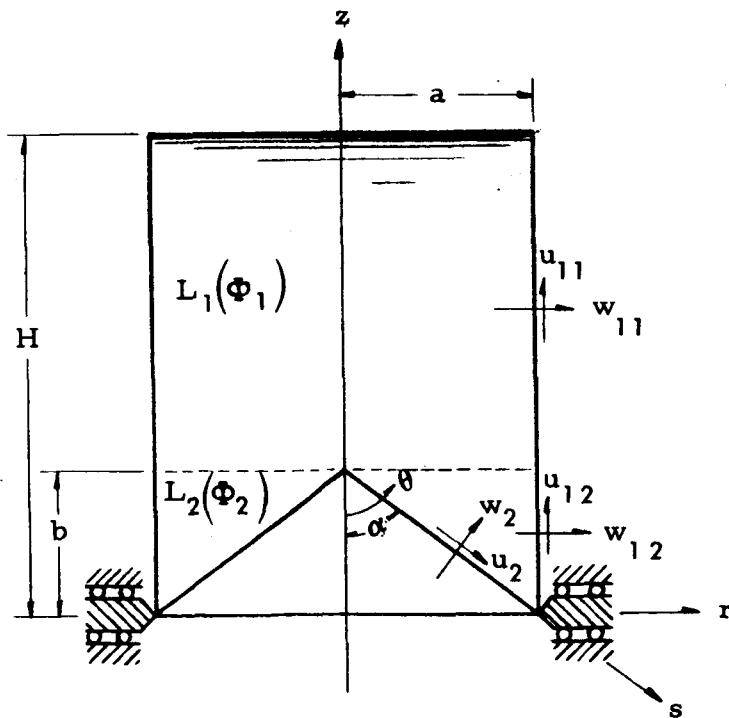


Figure 1. Liquid-Filled Circular Cylindrical Tank with An Inverted Conical Bulkhead at the Bottom

1. Liquid L_1

Motion of the liquid L_1 can be expressed in terms of a velocity potential $\Phi_1(r, z, t)$ which satisfies the axisymmetric Laplace equation in circular cylindrical coordinates

$$\nabla^2 \Phi_1 = \frac{\partial^2 \Phi_1}{\partial r^2} + \frac{1}{r} \frac{\partial \Phi_1}{\partial r} + \frac{\partial^2 \Phi_1}{\partial z^2} = 0 \quad (1)$$

where r and z are circular cylindrical coordinates and shown in Figure 1.

Two of the boundary conditions associated with this element are

$$\left. \frac{\partial \Phi_1}{\partial r} \right|_{r=0} = 0 \quad (2)$$

which is required to avoid a singular solution at the axis, and

$$\left[\frac{\partial^2 \Phi_1}{\partial t^2} + g \frac{\partial \Phi_1}{\partial z} \right]_{z=H} = 0 \quad (3)$$

which is the linearized boundary condition that is to be satisfied at the free surface.

The remaining boundary conditions are interface conditions coupling this element to other elements of the system. At the liquid-shell interface, the boundary condition is

$$\left. \frac{\partial \Phi_1}{\partial r} \right|_{r=a} = \frac{\partial w_{11}}{\partial t} \quad (4)$$



where $\frac{\partial \Phi_1}{\partial r}$ is taken positively in the direction of increasing r and w_{11} is the radial displacement of the cylindrical shell to be determined in terms of Φ_1 . Between the two elements of liquid, pressure, and normal velocity of the two elements must match, namely,

$$\left. \frac{\partial \Phi_1}{\partial t} \right|_{z=b} = \left. \frac{\partial \Phi_2}{\partial t} \right|_{\theta=\frac{\pi}{2}} \quad (5)$$

and

$$\left. \frac{\partial \Phi_1}{\partial z} \right|_{z=b} = \frac{1}{s} \left. \frac{\partial \Phi_2}{\partial \theta} \right|_{\theta=\frac{\pi}{2}} \quad (6)$$

where Φ_2 is a velocity potential of the liquid L_2 , and s and θ are spherical polar coordinates.

2. Liquid L_2

Motion of the liquid L_2 can be expressed in terms of a velocity potential $\Phi_2(s, \theta, t)$ which satisfies the Laplace equation in spherical polar coordinates

$$\nabla^2 \Phi_2 = \frac{\partial^2 \Phi_2}{\partial s^2} + \frac{2}{s} \frac{\partial \Phi_2}{\partial s} + \frac{1}{s^2} \frac{\partial^2 \Phi_2}{\partial \theta^2} + \frac{\cot \theta}{s^2} \frac{\partial \Phi_2}{\partial \theta} = 0 \quad (7)$$

The boundary conditions for Φ_2 are interface conditions and two of them are given by Eqs. (5) and (6). The other interface conditions are



$$\left[\frac{r}{\sqrt{r^2 + (b-z)^2}} \frac{\partial \Phi_2}{\partial s} + \frac{b-z}{r^2 + (b-z)^2} \frac{\partial \Phi_2}{\partial \theta} \right] = \frac{\partial w_{12}}{\partial t}$$

$$\begin{cases} r = a \\ s = \sqrt{a^2 + (b-z)^2} \\ \theta = \tan^{-1} \frac{a}{b-z} \end{cases}$$

(8)

and

$$\frac{1}{s} \left[\frac{\partial \Phi_2}{\partial \theta} \right]_{\theta=\alpha} = \frac{\partial w_2}{\partial t}$$

(9)

where $\frac{\partial \Phi_2}{\partial s}$ and $\frac{1}{s} \frac{\partial \Phi_2}{\partial \theta}$ are taken positively in the direction of increasing s and θ respectively and the left-hand side of Eq. (8) represents the component of Φ_2 in the direction normal to the cylindrical shell surface. It is obtained from

$$\frac{\partial \Phi_2}{\partial n} = \vec{n} \cdot \nabla \Phi_2 = \frac{\partial \Phi_2}{\partial s} \sin \theta + \frac{1}{s} \frac{\partial \Phi_2}{\partial \theta} \cos \theta$$

(10)

where \vec{n} is the unit outward normal vector.

3. Circular Cylindrical Shell Filled with Liquid L_1

Let the coordinate axes be chosen so that z is in the direction of the generatrix and r is in the direction of the normal to middle surface of the shell as shown in Figure 1. The axisymmetric differential equations (Reference 6) for this shell can then be expressed in terms of displacements $u_{11}(z,t)$ and $w_{11}(z,t)$ as



$$a \frac{\partial^2 u_{11}}{\partial z^2} + v_1 \frac{\partial w_{11}}{\partial z} = \frac{\rho_1 h_1 a}{D_1} \frac{\partial^2 u_{11}}{\partial t^2} \quad (11)$$

$$v_1 a \frac{\partial u_{11}}{\partial z} + w_{11} = - \frac{\rho_1 h_1 a^2}{D_1} \frac{\partial^2 w_{11}}{\partial t^2} + \frac{a^2 p_1}{D_1} \quad (12)$$

where v_1 is the Poisson's ratio, ρ_1 is the mass density, h_1 is the thickness, and D_1 is the extensional modulus ($= E_1 h_1 / (1-v_1^2)$) of the circular cylindrical shell, and p_1 is the liquid pressure acting at the shell wall.

From the linearized Bernoulli equation, the liquid pressure p_1 may be expressed as

$$p_1 = - \rho_f \frac{\partial \Phi_1}{\partial t} \quad (13)$$

where ρ_f is the mass density of the liquid.

The boundary conditions are

$$\left. \begin{array}{l} u_{11} = u_{12} \\ \frac{\partial u_{11}}{\partial z} = \frac{\partial u_{12}}{\partial z} \end{array} \right\} \text{at } z = b \quad (14)$$

which indicates that the shell is continuous at $z = b$ and

$$\left. N_z \right|_{z=H} = D_1 \left(\frac{\partial u_{11}}{\partial z} + v \frac{w_{11}}{a} \right) \Big|_{z=H} = 0 \quad (15)$$

which specifies that there is no axial force acting on the upper end of the shell. There is no boundary condition which can be applied to w_{11} because with bending stiffness neglected the differential equation is of zero order in w_{11} .



4. Circular Conical Shell

Let s be chosen in the direction of the generatrix and θ be measured in the direction of the normal to the middle surface of the shell as shown in Figure 1. In this case, the curvilinear coordinates are s and θ . Let $R_1 = R_s$ and $R_2 = R_\theta$ be the principal radii of curvature. Then $R_s = \infty$ and $R_\theta = s \tan\alpha$ where α is a semi-vertex angle of the cone.

From the general expression of the equations of motion of a thin shell in terms of curvilinear coordinates (Reference 5), the equations of motion of the circular conical membrane under axially symmetric conditions are obtained as

$$\frac{\partial s N_s}{\partial s} - N_\phi + s f_s = 0 \quad (16)$$

$$\frac{N_\phi}{\tan\alpha} + s(p_2 - f_\theta) = 0 \quad (17)$$

where N_s and N_ϕ are the stress resultants, f_s and f_θ are the inertia forces, and p_2 is the liquid dynamic pressure acting at both the circular cylindrical shell wall and the circular conical shell wall.

The stress-strain relations give

$$N_s = D_2(e_s + v_2 e_\phi) \quad (18)$$

$$N_\phi = D_2(e_\phi + v_2 e_s) \quad (19)$$

where e_s and e_ϕ are the extensional and angular strain, respectively, and D_2 is the extensional modulus ($= E_2 h_2 / [1 - v_2^2]$) of the circular conical shell.

The strain-displacement relations are obtained from the general expression (Reference 7):

$$e_s = \frac{\partial u_2}{\partial s} \quad (20)$$

$$e_\phi = \frac{1}{2}(u_2 + w_2 \cot\alpha) \quad (21)$$



Substituting Eqs. (20) and (21) into Eqs. (18) and (19) gives

$$N_s = D_2 \left[\frac{\partial u_2}{\partial s} + \frac{v}{s} \left(u_2 + w_2 \cot \alpha \right) \right] \quad (22)$$

$$N_\phi = D_2 \left[\frac{1}{s} \left(u_2 + w_2 \cot \alpha \right) + v \frac{\partial u_2}{\partial s} \right] \quad (23)$$

The inertia forces f_s and f_θ are given by

$$f_s = -\rho_2 h_2 \frac{\partial^2 u_2}{\partial t^2} \quad (24)$$

$$f_\theta = -\rho_2 h_2 \frac{\partial^2 w_2}{\partial t^2} \quad (25)$$

where ρ_2 is the mass density and h_2 is the thickness of the circular conical shell.

From the linearized Bernoulli equation, the liquid pressure p_2 may be expressed as

$$p_2 = -\rho_f \frac{\partial \Phi_2}{\partial t} \quad (26)$$

Substituting Eqs. (22) through (26) into Eqs. (16) and (17) gives

$$s \frac{\partial^2 u_2}{\partial s^2} + \frac{\partial u_2}{\partial s} - \frac{u_2}{s} + v \frac{\partial w_2}{\partial s} \cot \alpha - \frac{w_2}{s} \cot \alpha = \frac{\rho_2 h_2 s}{D_2} \frac{\partial^2 u_2}{\partial t^2} \quad (27)$$

$$v \frac{\partial u_2}{\partial s} + \frac{u_2}{s} + \frac{w_2}{s} \cot \alpha = -\frac{\rho_2 h_2 s}{D_2} \frac{\partial^2 w_2}{\partial t^2} \tan \alpha - \frac{s p_{2\alpha}}{D_2} \tan \alpha \quad (28)$$

where $p_{2\alpha}$ is the pressure p_2 at $\theta = \alpha$.



The boundary conditions are

$$u_2 = \text{finite} \quad \text{at } s = 0 \quad (29)$$

which is required to avoid a singular solution at the apex and

$$u_2 - w_2 \tan\alpha = 0 \quad \text{at } s = \sqrt{a^2 + b^2} \quad (30)$$

which assumes that there is no axial displacement at the bottom of the conical shell.

5. Circular Cylindrical Shell Filled with Liquid L_2

The axisymmetric differential equations for this shell can be expressed as

$$a \frac{\partial^2 u_{12}}{\partial z^2} + v_1 \frac{\partial w_{12}}{\partial z} = \frac{\rho_1 h_1 a}{D_1} \frac{\partial^2 u_{12}}{\partial t^2} \quad (31)$$

$$v_1 a \frac{\partial u_{12}}{\partial z} + w_{12} = - \frac{\rho_1 h_1 a^2}{D_1} \frac{\partial^2 w_{12}}{\partial t^2} + \frac{a^2 p_2}{D_1} \quad (32)$$

The boundary conditions are Eqs. (14) and

$$u_{12} \Big|_{z=0} = 0 \quad (33)$$

which indicates that there is no axial motion of the base of the shell. There is again no boundary condition which can be applied to w_{12} .



III. SOLUTIONS OF VELOCITY POTENTIALS AND SHELL DISPLACEMENTS

1. Liquid L_1

The equations for determining the shell displacements are found by assuming a simple harmonic motion of the system. For such a simple harmonic motion, the velocity potential, Φ_1 , may be taken in the form

$$\Phi_1 = a^2 \omega \bar{\phi}_1(r, z) \cos \omega t \quad (34)$$

where $\bar{\phi}_1(r, z)$ is the dimensionless velocity potential for steady flow, ω is the natural frequency of the system.

By separation of variables, $\Phi(r, z)$ is obtained in the form

$$\begin{aligned} \bar{\phi}_1(r, z) &= e_{10} + e_{20} \ln\left(\frac{r}{a}\right) + \left[e_{30} + e_{40} \ln\left(\frac{r}{a}\right) \right] \frac{z}{a} \\ &+ \sum_{n=1}^{\infty} \frac{a}{g} A_n \left[a_n \sinh\left(\frac{\alpha_n z}{a}\right) + \cosh\left(\frac{\alpha_n z}{a}\right) \right] J_0\left(\frac{\alpha_n r}{a}\right) \\ &+ \sum_{n=1}^{\infty} B_n \left[b_n \sin\left(\frac{\beta_n z}{a}\right) + \cos\left(\frac{\beta_n z}{a}\right) \right] I_0\left(\frac{\beta_n r}{a}\right), \\ &\quad \left(\begin{array}{l} m = 1, 2, 3, \dots \\ n = 1, 2, 3, \dots \end{array} \right) \end{aligned} \quad (35)$$

where e_{10} , e_{20} , e_{30} , e_{40} , A_n , a_n , α_n , B_n , and β_n are constant parameters, and $J_0(\alpha_n r/a)$ and $I_0(\beta_n r/a)$ are the Bessel function of the first kind, of order zero and the modified Bessel function of the first kind, of order zero, respectively.



By the boundary condition, Eq. (2),

$$e_{20} = e_{40} = 0 \quad (36)$$

The solution $\bar{\phi}_1(r, z)$, Eq. (35), may be split into two parts as

$$\bar{\phi}_1(r, z) = \bar{\phi}_{1A}(r, z) + \bar{\phi}_{1B}(r, z) \quad (37)$$

where $\bar{\phi}_{1A}(r, z)$ corresponds to the terms $e_{10} + e_{30} \frac{z}{a}$ and the series with coefficients A_n and represents the flow through the intersection between two liquids L_1 and L_2 , while $\bar{\phi}_{1B}(r, z)$ corresponds to the series with coefficients B_n and represents the flow through the side of the cylinder. However, both velocity potentials produce pressure on both the intersection and the side.

The boundary conditions for $\bar{\phi}_{1A}(r, z)$ and $\bar{\phi}_{1B}(r, z)$ are, respectively,

$$\left. \frac{\partial \bar{\phi}_{1A}}{\partial r} \right|_{r=a} = 0 \quad (38)$$

and

$$\left. \frac{\partial \bar{\phi}_{1B}}{\partial z} \right|_{z=b} = 0 \quad (39)$$

By the boundary condition, Eq. (38), we have

$$J_1(\alpha_n) = 0 \quad (40)$$

from which the constant parameter α_n can be determined. By the boundary condition, Eq. (39), we find

$$b_n \cos\left(\frac{\beta_n b}{a}\right) - \sin\left(\frac{\beta_n b}{a}\right) = 0 \quad (41)$$



from which

$$b_n = \tan\left(\frac{\beta_n b}{a}\right) \quad (42)$$

Define

$$e_{30} = A_0 \quad (43)$$

Substituting Eqs. (36), (42), and (43) into Eq. (35) gives

$$\begin{aligned} \bar{\Phi}_1(r, z) &= e_{10} + A_0 \frac{z}{a} \\ &+ \sum_{n=1}^{\infty} A_n \left[a_n \sinh\left(\frac{\alpha_n t}{a}\right) + \cosh\left(\frac{\alpha_n z}{a}\right) \right] J_0\left(\frac{\alpha_n r}{a}\right) \\ &+ \sum_{n=1}^{\infty} B_n \frac{\cos\left(\beta_n \frac{z-b}{a}\right)}{\cos\left(\frac{\beta_n b}{a}\right)} I_0\left(\frac{\beta_n r}{a}\right) \end{aligned} \quad (44)$$

The velocity potential $\Phi_1(r, z)$ is then obtained by substituting Eq. (44) into Eq. (34), as

$$\begin{aligned} \Phi_1(r, z) &= a^2 \omega \left\{ e_{10} + A_0 \frac{z}{a} + \sum_{n=1}^{\infty} B_n \frac{\cos\left(\beta_n \frac{z-b}{a}\right)}{\cos\left(\frac{\beta_n b}{a}\right)} I_0\left(\frac{\beta_n r}{a}\right) \right. \\ &\left. + \sum_{n=1}^{\infty} A_n \left[a_n \sinh\left(\frac{\alpha_n z}{a}\right) + \cosh\left(\frac{\alpha_n z}{a}\right) \right] J_0\left(\frac{\alpha_n r}{a}\right) \right\} \cos \omega t \end{aligned} \quad (45)$$

The constant parameters e_{10} , β_n , and a_n are obtained in terms of ω from the boundary condition, Eq. (3), as follows:



$$\left. \begin{aligned}
 e_{10} &= \frac{A_0}{a} \left(\frac{g}{\omega^2} - H \right) \\
 g \frac{\beta_n}{a} \tan \left(\beta_n \frac{H-b}{a} \right) + \omega^2 &= 0 \\
 A_n &= \frac{\omega^2 - g \frac{\alpha_n}{a} \tanh \left(\frac{\alpha_n h}{a} \right)}{g \frac{\alpha_n}{a} - \omega^2 \tanh \left(\frac{\alpha_n h}{a} \right)}
 \end{aligned} \right\} \quad (46)$$

The remaining unknown values, A_0 , A_n and B_n of Eq. (45), will be determined later in the eigenvalue problem.

2. Liquid L_2

For a simple harmonic motion, the velocity potential, Φ_2 , may be taken in the form

$$\Phi_2 = a^2 \omega \bar{\phi}_2(s, \theta) \cos \omega t \quad (47)$$

where $\bar{\phi}_2(s, \theta)$ is the velocity potential for steady flow.

By separation of variables, $\bar{\phi}_2(s, \theta)$ is obtained in the form

$$\bar{\phi}_2(\bar{s}, \theta) = C_0 + \frac{F_0}{\bar{s}} + \sum_{n=1}^{\infty} \left[C_n \bar{s} + F_n \bar{s}^{-(n+1)} \right] \left[P_n(\mu) + G_n Q_n(\mu) \right] \quad (48)$$

where $\mu = \cos \theta$.



In order to avoid the singularity at $s = 0$ so that the velocity potential $\bar{\phi}_2$ will remain finite, the coefficients F_0 , F_n and G_n are assumed zero. In this case, Eq. (48) becomes

$$\bar{\phi}_2(s, \theta) = C_0 + \sum_{n=0}^{\infty} C_n s^n P_n(\mu) \quad (49)$$

The remaining unknown values, C_0 and C_n will be determined later in the eigenvalue problem.

3. Circular Cylindrical Shell Partially Filled with Liquid L_1

Let the nondimensionalized displacement components \bar{u}_{li} and \bar{w}_{li} ($i = 1$ and 2), pressure \bar{p}_1 and axial coordinate \bar{z} be defined by the relations:

$$u_{li} = \frac{\rho_1 a^4 \omega^2}{D_1} \bar{u}_{li} \sin \omega t, \quad i = 1, 2, \dots \quad (50)$$

$$w_{li} = \frac{\rho_1 a^4 \omega^2}{D_1} \bar{w}_{li} \sin \omega t, \quad i = 1, 2, \dots \quad (51)$$

$$p_1 = \rho_1 a^2 \omega^2 \bar{p}_1 \sin \omega t, \quad (52)$$

$$z = a \bar{z} \quad (53)$$

$$t = \frac{1}{\omega} \bar{t} \quad (54)$$

Substituting these expressions into Eqs. (11) and (12) gives, respectively,



$$\frac{d^2 \bar{u}_{11}}{dz^2} + v \frac{d\bar{w}_{11}}{dz} = - \kappa_1 \omega^2 \bar{u}_{11} \quad (55)$$

$$v \frac{d\bar{u}_{11}}{dz} + \bar{w}_{11} = \kappa_1 \omega^2 \bar{w}_{11} + \bar{p}_{11} \quad (56)$$

where

$$\kappa_1 = \frac{\rho_1 h_1 a^2}{D_1} \quad (57)$$

Let the nondimensionalized liquid density $\bar{\rho}_{f1}$, velocity potentials $\bar{\phi}_1$ and $\bar{\Phi}_1$ be defined by the relations:

$$\rho_f = \rho_1 \bar{\rho}_{f1} \quad (58)$$

$$\phi_1 = a^2 \omega \bar{\phi}_1 \quad (59)$$

$$\bar{\Phi}_1 = \bar{\phi}_1 \cos \omega t \quad (60)$$

Substituting these expressions into Eq. (13) gives

$$p_1 = \rho_1 a^2 \omega^2 \bar{\rho}_{f1} \bar{\phi}_1 \sin \omega t \quad (61)$$

From Eqs. (52) and (61),

$$\bar{p}_1 = \bar{\rho}_{f1} \bar{\phi}_1 \quad (62)$$

Substituting Eqs. (50), (51) and (53) into Eqs. (14) and (15) gives

$$\frac{d\bar{u}_{11}}{dz} = \frac{d\bar{u}_{12}}{dz} \Delta \bar{u}_{1c} \quad \text{at } \bar{z} = \frac{b}{a} \quad (63)$$



and

$$\frac{d\bar{u}_{11}}{d\bar{z}} + v \bar{w}_{11} = 0 \quad \text{at } \bar{z} = \frac{H}{a} \quad (64)$$

where \bar{u}_{1c} is a constant to be determined later.

From Eq. (56) we find

$$\bar{w}_{11} = - \frac{v}{1 - \kappa_1^2 \omega^2} \frac{d\bar{u}_{11}}{d\bar{z}} + \frac{1}{1 - \kappa_1^2 \omega^2} \bar{p}_1 \quad (65)$$

Substituting Eq. (65) into Eq. (55) gives

$$\frac{d^2 \bar{u}_{11}}{d\bar{z}^2} + g_{11} \bar{u}_{11} = - g_{12} \frac{d\bar{p}_1}{d\bar{z}} \quad (66)$$

where

$$\left. \begin{aligned} g_{11} &= \frac{\kappa_1^2 \omega^2 (1 - \kappa_1^2 \omega^2)}{1 - v^2 - \kappa_1^2 \omega^2} \\ g_{12} &= \frac{v}{1 - v^2 - \kappa_1^2 \omega^2} \end{aligned} \right\} \quad (67)$$

The homogeneous solution of Eq. (66) is

$$\bar{u}_{11H} = C_{u1} \cos(\sqrt{g_{11}} \bar{z}) + C_{u2} \sin(\sqrt{g_{11}} \bar{z}) \quad (68)$$

where C_{u1} and C_{u2} are arbitrary constants.



The particular solution of Eq. (66) is assumed as

$$\bar{u}_{11P} = K_{11}(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) + K_{12}(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) \quad (69)$$

where

$$K_{11}(\bar{z}) = - \int^{\bar{z}} \frac{f_1(\bar{z}) U_{12}(\bar{z})}{W_1[U_{11}(\bar{z}), U_{12}(\bar{z})]} d\bar{z}$$

$$K_{12}(\bar{z}) = \int^{\bar{z}} \frac{f_1(\bar{z}) U_{11}(\bar{z})}{W_1[U_{11}(\bar{z}), U_{12}(\bar{z})]} d\bar{z}.$$

$$f_1(\bar{z}) = - g_{12} \frac{d \bar{p}_1}{d \bar{z}}$$

$$U_{11}(\bar{z}) = \cos\left(\sqrt{g_{11}} \bar{z}\right)$$

$$U_{12}(\bar{z}) = \sin\left(\sqrt{g_{11}} \bar{z}\right)$$

$$W_1[U_{11}(\bar{z}), U_{12}(\bar{z})] = \begin{vmatrix} U_{11}(\bar{z}) & U_{12}(\bar{z}) \\ U_{11}'(\bar{z}) & U_{12}'(\bar{z}) \end{vmatrix} \quad (70)$$

The Wronskian, W_1 , is

$$W_1[U_{11}(\bar{z}), U_{12}(\bar{z})] = \begin{vmatrix} \cos\left(\sqrt{g_{11}} \bar{z}\right) & \sin\left(\sqrt{g_{11}} \bar{z}\right) \\ -\sqrt{g_{11}} \sin\left(\sqrt{g_{11}} \bar{z}\right) & \sqrt{g_{11}} \cos\left(\sqrt{g_{11}} \bar{z}\right) \end{vmatrix}$$

$$= \sqrt{g_{11}} \left[\cos^2\left(\sqrt{g_{11}} \bar{z}\right) + \sin^2\left(\sqrt{g_{11}} \bar{z}\right) \right] = \sqrt{g_{11}} \quad (71)$$



Thus, from Eqs. (70), we obtain

$$\left. \begin{aligned} K_{11}(\bar{z}) &= \frac{g_{12}}{\sqrt{g_{11}}} \bar{p}_1(z) \sin\left(\sqrt{g_{11}} \bar{z}\right) - g_{12} \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \\ K_{12}(\bar{z}) &= - \frac{g_{12}}{\sqrt{g_{11}}} \bar{p}_1(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) - g_{12} \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \end{aligned} \right\} \quad (72)$$

Substituting Eqs. (72) into Eq. (69) and adding the result to Eq. (68), gives the general solution of Eq. (66) as

$$\begin{aligned} \bar{u}_{11} &= C_{u1} \cos\left(\sqrt{g_{11}} \bar{z}\right) + C_{u2} \sin\left(\sqrt{g_{11}} \bar{z}\right) \\ &\quad - g_{12} \left[\cos\left(\sqrt{g_{11}} \bar{z}\right) \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right. \\ &\quad \left. + \sin\left(\sqrt{g_{11}} \bar{z}\right) \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right] \end{aligned} \quad (73)$$

Differentiating Eq. (73) with respect to \bar{z} gives



$$\begin{aligned}
 \frac{d\bar{u}_{11}}{d\bar{z}} = & - C_{u1} \sqrt{g_{11}} \sin(\sqrt{g_{11}} \bar{z}) + C_{u2} \sqrt{g_{11}} \cos(\sqrt{g_{11}} \bar{z}) \\
 & + g_{12} \sqrt{g_{11}} \left[\sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\
 & \left. - \cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] - g_{12} \bar{p}_1(\bar{z}) \quad (74)
 \end{aligned}$$

Substituting Eq. (74) into Eq. (65) gives

$$\begin{aligned}
 \bar{w}_{11} = & \frac{v C_{u1}}{1 - \kappa_1 \omega^2} \sqrt{g_{11}} \sin(\sqrt{g_{11}} \bar{z}) - \frac{v C_{u2}}{1 - \kappa_1 \omega^2} \sqrt{g_{11}} \cos(\sqrt{g_{11}} \bar{z}) \\
 & - \frac{v g_{12} \sqrt{g_{11}}}{1 - \kappa_1 \omega^2} \left[\sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\
 & \left. - \cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] \\
 & + \frac{v g_{12}}{1 - \kappa_1 \omega^2} \bar{p}_1(\bar{z}) + \frac{1}{1 - \kappa_1 \omega^2} \bar{p}_1(\bar{z}) \quad (75)
 \end{aligned}$$

From Eqs. (74) and (75),



$$\begin{aligned}
 \frac{d\bar{u}_{11}}{dz} + v \bar{w}_{11} &= - \left(1 - \frac{v^2}{1 - \kappa_1 \omega^2} \right) c_{u1} \sqrt{g_{11}} \sin(\sqrt{g_{11}} \bar{z}) \\
 &\quad + \left(1 - \frac{v^2}{1 - \kappa_1 \omega^2} \right) c_{u2} \sqrt{g_{11}} \cos(\sqrt{g_{11}} \bar{z}) \\
 &\quad + \left(1 - \frac{v^2}{1 - \kappa_1 \omega^2} \right) g_{12} \sqrt{g_{11}} \left[\sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\
 &\quad \left. - \cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] \\
 &\quad - \left(1 - \frac{v^2}{1 - \kappa_1 \omega^2} \right) g_{12} \bar{p}_1(z) + \frac{v}{1 - \kappa_1 \omega^2} \bar{p}_1(\bar{z})
 \end{aligned} \tag{76}$$

By the boundary condition, Eq. (64),



$$\begin{aligned}
 & - \left(1 - \frac{v^2}{1-\kappa_1 \omega^2} \right) \sqrt{g_{11}} \sin\left(\frac{\sqrt{g_{11}} H}{a}\right) C_{u1} + \left(1 - \frac{v^2}{1-\kappa_1 \omega^2} \right) \sqrt{g_{11}} \cos\left(\frac{\sqrt{g_{11}} H}{a}\right) C_{u2} \\
 & = - \left(1 - \frac{v^2}{1-\kappa_1 \omega^2} \right) g_{12} \sqrt{g_{11}} \left[\sin\left(\frac{\sqrt{g_{11}} H}{a}\right) \left| \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right|_{\bar{z}=\frac{H}{a}} \right. \\
 & \quad \left. - \cos\left(\frac{\sqrt{g_{11}} H}{a}\right) \left| \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right|_{\bar{z}=\frac{H}{a}} \right] \\
 & + \left(1 - \frac{v^2}{1-\kappa_1 \omega^2} \right) g_{12} \bar{p}_1\left(\frac{H}{a}\right) - \frac{v}{1-\kappa_1 \omega^2} \bar{p}_1\left(\frac{H}{a}\right)
 \end{aligned} \tag{77}$$

By the boundary condition, Eq. (63),

$$\begin{aligned}
 & - \sqrt{g_{11}} \sin\left(\frac{\sqrt{g_{11}} b}{a}\right) C_{u1} + \sqrt{g_{11}} \cos\left(\frac{\sqrt{g_{11}} b}{a}\right) C_{u2} \\
 & = - g_{12} \sqrt{g_{11}} \left[\sin\left(\frac{\sqrt{g_{11}} b}{a}\right) \left| \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right|_{\bar{z}=\frac{b}{a}} \right. \\
 & \quad \left. - \cos\left(\frac{\sqrt{g_{11}} b}{a}\right) \left| \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right|_{\bar{z}=\frac{b}{a}} \right] \\
 & + g_{12} \bar{p}_1\left(\frac{b}{a}\right) + \bar{u}_{1c}
 \end{aligned} \tag{78}$$



$$\text{Let } - \left(1 - \frac{v^2}{1-\kappa_1 \omega^2} \right) \sqrt{g_{11}} \sin \left(\frac{\sqrt{g_{11}} H}{a} \right) = \ell_{11}$$

$$\left(1 - \frac{v^2}{1-\kappa_1 \omega^2} \right) \sqrt{g_{11}} \cos \left(\frac{\sqrt{g_{11}} H}{a} \right) = \ell_{12}$$

$$- \left(1 - \frac{v^2}{1-\kappa_1 \omega^2} \right) g_{12} \sqrt{g_{11}} \left[\sin \left(\frac{\sqrt{g_{11}} H}{a} \right) \middle| \int_{\bar{z}=H/a}^{\bar{z}} \bar{p}_1(\bar{z}) \cos \left(\sqrt{g_{11}} \bar{z} \right) d\bar{z} \right]$$

$$- \cos \left(\frac{\sqrt{g_{11}} H}{a} \right) \left| \int_{\bar{z}=H/a}^{\bar{z}} \bar{p}_1(\bar{z}) \sin \left(\sqrt{g_{11}} \bar{z} \right) d\bar{z} \right]$$

$$+ \left(1 - \frac{v^2}{1-\kappa_1 \omega^2} \right) g_{12} \bar{p}_1 \left(\frac{H}{a} \right) - \frac{v}{1-\kappa_1 \omega^2} \bar{p}_1 \left(\frac{H}{a} \right) = h_{11}$$

(79)

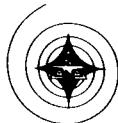
$$- \sqrt{g_{11}} \sin \left(\frac{\sqrt{g_{11}} b}{a} \right) = \ell_{13}$$

$$\sqrt{g_{11}} \cos \left(\frac{\sqrt{g_{11}} b}{a} \right) = \ell_{14}$$

$$- g_{12} \sqrt{g_{11}} \left[\sin \left(\frac{\sqrt{g_{11}} b}{a} \right) \middle| \int_{\bar{z}=b/a}^{\bar{z}} \bar{p}_1(\bar{z}) \cos \left(\sqrt{g_{11}} \bar{z} \right) d\bar{z} \right]$$

$$- \cos \left(\frac{\sqrt{g_{11}} b}{a} \right) \left| \int_{\bar{z}=b/a}^{\bar{z}} \bar{p}_1(\bar{z}) \sin \left(\sqrt{g_{11}} \bar{z} \right) d\bar{z} \right]$$

$$+ g_{12} \bar{p}_1 \left(\frac{b}{a} \right) = h_{12}$$



Eqs. (77) and (78) are then written as

$$\left. \begin{aligned} \ell_{11} c_{u1} + \ell_{12} c_{u2} &= h_{11} \\ \ell_{13} c_{u1} + \ell_{14} c_{u2} &= h_{12} + \bar{u}_{1c} \end{aligned} \right\} \quad (80)$$

from which

$$\left. \begin{aligned} c_{u1} &= \frac{\ell_{14} h_{11} - \ell_{12} h_{12} - \ell_{12} \bar{u}_{1c}}{\ell_{11} \ell_{14} - \ell_{12} \ell_{13}} \\ c_{u2} &= \frac{\ell_{11} h_{12} - \ell_{13} h_{11} + \ell_{11} \bar{u}_{1c}}{\ell_{11} \ell_{14} - \ell_{12} \ell_{13}} \end{aligned} \right\} \quad (81)$$

Substituting Eqs. (81) into Eqs. (73) through (75) gives

$$\begin{aligned} \bar{u}_{11} &= \frac{\ell_{14} h_{11} - \ell_{12} h_{12} - \ell_{12} \bar{u}_{1c}}{\ell_{11} \ell_{14} - \ell_{12} \ell_{13}} \cos(\sqrt{g_{11}} \bar{z}) \\ &\quad + \frac{\ell_{11} h_{12} - \ell_{13} h_{11} + \ell_{11} \bar{u}_{1c}}{\ell_{11} \ell_{14} - \ell_{12} \ell_{13}} \sin(\sqrt{g_{11}} \bar{z}) \\ &\quad - g_{12} \left[\cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\ &\quad \left. + \sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] \end{aligned} \quad (82)$$



$$\begin{aligned}
 \frac{d \bar{u}_{11}}{d \bar{z}} = & - \frac{\ell_{14} h_{11} - \ell_{12} h_{12} - \ell_{12} \bar{u}_{1c}}{\ell_{11} \ell_{14} - \ell_{12} \ell_{13}} \sqrt{g_{11}} \sin(\sqrt{g_{11}} \bar{z}) \\
 & + \frac{\ell_{11} h_{12} - \ell_{13} h_{11} + \ell_{11} \bar{u}_{1c}}{\ell_{11} \ell_{14} - \ell_{12} \ell_{13}} \sqrt{g_{11}} \cos(\sqrt{g_{11}} \bar{z}) \\
 & + g_{12} \sqrt{g_{11}} \left[\sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\
 & \left. - \cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] - g_{12} \bar{p}_1(\bar{z}) \quad (83)
 \end{aligned}$$

$$\begin{aligned}
 \bar{w}_{11} = & \frac{\nu \sqrt{g_{11}}}{1 - \kappa_1 \omega^2} \frac{\ell_{14} h_{11} - \ell_{12} h_{12} - \ell_{12} \bar{u}_{1c}}{\ell_{11} \ell_{14} - \ell_{12} \ell_{13}} \sin(\sqrt{g_{11}} \bar{z}) \\
 & - \frac{\nu \sqrt{g_{11}}}{1 - \kappa_1 \omega^2} \frac{\ell_{11} h_{12} - \ell_{13} h_{11} + \ell_{11} \bar{u}_{1c}}{\ell_{11} \ell_{14} - \ell_{12} \ell_{13}} \cos(\sqrt{g_{11}} \bar{z}) \\
 & - \frac{\nu g_{12} \sqrt{g_{11}}}{1 - \kappa_1 \omega^2} \left[\sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\
 & \left. - \cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] \\
 & + \frac{\nu g_{12}}{1 - \kappa_1 \omega^2} \bar{p}_1(\bar{z}) + \frac{1}{1 - \kappa_1 \omega^2} \bar{p}_1(\bar{z}) \quad (84)
 \end{aligned}$$



4. Circular Conical Shell

Let the nondimensionalized displacement components \bar{u}_2 and \bar{w}_2 , pressure $\bar{p}_{2\alpha}$ and generatrix \bar{s} be denoted by the relations:

$$u_2 = \frac{\rho_2 a^4 \omega^2}{D_2} \bar{u}_2 \sin \omega t \quad (85)$$

$$w_2 = \frac{\rho_2 a^4 \omega^2}{D_2} \bar{w}_2 \sin \omega t \quad (86)$$

$$p_{2\alpha} = \rho_2 a^2 \omega^2 \bar{p}_{2\alpha} \sin \omega t \quad (87)$$

$$s = a \bar{s} \quad (88)$$

where $P_{2\alpha}$ is the pressure acting at the conical shell surface.

Substituting these expressions into Eqs. (27) and (28) gives, respectively,

$$\bar{s} \frac{d^2 \bar{u}_2}{ds^2} + \frac{du_2}{ds} - \frac{\bar{u}_2}{\bar{s}} + \nu \frac{d\bar{w}_2}{ds} \cot \alpha - \frac{\bar{w}_2}{\bar{s}} \cot \alpha = - \kappa_2 \omega^2 \bar{s} \bar{u}_2 \quad (89)$$

$$\nu \frac{d\bar{u}_2}{ds} + \frac{\bar{u}_2}{\bar{s}} + \frac{\bar{w}_2}{\bar{s}} \cot \alpha = \kappa_2 \omega^2 \bar{s} \bar{w}_2 \tan \alpha - \bar{s} \bar{p}_{2\alpha} \tan \alpha \quad (90)$$

where

$$\kappa_2 = \frac{\rho_2 h_2 a^2}{D_2}$$

Let the nondimensionalized liquid density $\bar{\rho}_{f2}$ and velocity potentials $\bar{\phi}_2$ and $\bar{\Phi}_2$ be defined by the relations:



$$\rho_f = \rho_2 \bar{\rho}_{f2} \quad (91)$$

$$\phi_2 = a^2 \omega \bar{\phi}_2 \quad (92)$$

$$\Phi_2 = \bar{\phi}_2 \cos \omega t \quad (93)$$

Substituting these expressions into Eq. (26) gives

$$p_2 = \rho_2 a^2 \omega^2 \bar{\rho}_{f2} \bar{\phi}_2 \sin \omega t \quad (94)$$

From Eqs. (87) and (94)

$$\bar{p}_2 = \bar{\rho}_{f2} \bar{\phi}_2 \quad (95)$$

Substituting Eqs. (85) and (86) into boundary conditions (29) and (30) gives

$$\bar{u}_2 = \text{finite} \quad \text{at } \bar{s} = 0 \quad (96)$$

and

$$\bar{u}_2 - \bar{w}_2 \tan \alpha = 0 \quad \text{at } \bar{s} = \frac{\ell}{a} \quad (97)$$

where

$$\ell = \sqrt{a^2 + b^2}.$$

It will be indicated below how the two Equations (89) and (90) involving two dependent variables, \bar{u}_2 and \bar{w}_2 , can be reduced to a single differential equation of \bar{u}_2 .



First, from Eq. (90)

$$\bar{w}_2 = - \frac{v \bar{s} \frac{d\bar{u}_2}{d\bar{s}} + \bar{u}_2 + \bar{s}^2 \bar{p}_{2\alpha} \tan\alpha}{\cot\alpha - \kappa_2^2 \omega^{2-2} \bar{s} \tan\alpha} \quad (98)$$

Second, substitute Eq. (98) into Eq. (89),

$$\begin{aligned}
 & \left[1 - v^2 - (2-v^2) \kappa_2^2 \omega^{2-2} \bar{s} \tan^2\alpha + \kappa_2^2 \omega^{4-4} \bar{s} \tan^4\alpha \right] \bar{s}^2 \frac{d^2 \bar{u}_2}{d\bar{s}^2} \\
 & + \left[1 - v^2 - (2+v^2) \kappa_2^2 \omega^{2-2} \bar{s} \tan^2\alpha + \kappa_2^2 \omega^{4-4} \bar{s} \tan^4\alpha \right] \bar{s} \frac{d\bar{u}_2}{d\bar{s}} \\
 & + \left[(1-2v + \cot^2\alpha) \kappa_2^2 \omega^{2-2} \bar{s} \tan^2\alpha - (2 + \tan^2\alpha) \kappa_2^2 \omega^{4-4} \bar{s} \tan^2\alpha \right. \\
 & \left. + \kappa_2^3 \omega^{6-6} \bar{s} \tan^4\alpha \right] \bar{u}_2 \\
 & = \left(-\bar{s}^2 + \kappa_2^2 \omega^{2-4} \bar{s} \tan^2\alpha + 2v\bar{s}^2 \right) \bar{p}_{2\alpha} \tan\alpha \\
 & + \left(1 - \kappa_2^2 \omega^{2-2} \bar{s} \tan^2\alpha \right) v \bar{s}^3 \frac{d\bar{p}_{2\alpha}}{d\bar{s}} \tan\alpha \quad (99)
 \end{aligned}$$



Simplifying Eq. (99) for the case $\kappa_2 \omega^2 \bar{s}^2 \tan^2 \alpha \ll 1$ gives

$$\frac{d^2 \bar{u}_2}{d\bar{s}^2} + \frac{1}{\bar{s}} \frac{d\bar{u}_2}{d\bar{s}} - \lambda^2 \bar{u}_2 = \bar{f}_2(\bar{s}) \quad (100)$$

where

$$\left. \begin{aligned} \lambda^2 &= \frac{1 - 2v + \cot^2 \alpha}{1 - v^2} \kappa_2 \omega^2 \tan^2 \alpha \\ \bar{f}_2(\bar{s}) &= - \frac{1 - 2v}{1 - v^2} \bar{p}_{2\alpha} \tan \alpha + \frac{v}{1 - v^2} \bar{s} \frac{d\bar{p}_{2\alpha}}{d\bar{s}} \tan \alpha \end{aligned} \right\} \quad (101)$$

Multiplying both sides of Eq. (100) by $\bar{s} J_0(\mu_i \bar{s})$, and integrating from 0 to $\bar{\ell}$ gives

$$\int_0^{\bar{\ell}} \bar{s} \left(\frac{d^2 \bar{u}_2}{d\bar{s}^2} + \frac{1}{\bar{s}} \frac{d\bar{u}_2}{d\bar{s}} \right) J_0(\mu_i \bar{s}) d\bar{s} + \lambda^2 \bar{U}_2 = \bar{F}_2 \quad (102)$$

where \bar{U}_2 and \bar{F}_2 are the finite Hankel transforms of \bar{u}_2 and \bar{f}_2 , respectively, and defined as

$$\bar{U}_2(\mu_i) = \int_0^{\bar{\ell}} \bar{s} \bar{u}_2(\bar{s}) J_0(\mu_i \bar{s}) d\bar{s} \quad (103)$$

$$\bar{F}_2(\mu_i) = \int_0^{\bar{\ell}} \bar{s} \bar{f}_2(\bar{s}) J_0(\mu_i \bar{s}) d\bar{s} \quad (104)$$

where

$$\bar{\ell} = \sqrt{1 + \left(\frac{b}{a} \right)^2} \quad (105)$$



The first term on the left-hand side of Eq. (102) is written as

$$\begin{aligned}
 & \int_0^{\bar{\ell}} \bar{s} \left(\frac{d^2 \bar{u}_2}{ds^2} + \frac{1}{\bar{s}} \frac{d\bar{u}_2}{ds} \right) J_0(\mu_i \bar{s}) ds \\
 &= \int_0^{\bar{\ell}} \frac{d}{ds} \left(\bar{s} \frac{d\bar{u}_2}{ds} \right) J_0(\mu_i \bar{s}) ds \\
 &= J_0(\mu_i \bar{s}) \bar{s} \frac{d\bar{u}_2}{ds} \Big|_0^{\bar{\ell}} + \mu_i \int_0^{\bar{\ell}} \bar{s} J_1(\mu_i \bar{s}) \frac{d\bar{u}_2}{ds} ds \quad (106)
 \end{aligned}$$

Recognizing that $d\bar{u}_2/d\bar{s}$ must be finite from physical considerations, the first term on the right-hand side of Eq. (106) becomes

$$J_0(\mu_i \bar{s}) \bar{s} \frac{d\bar{u}_2}{ds} \Big|_0^{\bar{\ell}} = J_0(\mu_i \bar{\ell}) \bar{\ell} \frac{d\bar{u}_2}{ds} \Big|_{\bar{\ell}} \quad (107)$$

The second term on the right-hand side of Eq. (106) can be again integrated by parts as

$$\begin{aligned}
 & \mu_i \int_0^{\bar{\ell}} \bar{s} J_1(\mu_i \bar{s}) \frac{d\bar{u}_2}{ds} ds \\
 &= \mu_i \bar{u}_2 \bar{s} J_1(\mu_i \bar{s}) \Big|_0^{\bar{\ell}} - \mu_i^2 \int_0^{\bar{\ell}} \bar{s} J_0(\mu_i \bar{s}) \bar{u}_2 ds \quad (108)
 \end{aligned}$$

From the boundary condition, Eq. (29), the first term on the right-hand side of Eq. (108) becomes

$$\mu_i \bar{u}_2 \bar{s} J_1(\mu_i \bar{s}) \Big|_0^{\bar{\ell}} = \mu_i J_1(\mu_i \bar{\ell}) \bar{\ell} \bar{u}_2 \Big|_{\bar{\ell}} \quad (109)$$



Substituting Eqs. (107) through (109) into Eq. (106) and using the definition, Eq. (103), gives

$$\int_0^{\bar{\ell}} \bar{s} \left(\frac{d^2 \bar{u}_2}{d\bar{s}^2} + \frac{1}{\bar{s}} \frac{d\bar{u}_2}{d\bar{s}} \right) J_0(\mu_i \bar{s}) d\bar{s}$$

$$= J_0(\mu_i \bar{\ell}) \bar{\ell} \left[\frac{d\bar{u}_2}{d\bar{s}} \right]_{\bar{\ell}} + \mu_i J_1(\mu_i \bar{\ell}) \bar{\ell} \bar{u}_2 \Big|_{\bar{\ell}} - \mu_i^2 \bar{u}_2 \quad (110)$$

By choosing the μ_i so that

$$J_0(\mu_i \bar{\ell}) \bar{\ell} \left[\frac{d\bar{u}_2}{d\bar{s}} \right]_{\bar{\ell}} + \mu_i J_1(\mu_i \bar{\ell}) \bar{\ell} \bar{u}_2 \Big|_{\bar{\ell}} = 0 \quad (111)$$

Eq. (110) may be written as

$$\int_0^{\bar{\ell}} \bar{s} \left(\frac{d^2 \bar{u}_2}{d\bar{s}^2} + \frac{1}{\bar{s}} \frac{d\bar{u}_2}{d\bar{s}} \right) J_0(\mu_i \bar{s}) d\bar{s} = -\mu_i^2 \bar{u}_2 \quad (112)$$

Substituting this expression into Eq. (102) gives

$$\bar{u}_2 = \frac{1}{\lambda^2 - \mu_i^2} \bar{F}_2 \quad (113)$$

By applying an inversion formula of finite Hankel transform to Eq. (113),

$$\bar{u}_2 = \frac{2}{\bar{\ell}^2} \sum_i \frac{\mu_i^2 \bar{F}_2}{(\sigma^2 + \mu_i^2)(\lambda^2 - \mu_i^2)} \cdot \frac{J_0(\mu_i \bar{s})}{[J_0(\mu_i \bar{\ell})]^2} \quad (114)$$



where

$$\sigma = - \frac{\left[\frac{d\bar{u}_2}{ds} \right]_{\bar{l}}}{\bar{u}_2|_{\bar{l}}} \quad (115)$$

and the sum is taken over all the positive roots of Eq. (111).

The numerical value of σ can be determined by an iteration scheme. An estimated value of σ is first chosen. The numerical value of $\left[\frac{d\bar{u}_2}{ds} \right] / \bar{u}_2$ at $\bar{s} = \bar{l}$ is evaluated next. This process is repeated for successive values of σ until Eq. (115) is satisfied. An approximate value of σ is, however, determined by assuming that the resultant of the hydrostatic pressure acting on the conical shell surface in its tangential direction is negligibly small compared with the resultant of the dynamic pressure. For this case, the stress resultant, N_s , of Eq. (22) can be approximately put equal to zero at $\bar{s} = \bar{l}$, i.e.,

$$\left[\frac{d\bar{u}_2}{ds} \right]_{\bar{l}} + \frac{v}{l} \left(u_2 \left[\frac{}{} \right]_{\bar{l}} + \bar{w}_2 \left[\frac{}{} \right]_{\bar{l}} \cot \alpha \right) = 0 \quad (116)$$

Substituting the boundary condition, Eq. (97) into the above equation gives

$$\left[\frac{d\bar{u}_2}{ds} \right]_{\bar{l}} + \frac{v}{l} \bar{u}_2 \left[\frac{}{} \right]_{\bar{l}} \operatorname{cosec}^2 \alpha = 0 \quad (117)$$

from which

$$\frac{\left[\frac{d\bar{u}_2}{ds} \right]_{\bar{l}}}{\bar{u}_2|_{\bar{l}}} = - \frac{v}{l} \operatorname{cosec}^2 \alpha \quad (118)$$

Substituting this expression into Eq. (115) gives



$$\sigma = \frac{v}{\ell} \operatorname{cosec}^2 \alpha \quad (119)$$

Substituting Eq. (114) into Eq. (98) gives

$$\bar{w}_2 = \frac{2 \tan \alpha}{\ell^2} \sum_i \frac{\mu_i^2 \bar{F}_2}{(\sigma^2 + \mu_i^2)(\lambda^2 - \mu_i^2)} \cdot \frac{v \bar{s} \mu_i J_1(\mu_i \bar{s}) - J_0(\mu_i \bar{s})}{[J_0(\mu_i \ell)]^2} - \bar{s}^2 \bar{p}_{2\alpha} \tan^2 \alpha \quad (120)$$

where

$$\left. \begin{aligned} \bar{F}_2 &= \int_0^{\bar{\ell}} \bar{s} \bar{f}_2(\bar{s}) J_0(\mu_i \bar{s}) ds \\ \bar{p}_{2\alpha} &= \bar{\rho}_{f2} \left[C_0 + \sum_{n=1}^{\infty} C_n \bar{s}^n P_n(\mu) \right] \end{aligned} \right\} \quad (121)$$

Substituting the expression, $\bar{p}_{2\alpha}$ of Eqs. (121) into the second equation of Eqs. (101) gives

$$\begin{aligned} \bar{f}_2(\bar{s}) &= - \frac{1 - 2v}{1 - v^2} \bar{\rho}_{f2} C_0 \tan \alpha \\ &\quad - \frac{1 - 2v}{1 - v^2} \bar{\rho}_{f2} \sum_{n=1}^{\infty} C_n \bar{s}^n P_n(\mu) \tan \alpha \\ &\quad + \frac{v}{1 - v^2} \bar{\rho}_{f2} \sum_{n=1}^{\infty} n C_n \bar{s}^n P_n(\mu) \tan \alpha \end{aligned} \quad (122)$$

5. Circular Cylindrical Shell Filled with Liquid L_2

Referring to Eqs. (73) and (75), the general solutions of two simultaneous equations, Eqs. (31) and (32) are obtained as

$$\begin{aligned} \bar{u}_{12} = & C_{u3} \cos(\sqrt{g_{11}} \bar{z}) + C_{u4} \sin(\sqrt{g_{11}} \bar{z}) \\ & - g_{12} \left[\cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\ & \left. + \sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] \end{aligned} \quad (123)$$

$$\begin{aligned} \bar{w}_{12} = & \frac{\nu C_{u3}}{1 - \kappa_1 \omega^2} \sqrt{g_{11}} \sin(\sqrt{g_{11}} \bar{z}) - \frac{\nu C_{u4}}{1 - \kappa_1 \omega^2} \sqrt{g_{11}} \cos(\sqrt{g_{11}} \bar{z}) \\ & - \frac{\nu g_{12} \sqrt{g_{11}}}{1 - \kappa_1 \omega^2} \left[\sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\ & \left. - \cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] \\ & + \frac{\nu g_{12}}{1 - \kappa_1 \omega^2} \bar{p}_2(\bar{z}) + \frac{1}{1 - \kappa_1 \omega^2} \bar{p}_2(\bar{z}) \end{aligned} \quad (124)$$

where $\bar{p}_2(\bar{z})$ is obtained from Eqs. (49) and (95) and the coordinate transformation as



$$\bar{p}_2(z) = \bar{\rho}_{f2} \left\{ C_0 + \sum_{n=1}^{\infty} C_n \left[1 + (\bar{b} - \bar{z})^2 \right]^{\frac{n}{2}} P_n(\mu_{\bar{z}}) \right\} \quad (125)$$

$$\mu_{\bar{z}} = \frac{\bar{b} - \bar{z}}{\sqrt{1 + (\bar{b} - \bar{z})^2}}$$

The boundary conditions are

$$\frac{d\bar{u}_{12}}{d\bar{z}} = \bar{u}_{1c} \quad \text{at} \quad \bar{z} = \frac{b}{a} \quad (126)$$

$$\bar{u}_{12} = 0 \quad \text{at} \quad \bar{z} = 0 \quad (127)$$

From Eq. (123),

$$\begin{aligned} \frac{d\bar{u}_{12}}{d\bar{z}} &= -C_{u3} \sqrt{g_{11}} \sin\left(\sqrt{g_{11}} \bar{z}\right) + C_{u4} \sqrt{g_{11}} \cos\left(\sqrt{g_{11}} \bar{z}\right) \\ &\quad + g_{12} \sqrt{g_{11}} \left[\sin\left(\sqrt{g_{11}} \bar{z}\right) \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right. \\ &\quad \left. - \cos\left(\sqrt{g_{11}} \bar{z}\right) \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right] - g_{12} \bar{p}_2(\bar{z}) \end{aligned} \quad (128)$$

By the boundary condition, Eq. (126)



$$- C_{u3} \sqrt{g_{11}} \sin\left(\frac{b \sqrt{g_{11}}}{a}\right) + C_{u4} \sqrt{g_{11}} \cos\left(\frac{b \sqrt{g_{11}}}{a}\right)$$

$$+ g_{12} \sqrt{g_{11}} \left[\sin\left(\frac{b \sqrt{g_{11}}}{a}\right) \Big| \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \Big|_{\bar{z}=\frac{b}{a}} \right]$$

$$- \cos\left(\frac{b \sqrt{g_{11}}}{a}\right) \Big| \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \Big|_{\bar{z}=\frac{b}{a}} \Big]$$

$$- g_{12} \bar{p}_2\left(\frac{b}{a}\right) = \bar{u}_{1c} \quad (129)$$

By the boundary condition, Eq. (127)

$$C_{u3} = g_{12} \Big| \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \Big|_{\bar{z}=0} \quad (130)$$

Let



$$-\sqrt{g_{11}} \sin\left(\frac{b\sqrt{g_{11}}}{a}\right) = \ell_{23}$$

$$\sqrt{g_{11}} \cos\left(\frac{b\sqrt{g_{11}}}{a}\right) = \ell_{24}$$

$$g_{12} \left| \int_{z=0}^{\bar{z}} \bar{p}_2(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right| = h_{21}$$

(131)

$$- g_{12} \sqrt{g_{11}} \left[\sin\left(\frac{b\sqrt{g_{11}}}{a}\right) \right] \left| \int_{z=\frac{b}{a}}^{\bar{z}} \bar{p}_2(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right|$$

$$- \cos\left(\frac{b\sqrt{g_{11}}}{a}\right) \left| \int_{z=\frac{b}{a}}^{\bar{z}} \bar{p}_2(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right| + g_{12} \bar{p}_2\left(\frac{b}{a}\right) = h_{22}$$

From Eqs. (129) through (131),

$$\left. \begin{aligned} c_{u3} &= h_{21} \\ c_{u4} &= \frac{h_{22} - \ell_{23}h_{21} + \bar{u}_{1c}}{\ell_{24}} \end{aligned} \right\} \quad (132)$$

Substituting Eqs. (132) into Eqs. (123), (124), and (127) gives, respectively,



$$\begin{aligned}
 \bar{u}_{12} = & h_{21} \cos\left(\sqrt{g_{11}} \bar{z}\right) + \frac{h_{22} - \ell_{23}h_{21} + \bar{u}_{1c}}{\ell_{24}} \sin\left(\sqrt{g_{11}} \bar{z}\right) \\
 & - g_{12} \left[\cos\left(\sqrt{g_{11}} \bar{z}\right) \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right. \\
 & \left. + \sin\left(\sqrt{g_{11}} \bar{z}\right) \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right] \quad (133)
 \end{aligned}$$

$$\begin{aligned}
 \bar{w}_{12} = & \frac{v}{1 - \kappa_1 \omega^2} h_{21} \sqrt{g_{11}} \sin\left(\sqrt{g_{11}} \bar{z}\right) \\
 & - \frac{v}{1 - \kappa_1 \omega^2} \frac{h_{22} - \ell_{23}h_{21} + \bar{u}_{1c}}{\ell_{24}} \sqrt{g_{11}} \cos\left(\sqrt{g_{11}} \bar{z}\right) \\
 & - \frac{v g_{12} \sqrt{g_{11}}}{1 - \kappa_1 \omega^2} \left[\sin\left(\sqrt{g_{11}} \bar{z}\right) \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right. \\
 & \left. - \cos\left(\sqrt{g_{11}} \bar{z}\right) \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \right] \\
 & + \frac{v g_{12}}{1 - \kappa_1 \omega^2} \bar{p}_2(\bar{z}) + \frac{1}{1 - \kappa_1 \omega^2} \bar{p}_2(\bar{z}) \quad (134)
 \end{aligned}$$



$$\frac{d\bar{u}_2}{dz} = - h_{21} \sqrt{g_{11}} \sin(\sqrt{g_{11}} \bar{z}) + \frac{h_{22} - \ell_{23} h_{21} + \bar{u}_{1c}}{\ell_{24}} \sqrt{g_{11}} \cos(\sqrt{g_{11}} \bar{z})$$

$$\begin{aligned}
 & + g_{12} \sqrt{g_{11}} \left[\sin(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right. \\
 & \left. - \cos(\sqrt{g_{11}} \bar{z}) \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right] - g_{12} \bar{p}_2(\bar{z}) \quad (135)
 \end{aligned}$$

From the solutions Eqs. (82) and (133) and the first boundary condition of Eqs. (14),

$$\begin{aligned}
 \bar{u}_{1c} \cot \sqrt{g_{11}} (\bar{H} - \bar{b}) &= \bar{u}_{1c} \tan(\sqrt{g_{11}} \bar{b}) \\
 &= \frac{h_{11}}{\left(1 - \frac{v^2}{1 - \kappa_1 \omega^2} \right) \sin \sqrt{g_{11}} (\bar{H} - \bar{b})} - h_{12} \cot \sqrt{g_{11}} (\bar{H} - \bar{b}) + \frac{\sqrt{g_{11}} h_{21}}{\cos(\sqrt{g_{11}} \bar{b})} \\
 &+ h_{22} \tan(\sqrt{g_{11}} \bar{b}) + g_{12} \sqrt{g_{11}} \left\{ \left[f_{p11}(\bar{b}) - f_{p21}(\bar{b}) \right] \cos(\sqrt{g_{11}} \bar{b}) \right. \\
 &\left. + \left[f_{p12}(\bar{b}) - f_{p22}(\bar{b}) \right] \sin(\sqrt{g_{11}} \bar{b}) \right\} \quad (136)
 \end{aligned}$$



where f_{p11} , f_{p12} , f_{p21} , and f_{p22} are defined, respectively, as

$$\left. \begin{aligned} f_{p11}(\bar{b}) &= f_{p11}(\bar{z}) \Big|_{\bar{z}=\bar{b}} = \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \Big|_{\bar{z}=\bar{b}} \\ f_{p12}(\bar{b}) &= f_{p12}(\bar{z}) \Big|_{\bar{z}=\bar{b}} = \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \Big|_{\bar{z}=\bar{b}} \\ f_{p21}(\bar{b}) &= f_{p21}(\bar{z}) \Big|_{\bar{z}=\bar{b}} = \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \Big|_{\bar{z}=\bar{b}} \\ f_{p22}(\bar{b}) &= f_{p22}(\bar{z}) \Big|_{\bar{z}=\bar{b}} = \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \Big|_{\bar{z}=\bar{b}} \end{aligned} \right\} \quad (137)$$

From which

$$\begin{aligned} \bar{u}_{1c} &= \frac{\sin \sqrt{g_{11}} (\bar{H}-\bar{b}) \cos(\sqrt{g_{11}} \bar{b})}{\cos(\sqrt{g_{11}} \bar{H})} \\ &\times \left[\frac{h_{11}}{\left(1 - \frac{v^2}{1 - \kappa_1 \omega^2} \right) \sin \sqrt{g_{11}} (\bar{H}-\bar{b})} - h_{12} \cot \sqrt{g_{11}} (\bar{H}-\bar{b}) \right. \\ &+ \frac{\sqrt{g_{11}} h_{21}}{\cos(\sqrt{g_{11}} \bar{b})} + h_{22} \tan(\sqrt{g_{11}} \bar{b}) \\ &+ g_{12} \sqrt{g_{11}} \left\{ [f_{p11}(\bar{b}) - f_{p21}(\bar{b})] \cos(\sqrt{g_{11}} \bar{b}) \right. \\ &+ \left. \left. [f_{p12}(\bar{b}) - f_{p22}(\bar{b})] \sin(\sqrt{g_{11}} \bar{b}) \right\} \right] \end{aligned} \quad (138)$$



IV. MATRIX EIGENVALUE PROBLEM

The eigenvalue problem for determining the natural frequencies and corresponding mode shapes of the coupled liquid-shell system is formulated using five equations which were obtained as a result of imposing the necessary boundary conditions. These five equations are Eqs. (4), (5), (6), (8) and (9). Other than the trivial solution in which all the unknown constants are zero, it may be seen that the series in these five equations cannot be satisfied term by term. Hence, some approximate numerical method of solution must be used. The approximate method used herein to formulate the eigenvalue problem is based on the least squared error technique. Although the series representations for the velocity potentials, Eqs. (44) and (49), would represent an exact solution if an infinite number of terms were used, as a practical matter, the series must be truncated. Since the chosen velocity potentials satisfy the differential equations and non-interface boundary conditions exactly term by term, it is the five interface conditions that suffer. These conditions can be satisfied only approximately. The functional errors can then be defined as:

$$\bar{\varepsilon}_1 = \left[\frac{\partial \bar{\phi}_1}{\partial \bar{r}} \right]_{\bar{r}=1} - \beta_{K1} \bar{w}_{11} \quad (139)$$

$$\bar{\varepsilon}_2 = \left[\frac{\partial \bar{\phi}_1}{\partial \bar{z}} \right]_{\bar{z}=\frac{b}{a}} - \frac{1}{s} \left[\frac{\partial \bar{\phi}_2}{\partial \theta} \right]_{\theta=\frac{\pi}{2}} \quad (140)$$

$$\bar{\varepsilon}_3 = \left[\bar{\phi}_1 \right]_{\bar{z}=\frac{b}{a}} - \left[\bar{\phi}_2 \right]_{\theta=\frac{\pi}{2}} \quad (141)$$

$$\bar{\varepsilon}_4 = \left[\frac{1}{s} \frac{\partial \bar{\phi}_2}{\partial \theta} \right]_{\theta=\alpha} - \beta_{K2} \bar{w}_2 \quad (142)$$



$$\bar{\epsilon}_5 = \left[\frac{1}{\sqrt{1 + \left(\frac{b}{a} - \bar{z} \right)^2}} \frac{\partial \bar{\phi}_2}{\partial \bar{s}} + \frac{\frac{b}{a} - \bar{z}}{1 + \left(\frac{b}{a} - \bar{z} \right)^2} \frac{\partial \bar{\phi}_2}{\partial \theta} \right] - \beta_{K1} \bar{w}_{12}$$

$$\left. \begin{aligned} \bar{s} &= \sqrt{1 + \left(\frac{b}{a} - \bar{z} \right)^2} \\ \theta &= \tan^{-1} \frac{1}{\frac{b}{a} - \bar{z}} \end{aligned} \right\} \quad (143)$$

where $\bar{\epsilon}_k$ ($k = 1, 2, 3, 4$, and 5) is defined by the relation:

$$\left. \begin{aligned} \epsilon_k &= a \omega \bar{\epsilon}_k \cos \omega t & (k=1,2,4,5) \\ \epsilon_3 &= -a^2 \omega^2 \bar{\epsilon}_3 \sin \omega t \end{aligned} \right\} \quad (144)$$

and

$$\left. \begin{aligned} \beta_{K1} &= \frac{\rho_1 a^3 \omega^2}{D_1} \\ \beta_{K2} &= \frac{\rho_2 a^3 \omega^2}{D_2} \end{aligned} \right\} \quad (145)$$

With these expressions, the total integrated squared error, S_T , over the boundaries involved can be expressed as

$$S_T = S_1 + S_2 + S_3 + S_4 + S_5 \quad (146)$$

where



where

$$\left. \begin{aligned}
 S_1 &= 2\pi a^2 \int_{\bar{b}}^{\bar{H}} \bar{\varepsilon}_1^2 (A_n, B_n, C_{\hat{n}}, C_{\bar{n}}, \bar{z}) d\bar{z} \\
 S_2 &= 2\pi a^2 \int_0^1 \bar{\varepsilon}_2^2 (A_n, C_{\bar{n}}, \bar{r}) \bar{r} d\bar{r} \\
 S_3 &= 2\pi a^2 \int_0^1 \bar{\varepsilon}_3^2 (A_n, B_n, C_{\hat{n}}, \bar{r}) \bar{r} d\bar{r} \\
 S_4 &= 2\pi a^2 \int_0^{\bar{\ell}} \bar{\varepsilon}_4^2 (C_{\hat{n}}, C_{\bar{n}}, \bar{s}) \bar{s} \sin \omega \bar{s} d\bar{s} \\
 S_5 &= 2\pi a^2 \int_0^{\bar{b}} \bar{\varepsilon}_5^2 (A_n, B_n, C_{\hat{n}}, C_{\bar{n}}, \bar{z}) d\bar{z}
 \end{aligned} \right\} \quad (147)$$

$n = 0, 1, 2, \dots, N$
 $\hat{n} = 2n$
 $\bar{n} = 2n + 1$

but $B_0 = 0$, and $\bar{H} = H/a$, $\bar{\ell} = \ell/a$, and $\bar{b} = b/a$.

The frequency, ω and constants, A_n , B_n , $C_{\hat{n}}$, and $C_{\bar{n}}$ are then determined by minimizing the total integrated squared error S_T . The conditions for this minimum are

$$\left. \begin{aligned}
 \frac{\partial S_T}{\partial A_m} &= 0 \\
 \frac{\partial S_T}{\partial B_m} &= 0 \\
 \frac{\partial S_T}{\partial C_{\hat{m}}} &= 0 \\
 \frac{\partial S_T}{\partial C_{\bar{m}}} &= 0
 \end{aligned} \right\} \quad (148)$$

$m = 0, 1, 2, \dots, N$
 $\hat{m} = 2m$
 $\bar{m} = 2m + 1$



Substituting Eq. (146) into Eqs. (148) gives

$$\left. \begin{aligned} \frac{\partial S_1}{\partial A_m} + \frac{\partial S_2}{\partial A_m} + \frac{\partial S_3}{\partial A_m} + \frac{\partial S_5}{\partial A_m} &= 0 \\ \frac{\partial S_1}{\partial B_m} + \frac{\partial S_3}{\partial B_m} + \frac{\partial S_5}{\partial B_m} &= 0 \\ \frac{\partial S_1}{\partial C_{\hat{m}}} + \frac{\partial S_3}{\partial C_{\hat{m}}} + \frac{\partial S_4}{\partial C_{\hat{m}}} + \frac{\partial S_5}{\partial C_{\hat{m}}} &= 0 \\ \frac{\partial S_1}{\partial C_{\bar{m}}} + \frac{\partial S_2}{\partial C_{\bar{m}}} + \frac{\partial S_4}{\partial C_{\bar{m}}} + \frac{\partial S_5}{\partial C_{\bar{m}}} &= 0 \end{aligned} \right\} \quad (149)$$

Substituting Eqs. (147) into Eqs. (149) gives

$$\left. \begin{aligned} \int_{\bar{b}}^{\bar{H}} \bar{\epsilon}_1 \frac{\partial \bar{\epsilon}_1}{\partial A_m} d\bar{z} + \int_0^1 \bar{\epsilon}_2 \frac{\partial \bar{\epsilon}_2}{\partial A_m} \bar{r} d\bar{r} + \int_0^1 \bar{\epsilon}_3 \frac{\partial \bar{\epsilon}_3}{\partial A_m} \bar{r} d\bar{r} + \int_0^{\bar{b}} \bar{\epsilon}_5 \frac{\partial \bar{\epsilon}_5}{\partial A_m} d\bar{z} &= 0 \\ \int_{\bar{b}}^{\bar{H}} \bar{\epsilon}_1 \frac{\partial \bar{\epsilon}_1}{\partial B_m} d\bar{z} + \int_0^1 \bar{\epsilon}_3 \frac{\partial \bar{\epsilon}_3}{\partial B_m} \bar{r} d\bar{r} + \int_0^{\bar{b}} \bar{\epsilon}_5 \frac{\partial \bar{\epsilon}_5}{\partial B_m} d\bar{z} &= 0 \\ \int_{\bar{b}}^{\bar{H}} \bar{\epsilon}_1 \frac{\partial \bar{\epsilon}_1}{\partial C_{\hat{m}}} d\bar{z} + \int_0^1 \bar{\epsilon}_3 \frac{\partial \bar{\epsilon}_3}{\partial C_{\hat{m}}} \bar{r} d\bar{r} + \int_0^{\bar{l}} \bar{\epsilon}_4 \frac{\partial \bar{\epsilon}_4}{\partial C_{\hat{m}}} \bar{s} \sin\alpha d\bar{s} + \int_0^{\bar{b}} \bar{\epsilon}_5 \frac{\partial \bar{\epsilon}_5}{\partial C_{\hat{m}}} d\bar{z} &= 0 \\ \int_{\bar{b}}^{\bar{H}} \bar{\epsilon}_1 \frac{\partial \bar{\epsilon}_1}{\partial C_{\bar{m}}} d\bar{z} + \int_0^1 \bar{\epsilon}_2 \frac{\partial \bar{\epsilon}_2}{\partial C_{\bar{m}}} \bar{r} d\bar{r} + \int_0^{\bar{l}} \bar{\epsilon}_4 \frac{\partial \bar{\epsilon}_4}{\partial C_{\bar{m}}} \bar{s} \sin\alpha d\bar{s} + \int_0^{\bar{b}} \bar{\epsilon}_5 \frac{\partial \bar{\epsilon}_5}{\partial C_{\bar{m}}} d\bar{z} &= 0 \end{aligned} \right\} \quad (150)$$



Eqs. (150) are a set of homogeneous algebraic equations for the constants A_n , B_n , $C_n^{\hat{}}$ and C_n^{-} with ω as a parameter. These equations can be put into the form of a matrix eigenvalue problem.

The substitution of solutions for the error functions, Eqs. (139) through (143), together with solutions for the velocity potentials, Eqs. (45) and (49), and the displacement components, Eqs. (82), (120) and (134), into Eqs. (150) yields the following matrix form:

$$\begin{bmatrix} I_{11mn} & I_{12mn} & I_{13m\hat{n}} & I_{13m\bar{n}} \\ I_{21mn} & I_{22mn} & I_{23m\hat{n}} & I_{23m\bar{n}} \\ I_{31m\hat{n}} & I_{32m\hat{n}} & I_{33m\hat{\hat{n}}} & I_{33m\hat{-n}} \\ I_{31m\bar{n}} & I_{32m\bar{n}} & I_{33m\bar{\hat{n}}} & I_{33m\bar{-n}} \end{bmatrix} \begin{Bmatrix} A_m \\ B_m \\ C_m^{\hat{}} \\ C_m^{-} \end{Bmatrix} = 0 \quad (151)$$

where I_{ijmn} ($i,j = 1,2,3$) is a function of ω .

This linear algebraic system of $4 \times (N+1)$ homogeneous simultaneous equations has a non-trivial solution only if the determinant is zero.

The matrix eigenvalue problem formulated above can be solved only with the aid of a digital computer. The method of solution consists of searching for values of ω that make the determinant of a set of simultaneous equations, Eqs. (150) vanish. An estimated value of the frequency ω is first chosen. Numerical values of the coefficients, A_n , B_n , $C_n^{\hat{}}$, and C_n^{-} are evaluated next. This process is repeated for successive successive values of ω until a zero value of the determinant is found to a desired degree of accuracy.



V. PROCEDURE FOR NUMERICAL CALCULATIONS

The eigenvalue problem for determining the natural frequencies, ω , and the corresponding mode shapes, A_n , B_n , $C_{\hat{n}}$, and $C_{\bar{n}}$ has been formulated as a linear algebraic system of $4 \times (N+1)$ homogeneous simultaneous equations. Such $4 \times (N+1)$ algebraic equations with respect to $4 \times (N+1)$ unknowns, A_n , B_n , $C_{\hat{n}}$ and $C_{\bar{n}}$ written from Eqs. (151) as

$$\left. \begin{aligned} \sum_{n=0}^N (I_{11mn} A_n + I_{12mn} B_n + I_{13m\hat{n}} C_{\hat{n}} + I_{13m\bar{n}} C_{\bar{n}}) &= 0 \\ \sum_{n=0}^N (I_{21mn} A_n + I_{22mn} B_n + I_{23m\hat{n}} C_{\hat{n}} + I_{23m\bar{n}} C_{\bar{n}}) &= 0 \\ \sum_{n=0}^N (I_{31m\hat{n}} A_n + I_{32m\hat{n}} B_n + I_{33m\hat{n}} C_{\hat{n}} + I_{33m\bar{n}} C_{\bar{n}}) &= 0 \\ \sum_{n=0}^N (I_{31m\bar{n}} A_n + I_{32m\bar{n}} B_n + I_{33m\bar{n}} C_{\hat{n}} + I_{33m\bar{n}} C_{\bar{n}}) &= 0 \end{aligned} \right\} \quad (152)$$

$$\hat{m} = 2m, \bar{m} = 2m + 1; m = 0, 1, 2, \dots, N$$

where I_{ijmn} (i and $j = 1, 2, 3$; m and $n = 0, 1, 2, \dots, N$) is a function of ω .

Obviously, $A_n = B_n = C_{\hat{n}} = C_{\bar{n}} = 0$ for $n = 0, 1, 2, \dots, N$ is a solution of Eqs. (1). There exist, however, other solutions of the determinant of Eqs. (152) is zero. Thus, from the condition that the determinant of Eqs. (152) is zero, the frequencies ω are found. Hence.



$$\begin{vmatrix} I_{11mn} & I_{12mn} & I_{13\hat{m}\hat{n}} & I_{13\bar{m}\bar{n}} \\ I_{21mn} & I_{22mn} & I_{23\hat{m}\hat{n}} & I_{23\bar{m}\bar{n}} \\ I_{31\hat{m}\hat{n}} & I_{32\hat{m}\hat{n}} & I_{33\hat{m}\hat{n}} & I_{33\bar{m}\bar{n}} \\ I_{31\bar{m}\bar{n}} & I_{32\bar{m}\bar{n}} & I_{33\hat{m}\hat{n}} & I_{33\bar{m}\bar{n}} \end{vmatrix} = 0 \quad (153)$$

Let $A_0 = k$, and solve any $4 \times (N+1) - 1$ of Eqs. (152) for the unknowns A_n ($n = 1, 2, \dots, N$), B_n , $C_{\hat{n}}$ and $C_{\bar{n}}$. For an example, solve the following $4 \times (N+1) - 1$ simultaneous equations in $4 \times (N+1) - 1$ unknowns:

$$\left. \begin{array}{l} \sum_{n=1}^N I_{11mn} \bar{A}_n + \sum_{n=0}^N (I_{12mn} \bar{B}_n + I_{13\hat{m}\hat{n}} \hat{C}_{\hat{n}} \\ \quad + I_{13\bar{m}\bar{n}} \bar{C}_{\bar{n}}) = - I_{11m0} \\ \\ \sum_{n=1}^N I_{21mn} \bar{A}_n + \sum_{n=0}^N (I_{22mn} \bar{B}_n + I_{23\hat{m}\hat{n}} \hat{C}_{\hat{n}} \\ \quad + I_{23\bar{m}\bar{n}} \bar{C}_{\bar{n}}) = - I_{21m0} \\ \\ \sum_{n=1}^N I_{31\hat{m}\hat{n}} \bar{A}_n + \sum_{n=0}^N (I_{32\hat{m}\hat{n}} \bar{B}_n + I_{33\hat{m}\hat{n}} \hat{C}_{\hat{n}} \\ \quad + I_{33\bar{m}\bar{n}} \bar{C}_{\bar{n}}) = - I_{31\hat{m}0} \\ \\ \sum_{n=1}^N I_{31\bar{m}\bar{n}} \bar{A}_n + \sum_{n=0}^N (I_{32\bar{m}\bar{n}} \bar{B}_n + I_{33\bar{m}\bar{n}} \bar{C}_{\bar{n}} \\ \quad + I_{33\hat{m}\hat{n}} \hat{C}_{\hat{n}}) = - I_{31\bar{m}0} \end{array} \right\} \begin{array}{l} (m = 0, 1, 2, \dots, N) \\ (m = 1, 2, \dots, N) \end{array} \quad (154)$$



where $\bar{A}_n = A_n/k$, $\bar{B}_n = B_n/k$, $\bar{C}_{\hat{n}} = C_{\hat{n}}/k$ and $\bar{C}_{\tilde{n}} = C_{\tilde{n}}/k$.

Eqs. (3) can be solved for \bar{A}_n , \bar{B}_n , $\bar{C}_{\hat{n}}$ and $\bar{C}_{\tilde{n}}$ by the aid of a digital computer. Known coefficients, I_{ijmn} of Eqs. (3) are calculated in Appendices A and B.



CONCLUDING REMARKS AND RECOMMENDATIONS

The solutions given by Equations (152) have been programmed and computed for the first three terms ($n = 0, 1, 2$) of the velocity potentials. This computer program has solved for the natural frequencies and the corresponding mode shapes. The results of the numerical computation are shown in Figures 2 through 4.

In the numerical calculation of the natural frequencies of the system shown in Figure 1, some difficulties have been encountered with the computer program to find zeros of the eigenvalue determinant accurately. With the present method of analysis, however, the zeros of the eigenvalue determinant are not necessarily points where the determinant changes its sign but are the minimum points where the determinant approaches zero. Thus, refining the computer program to improve the search for the minimum values of ω , three natural frequencies: $\omega = 14, 25$, and 39 rad/sec, have been found. From the solutions for the free vibration of a complete circular ring (Reference 8) and those of a circular cylinder with rigid bottom (Reference 9), the natural frequencies of the present system ($a/h_1 = 2000$) can be roughly estimated. Their numerical results show that the natural frequencies for the first three modes range from 13 to 45 rad/sec., approximately. Thus, the present method of analysis yields fairly good results.

In the numerical calculation of the mode shapes of the circular cylindrical shell, it has been found that the normal displacement does not satisfy a continuity condition at the boundary where a cylinder of liquid and an inverted cone of liquid match. This unsatisfactory result is obviously due to the fact that there is no boundary condition which can be applied to the normal displacement because with bending stiffness neglected the differential equation of motion, (Equation (12)), is of zero order in the normal displacement. To satisfy this continuity condition, a constraint, $|w_{11} - w_{12}|_{z=\bar{b}} = 0$, has been imposed on the matrix eigenvalue equation,

Equation (151), by introducing a Lagrange multiplier λ_1 . On the upper edge of the circular cylindrical shell, $z = H$, one more constraint, $w_{11}|_{z=\bar{H}} = 0$,

The corresponding Lagrange multiplier is denoted by λ_2 . With these two constraints, the matrix eigenvalue equation, Equation (151) has been replaced for the present numerical calculation by



$$\begin{array}{|c|c|c|c|c|c|c|c|} \hline & & q_{ilbn} & q_{11Hn} & & A_n \\
 I_{ijmn} & & q_{i2bn} & q_{12Hn} & & B_n \\
 & & q_{i3b\hat{n}} & q_{13H\hat{n}} & & C_{\hat{n}} \\
 & & q_{i3b\bar{n}} & q_{13H\bar{n}} & & C_{\bar{n}} \\
 \hline q_{ilbn} & q_{i2bn} & q_{i3b\hat{n}} & q_{i3b\bar{n}} & 0 & 0 & \\
 \hline q_{11Hn} & q_{12Hn} & q_{13H\hat{n}} & q_{13H\bar{n}} & 0 & 0 & \\
 \hline \end{array} = 0$$

$\frac{\lambda_1}{\lambda_2}$

where

$$\left\{ \begin{array}{l} q_{ilbn} = q_{11bn} - q_{31bn} \\ q_{i2bn} = q_{12bn} - q_{32bn} \\ q_{i3b\hat{n}} = q_{13b\hat{n}} - q_{33b\hat{n}} \\ q_{i3b\bar{n}} = q_{13b\bar{n}} - q_{33b\bar{n}} \end{array} \right.$$

The continuity condition for the normal displacement at $\bar{z}=\bar{b}$ are then well satisfied as shown in Figures 2 through 4. Computing the unknown constants A_n , B_n , C_n and $C_{\bar{n}}$ and two Lagrange multipliers λ_1 and λ_2 for the natural frequencies $\omega = 14, 25$, and 39 , the three corresponding mode shapes have been found and shown in Figures 2 through 4.

The numerical results show that the present method of analysis yields reasonable results for both the frequencies and the mode shapes. An advantage of the present method is that the solutions do not necessarily have to satisfy the boundary conditions exactly but they need to satisfy the functional errors derived at the boundaries approximately. By introducing such functional errors, the present method can thus be developed to a problem with the more complicated boundary conditions. A tank partially filled with a liquid is known to be a very difficult problem because of the complicated nature of its boundary conditions. As the next problem, therefore, it is recommended that the longitudinal oscillation of an elastic cylindrical tank with a flexible conical bulkhead, partially filled with liquid, be investigated by the method of analysis presented here.



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TABLE 1

Physical Characteristics of A Liquid-Filled Circular
Cylindrical Tank with An Inverted Conical Bulkhead

<u>Constants</u>	<u>Numerical Values</u>	<u>Units</u>
a	200	In.
b	200	In.
g	32.2×12	In/sec^2
h_i ($i = 1, 2$)	0.1	In.
ℓ	141.4×2	In.
D	1.099×10^6	$L_{bf}/\text{In.}$
E	10^7	$L_{bf}/\text{In.}^2$
H	400	In.
v	0.3	
ρ_i ($i = 1, 2$)	2.59×10^{-4}	$L_{bf}\text{Sec}^2/\text{In.}^4$
ρ_{fi} ($i = 1, 2$)	1.06×10^{-4}	$L_{bf}\text{Sec}^2/\text{In.}^4$
α	$\pi/4$	

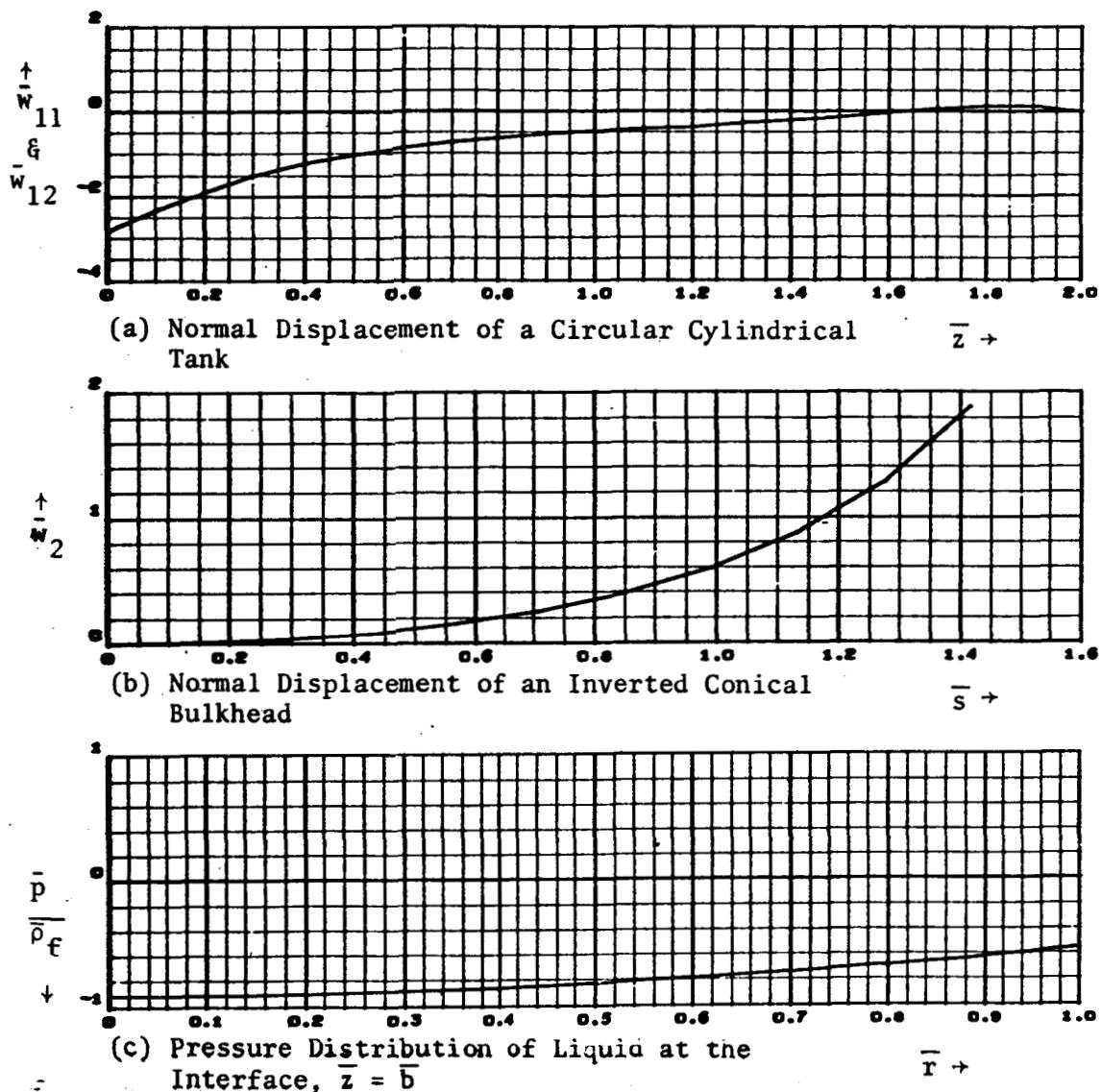


Figure 2. Axisymmetric Mode Shapes for a Liquid-Filled Circular Cylindrical Tank with an Inverted Conical Bulkhead, $f = 2.23$ cps

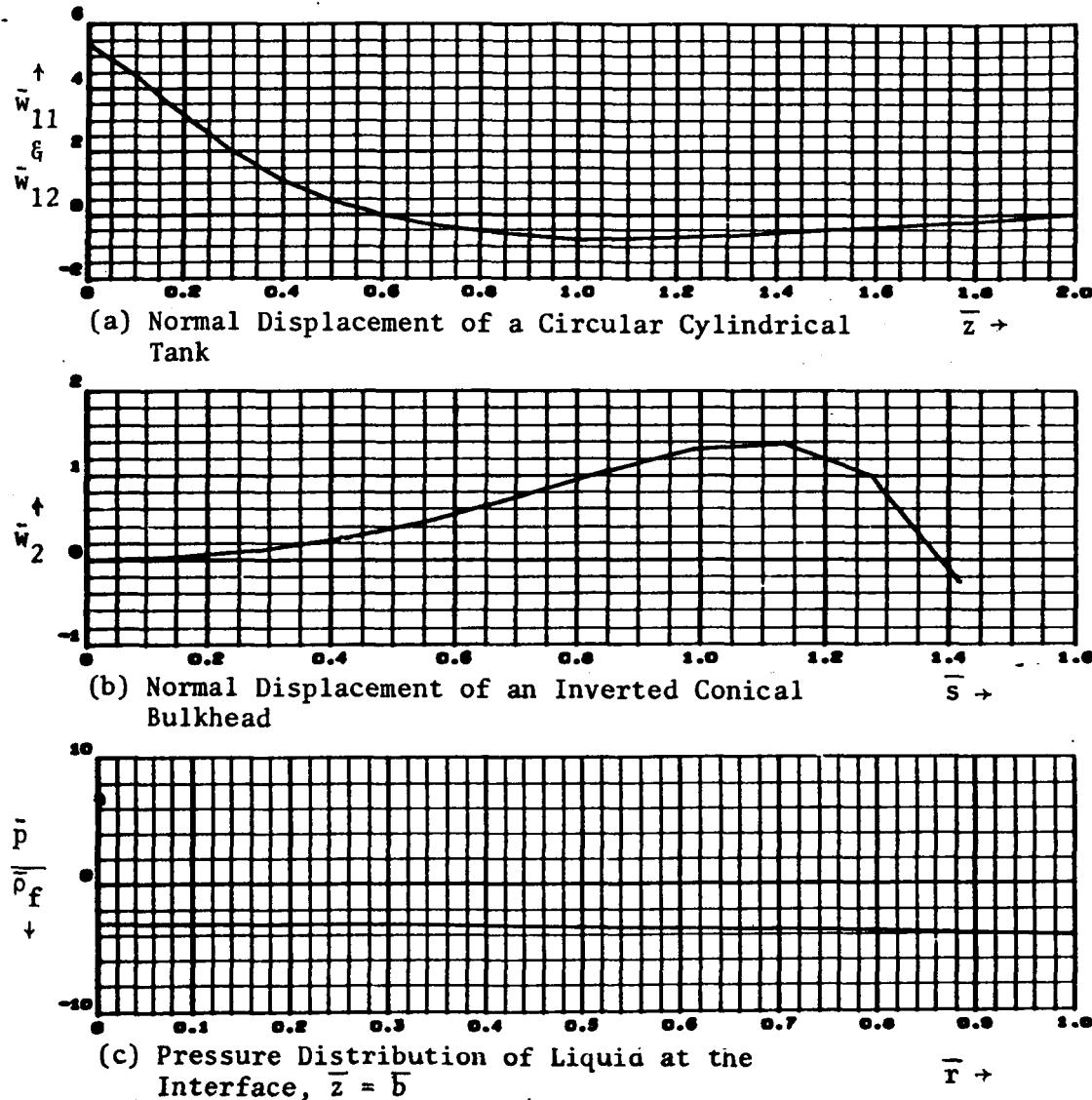


Figure 3. Axisymmetric Mode Shapes for a Liquid-Filled Circular Cylindrical Tank with an Inverted Conical Bulkhead, $f = 3.98$ cps

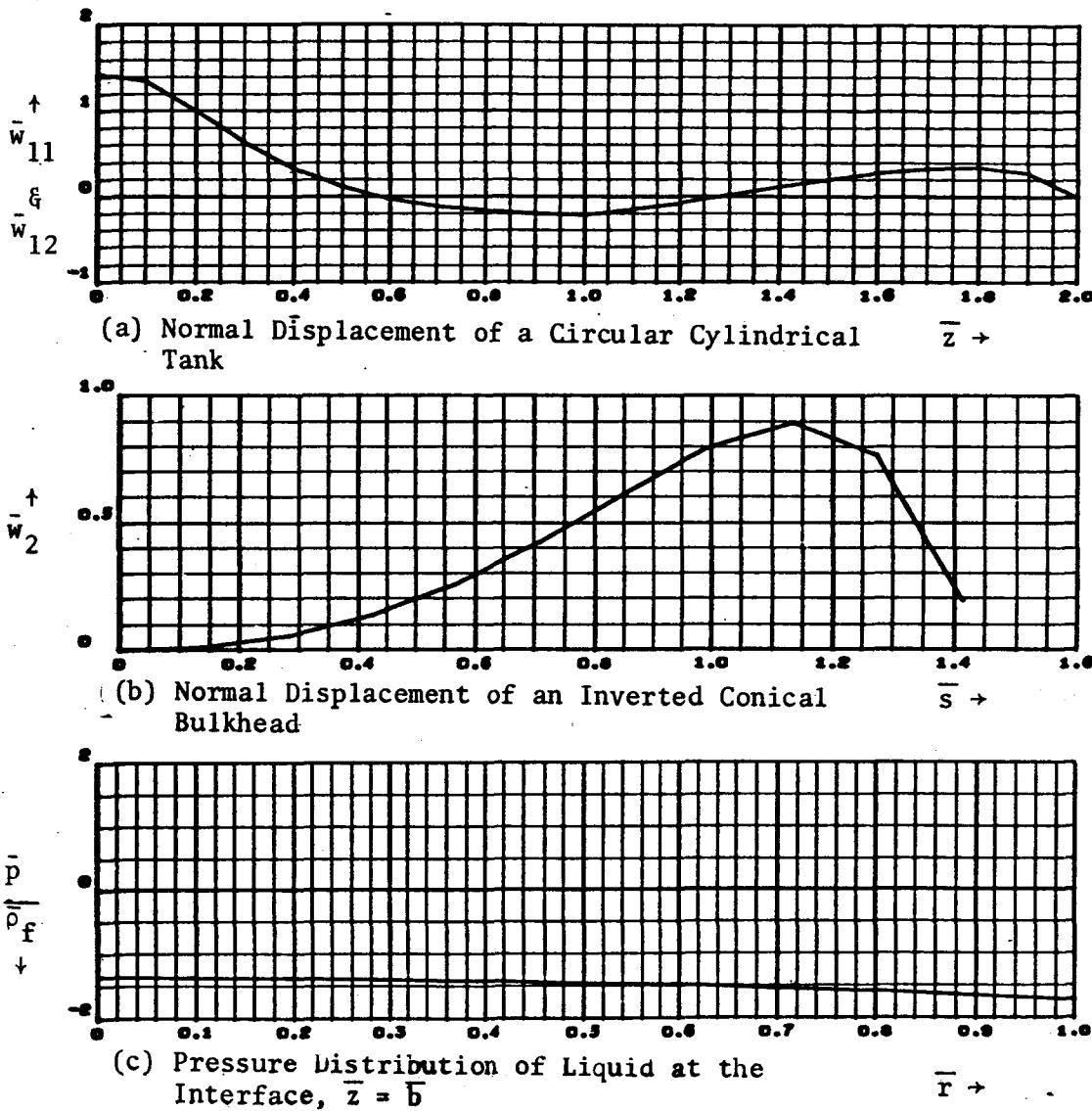


Figure 4. Axisymmetric Mode Shapes for a Liquid-Filled Circular Cylindrical Tank with an Inverted Conical Bulkhead, $f = 6.21$ cps



APPENDICES

A. Liquid Velocity Potentials $\bar{\phi}_1(\bar{r}, \bar{z})$ and $\bar{\phi}_2(\bar{s}, \mu)$
and Shell Displacement Components $\bar{w}_{11}(\bar{z})$,
 $\bar{w}_{12}(\bar{z})$ and $\bar{w}_2(\bar{s})$

The liquid velocity potentials, $\bar{\phi}_1(\bar{r}, \bar{z})$ and $\bar{\phi}_2(\bar{s}, \mu)$ and the shell displacement components, $\bar{w}_{11}(\bar{z})$, $\bar{w}_2(\bar{s})$, and $\bar{w}_{12}(\bar{z})$ have been obtained and expressed by Eqs. (44), (49), (84), (120), and (124), respectively. These are put in the following forms:

$$\left. \begin{aligned} \bar{\phi}_1(r, z) &= \sum_{n=0}^{\infty} (q_{4rzn} A_n + q_{5rzn} B_n) \\ \bar{\phi}_2(s, \mu) &= \sum_{n=0}^{\infty} (q_{6s\mu n} C_n + q_{6s\mu n}^{-} C_n^{-}) \\ \bar{w}_{11}(\bar{z}) &= \sum_{n=0}^{\infty} (q_{11zn} A_n + q_{12zn} B_n + q_{13z\hat{n}} C_n + q_{13z\hat{n}}^{-} C_n^{-}) \\ \bar{w}_2(\bar{s}) &= \sum_{n=0}^{\infty} (q_{21sn} C_n + q_{21sn}^{-} C_n^{-}) \\ \bar{w}_{12}(\bar{z}) &= \sum_{n=0}^{\infty} (q_{31zn} A_n + q_{32zn} B_n + q_{33z\hat{n}} C_n + q_{33z\hat{n}}^{-} C_n^{-}) \end{aligned} \right\} \quad (A-1)$$

where q_{4rzn} , q_{5rzn} , \dots , $q_{33z\hat{n}}$ are the function of independent variables, $\bar{r}, \bar{z}, \bar{s}$ and μ which are shown by the subscripts, r , z , s , and μ , respectively. These functions will be derived herein from Eqs. (44), (49), (84), (120), and (124) by integrating the pressure terms given in Eqs. (121) and (137).

One of the pressure terms involved in the variables, q_{21sn} and q_{21sn}^{-} is $\bar{F}_2(\mu_i)$. Thus, substituting the second equation of Eqs. (101) into the first equation of Eqs. (121) gives



$$\bar{F}_2(\mu_i) = -\frac{1-2v}{1-v} \bar{\rho}_{f2} C_0 \tan \alpha \int_0^{\bar{\ell}} \bar{s} J_0(\mu_i \bar{s}) d\bar{s}$$

$$- \sum_{n=1}^{\infty} \frac{1-2v-nv}{1-v} \bar{\rho}_{f2} C_n P_n(\mu_\alpha) \tan \alpha \int_0^{\bar{\ell}} \bar{s}^{n+1} J_0(\mu_i \bar{s}) d\bar{s} \quad (A-2)$$

The first and second integrations of Eq. (A-2) are, respectively,

$$\int_0^{\bar{\ell}} \bar{s} J_0(\mu_i \bar{s}) d\bar{s} = \frac{\bar{\ell}}{\mu_i} J_1(\mu_i \bar{\ell}) \quad \left. \begin{array}{l} \\ \end{array} \right\}$$

and

$$\int_0^{\bar{\ell}} \bar{s}^{n+1} J_0(\mu_i \bar{s}) d\bar{s} = \frac{1}{\mu_i^{n+2}} \left[n \mu_i \bar{\ell} J_0(\mu_i \bar{\ell}) + S_{n,-1}(\mu_i \bar{\ell}) + \mu_i \bar{\ell} J_1(\mu_i \bar{\ell}) S_{n+1,0}(\mu_i \bar{\ell}) + 2^{n+1} \frac{\Gamma(1 + \frac{n}{2})}{\Gamma(-\frac{n}{2})} \right] \quad \left. \begin{array}{l} \\ \end{array} \right\} \quad (A-3)$$

$$n = 1, 2, \dots, \infty$$

where $S_{n,-1}$ and $S_{n+1,0}$ are Lommel functions and expressed, respectively, as:

$$S_{n,-1}(\mu_i \bar{\ell}) = (\mu_i \bar{\ell})^{n-1} \sum_{m=0}^{\infty} \frac{(-1)^m \left(\frac{\mu_i \bar{\ell}}{2}\right)^{2m+2} \Gamma\left(1 + \frac{n}{2}\right) \Gamma\left(\frac{n}{2}\right)}{\Gamma\left(2 + \frac{n}{2} + m\right) \Gamma\left(1 + \frac{n}{2} + m\right)}$$

$$+ 2^{n-1} \Gamma\left(1 + \frac{n}{2}\right) \Gamma\left(\frac{n}{2}\right) \left\{ \sin\left[\frac{1}{2}(n+1)\pi\right] J_{-1}(\mu_i \bar{\ell}) - \cos\left[\frac{1}{2}(n+1)\pi\right] N_{-1}(\mu_i \bar{\ell}) \right\}$$

$$n = 2, 4, 6, \dots$$



$$s_{n,-1}(\mu_i \bar{\ell}) = (\mu_i \bar{\ell})^{n-1} \sum_{m=0}^{p-1} \frac{(-1)^m \Gamma\left(-\frac{1}{2}n+m\right) \Gamma\left(1 - \frac{1}{2}n+m\right)}{\left(\frac{\mu_i \bar{\ell}}{2}\right)^m \Gamma\left(-\frac{1}{2}n\right) \Gamma\left(1 - \frac{1}{2}n\right)}$$

$$+ 0 \left[(\mu_i \bar{\ell})^{n-2p} \right] \quad n = 1, 3, 5, \dots$$

$N_{-1}(\mu_i \bar{\ell})$ [Bessel functions of the second kind]

$$= -\frac{2}{\pi} \left[\left(\ln \frac{\mu_i \bar{\ell}}{2} + 0.5772157\dots \right) J_1(\mu_i \bar{\ell}) - \frac{1}{\mu_i \bar{\ell}} \right. \\ \left. - \frac{\mu_i \bar{\ell}}{4} + \frac{1 + \left(1 + \frac{1}{2}\right)}{2} \frac{(\mu_i \bar{\ell})^3}{2^3 2!} - \dots \right]$$

$$s_{n+1,0}(\mu_i \bar{\ell}) = (\mu_i \bar{\ell})^n \sum_{m=0}^{\infty} \frac{(-1)^m \left(\frac{\mu_i \bar{\ell}}{2}\right)^{2m+2} \Gamma\left(1 + \frac{n}{2}\right) \Gamma\left(1 + \frac{n}{2}\right)}{\Gamma\left(2 + \frac{n}{2} + m\right) \Gamma\left(2 + \frac{n}{2} + m\right)}$$

$$+ 2^n \Gamma\left(1 + \frac{n}{2}\right) \Gamma\left(1 + \frac{n}{2}\right) \left\{ \sin\left[\frac{1}{2}(n+1)\pi\right] J_0(\mu_i \bar{\ell}) - \cos\left[\frac{1}{2}(n+1)\pi\right] N_0(\mu_i \bar{\ell}) \right\}$$

$$n = 2, 4, 6, \dots$$

$$N_0(\mu_i \bar{\ell}) = \frac{2}{\pi} \left\{ \left(\ln \frac{\mu_i \bar{\ell}}{2} + 0.5772157\dots \right) J_0(\mu_i \bar{\ell}) \right. \\ \left. + \left[\frac{(\mu_i \bar{\ell})^2}{2^2} - \frac{(\mu_i \bar{\ell})^4}{2^4 (2!)^2} \left(1 + \frac{1}{2}\right) + \dots \right] \right\}$$



$$s_{n+1,0}(\mu_i \bar{\ell}) = (\mu_i \bar{\ell})^n \sum_{m=0}^{p-1} \frac{(-1)^m \Gamma\left(-\frac{1}{2}n+m\right) \Gamma\left(-\frac{1}{2}n+m\right)}{\left(\frac{\mu_i \bar{\ell}}{2}\right)^m \Gamma\left(-\frac{1}{2}n\right) \Gamma\left(-\frac{1}{2}n\right)} + {}_0\left[(\mu_i \bar{\ell})^{n+1-2p} \right]$$

$$n = 1, 3, 5, \dots$$

Substituting Eqs. (A-3) into Eqs. (A-2) and letting

$$\left. \begin{aligned} & -\frac{1-2v}{1-v^2} \frac{\bar{\ell}}{\mu_i} J_1(\mu_i \bar{\ell}) \tan \alpha = f_{J_{i0}}(\mu_i), \\ & -\sum_{n=1}^{\infty} \left\{ \frac{1-2v}{1-v^2} P_n(\mu_i \alpha) \frac{1}{\mu_i^{n+2}} \tan \alpha \left[n \mu_i \bar{\ell} J_0(\mu_i \bar{\ell}) \right. \right. \\ & \left. \left. + s_{n,-1}(\mu_i \bar{\ell}) - \mu_i \bar{\ell} J_{-1}(\mu_i \bar{\ell}) s_{n+1,0}(\mu_i \bar{\ell}) + 2^{n+1} \frac{\Gamma(1 + \frac{n}{2})}{\Gamma(-\frac{n}{2})} \right] \right\} \\ & = f_{J_{in}}(\mu_i), \end{aligned} \right\} \quad (A-4)$$

the pressure term $\bar{F}_2(\mu_i)$ of Eq. (A-2) may be written in the form:

$$\bar{F}_2(\mu_i) = \bar{\rho}_{f2} \sum_{n=0}^{\infty} \left[f_{J_{i\hat{n}}}(\mu_i) C_{\hat{n}} + f_{J_{i\bar{n}}}(\mu_i) C_{\bar{n}} \right] \quad (A-5)$$

The other pressure terms are defined by Eqs. (137) and they are written as follows:



$$\left. \begin{aligned} f_{p11}(\bar{z}) &= \int^{\bar{z}} \bar{p}_1(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \\ f_{p12}(\bar{z}) &= \int^{\bar{z}} \bar{p}_1(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \\ f_{p21}(\bar{z}) &= \int^{\bar{z}} \bar{p}_2(\bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \\ f_{p22}(\bar{z}) &= \int^{\bar{z}} \bar{p}_2(\bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \end{aligned} \right\} \quad (A-6)$$

Substituting the pressure, $\bar{p}_1(\bar{z})$ from Eq. (62), together with Eq. (44) into the first two equations of Eqs. (A-6) gives:

$$\begin{aligned} f_{p11}(\bar{z}) &= \bar{\rho}_{f1} \left[\left[\frac{\bar{e}_{10}}{\sqrt{g_{11}}} \sin(\sqrt{g_{11}} \bar{z}) + \frac{1}{g_{11}} \cos(\sqrt{g_{11}} \bar{z}) + \frac{\bar{z}}{\sqrt{g_{11}}} \sin(\sqrt{g_{11}} \bar{z}) \right] A_0 \right. \\ &\quad + \sum_{n=1}^{\infty} A_n \left[\frac{a_n \alpha_n}{\alpha_n^2 + g_{11}} \cosh(\alpha_n \bar{z}) \cos(\sqrt{g_{11}} z) \right. \\ &\quad \left. + \frac{a_n \sqrt{g_{11}}}{\alpha_n^2 + g_{11}} \sinh(\alpha_n \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) + \frac{\alpha_n}{\alpha_n^2 + g_{11}} \sinh(\alpha_n \bar{z}) \cos(\sqrt{g_{11}} \bar{z}) \right. \\ &\quad \left. + \frac{\sqrt{g_{11}}}{\alpha_n^2 + g_{11}} \cosh(\alpha_n \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) \right] J_0(\alpha_n) \\ &\quad + \sum_{n=1}^{\infty} \frac{B_n I_0(\beta_n)}{\cos(\beta_n b)} \left\{ \frac{\sin[(\beta_n - \sqrt{g_{11}})\bar{z} - \beta_n b]}{2(\beta_n - \sqrt{g_{11}})} \right. \\ &\quad \left. + \frac{\sin[(\beta_n + \sqrt{g_{11}})\bar{z} - \beta_n b]}{2(\beta_n + \sqrt{g_{11}})} \right\} \quad (A-7) \end{aligned}$$



$$\begin{aligned}
 f_{p12}(\bar{z}) = & \bar{\rho}_{f1} \left[\left[-\frac{\bar{e}_{10}}{\sqrt{g_{11}}} \cos(\sqrt{g_{11}} \bar{z}) + \frac{1}{g_{11}} \sin(\sqrt{g_{11}} \bar{z}) \right. \right. \\
 & \left. \left. - \frac{\bar{z}}{\sqrt{g_{11}}} \cos(\sqrt{g_{11}} \bar{z}) \right] A_0 \right. \\
 & + \sum_{n=1}^{\infty} A_n \left[\frac{a_n \alpha_n}{\alpha_n^2 + g_{11}} \cosh(\alpha_n \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) - \frac{a_n \sqrt{g_{11}}}{\alpha_n^2 + g_{11}} \sinh(\alpha_n \bar{z}) \cos(\sqrt{g_{11}} \bar{z}) \right. \\
 & \left. + \frac{\alpha_n}{\alpha_n^2 + g_{11}} \sinh(\alpha_n \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) - \frac{\sqrt{g_{11}}}{\alpha_n^2 + g_{11}} \cosh(\alpha_n \bar{z}) \cos(\sqrt{g_{11}} \bar{z}) \right] J_0(\alpha_n) \\
 & - \sum_{n=1}^{\infty} \frac{B_n I_0(\beta_n)}{\cos(\beta_n \bar{b})} \left\{ \frac{\cos[(\beta_n - \sqrt{g_{11}})\bar{z} - \beta_n \bar{b}]}{2(\beta_n - \sqrt{g_{11}})} \right. \\
 & \left. + \frac{\cos[(\beta_n + \sqrt{g_{11}})\bar{z} - \beta_n \bar{b}]}{2(\beta_n + \sqrt{g_{11}})} \right\} \quad (A-8)
 \end{aligned}$$

For $\beta_n = \sqrt{g_{11}}$, however, the last terms of Eqs. (A-7) and (A-8) become, respectively,

$$\begin{aligned}
 & + \sum_{n=1}^{\infty} \frac{B_n I_0(\beta_n)}{\cos(\beta_n \bar{b})} \left[\frac{1}{2} \bar{z} \cos(\beta_n \bar{b}) + \frac{1}{4\beta_n} \sin \beta_n (2\bar{z} - \bar{b}) \right] \text{ and} \\
 & - \sum_{n=1}^{\infty} \frac{B_n I_0(\beta_n)}{\cos(\beta_n \bar{b})} \left[\frac{1}{2} \bar{z} \sin(\beta_n \bar{b}) - \frac{1}{4\beta_n} \cos \beta_n (2\bar{z} - \bar{b}) \right].
 \end{aligned}$$



Similarly, substituting the pressure, $\bar{p}_2(\bar{z})$ from Eq. (95) together with Eq. (49) into the last two equations of Eqs. (A-6) gives:

$$f_{p21}(\bar{z}) = \bar{\rho}_{f2} \left\{ C_0 \int^{\bar{z}} \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} + \sum_{n=1}^{\infty} C_n \int^{\bar{z}} \left[1 + (\bar{b}-\bar{z})^2 \right]^{\frac{n}{2}} p_n(\mu \bar{z}) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \right\} \quad (A-9)$$

$$f_{p22}(\bar{z}) = \bar{\rho}_{f2} \left\{ C_0 \int^{\bar{z}} \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} + \sum_{n=1}^{\infty} C_n \int^{\bar{z}} \left[1 + (\bar{b}-\bar{z})^2 \right]^{\frac{n}{2}} p_n(\mu \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \right\} \quad (A-10)$$

where

$$\mu \bar{z} = \mu(z) = \frac{\bar{b} - \bar{z}}{\sqrt{1 + (b-z)^2}} \quad (A-11)$$

If the semivertex angle of cone, α is greater than $\pi/6$, the following approximation holds for the present system where the tank is fully filled with a liquid:

$$\cos(\sqrt{g_{11}} \bar{z}) \approx \cos(\sqrt{g_{11}} \bar{b}) - \frac{g_{11}}{2} \frac{\mu^2}{1-\mu^2} \cos(\sqrt{g_{11}} \bar{b})$$

$$+ \sqrt{g_{11}} \frac{\mu}{\sqrt{1-\mu^2}} \sin(\sqrt{g_{11}} \bar{b}) - \frac{g_{11}}{6} \frac{\mu^3}{(1-\mu^2)^{\frac{3}{2}}} \sin(\sqrt{g_{11}} \bar{b})$$



$$\sin(\sqrt{g_{11}} \bar{z}) \cong \sin(\sqrt{g_{11}} \bar{b}) - \frac{g_{11}}{2} \frac{\mu^2}{1-\mu^2} \sin(\sqrt{g_{11}} \bar{b})$$

$$- \sqrt{g_{11}} \frac{\mu}{\sqrt{1-\mu^2}} \cos(\sqrt{g_{11}} \bar{b}) + \frac{g_{11}}{6} \frac{\mu^3}{(1-\mu^2)^{\frac{3}{2}}} \cos(\sqrt{g_{11}} \bar{b})$$

For this approximation, the second integrals of Eqs. (A-9) and (A-10) become, respectively,

$$\left. \begin{aligned} & \int^{\bar{z}} \left[1 + (\bar{b} - \bar{z})^2 \right]^{\frac{n}{2}} P_n\left(\mu \frac{\bar{z}}{\bar{b}}\right) \cos(\sqrt{g_{11}} \bar{z}) d\bar{z} \\ &= \sum_{i=0}^5 H_{c_i} \int^{\mu} \frac{\mu^i}{(1-\mu^2)^{\frac{n+i+1}{2}}} P_n(\mu) d\mu \end{aligned} \right\} \quad (A-12)$$

and

$$\left. \begin{aligned} & \int^{\bar{z}} \left[1 + (\bar{b} - \bar{z})^2 \right]^{\frac{n}{2}} P_n\left(\mu \frac{\bar{z}}{\bar{b}}\right) \sin(\sqrt{g_{11}} \bar{z}) d\bar{z} \\ &= \sum_{i=0}^5 H_{s_i} \int^{\mu} \frac{\mu^i}{(1-\mu^2)^{\frac{n+i+1}{2}}} P_n(\mu) d\mu \end{aligned} \right\} .$$

where



$$H_{c0} = - \cos(\sqrt{g_{11}} \bar{b})$$

$$H_{c1} = - \sqrt{g_{11}} \sin(\sqrt{g_{11}} \bar{b})$$

$$H_{c2} = - \left(1 - \frac{1}{2}g_{11} \right) \cos(\sqrt{g_{11}} \bar{b})$$

$$H_{c3} = - \left(\sqrt{g_{11}} - \frac{1}{6}g_{11}^{\frac{3}{2}} \right) \sin(\sqrt{g_{11}} \bar{b})$$

$$H_{c4} = \frac{1}{2}g_{11} \cos(\sqrt{g_{11}} \bar{b})$$

$$H_{c5} = \frac{1}{6}g_{11}^{\frac{3}{2}} \sin(\sqrt{g_{11}} \bar{b})$$

$$H_{s0} = - \sin(\sqrt{g_{11}} \bar{b})$$

$$H_{s1} = \sqrt{g_{11}} \cos(\sqrt{g_{11}} \bar{b})$$

$$H_{s2} = - \left(1 - \frac{1}{2}g_{11} \right) \sin(\sqrt{g_{11}} \bar{b})$$

$$H_{s3} = \left(g_{11} - \frac{1}{6}g_{11}^{\frac{3}{2}} \right) \cos(\sqrt{g_{11}} \bar{b})$$

$$H_{s4} = \frac{1}{2}g_{11} \sin(\sqrt{g_{11}} \bar{b})$$

$$H_{s5} = - \frac{1}{6}g_{11}^{\frac{3}{2}} \cos(\sqrt{g_{11}} \bar{b})$$

The general expression for the Legendre polynomials is expressed as

$$P_n(\mu) = \sum_{k=0}^N \frac{(-1)^k (2n-2k)!}{2^n k! (n-k)! (n-2k)!} \mu^{n-2k} \quad (A-13)$$



$$\begin{aligned}
 & \int_{\bar{z}}^{\bar{z}} \left[1 + (\bar{b} - \bar{z})^2 \right]^{\frac{n}{2}} P_n\left(\mu_{\bar{z}}\right) \cos\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \\
 &= \sum_{i=0}^{5} \sum_{k=0}^{N} \frac{(-1)^k (2n-2k)! H_{ci}}{2^n k! (n-k)! (n-2k)!} \left. \int_{\bar{z}}^{\mu} \frac{\mu^{i+n-2k}}{\left(1-\mu^2\right)^{\frac{n+i+1}{2}}} d\mu \right\} \\
 & \int_{\bar{z}}^{\bar{z}} \left[1 + (\bar{b} - \bar{z})^2 \right]^{\frac{n}{2}} P_n\left(\mu_{\bar{z}}\right) \sin\left(\sqrt{g_{11}} \bar{z}\right) d\bar{z} \\
 &= \sum_{i=0}^{5} \sum_{k=0}^{N} \frac{(-1)^k (2n-2k)! H_{si}}{2^n k! (n-k)! (n-2k)!} \left. \int_{\bar{z}}^{\mu} \frac{\mu^{i+n-2k}}{\left(1-\mu^2\right)^{\frac{n+i+1}{2}}} d\mu \right\}
 \end{aligned} \tag{A-14}$$

From the integrated results: Eqs. (A-7), (A-8), (A-9), and (A-10), together with Eqs. (A-14), the pressure terms: $f_{p11}(\bar{z})$, $f_{p12}(\bar{z})$, $f_{p21}(\bar{z})$ and $f_{p22}(\bar{z})$ of Eqs. (A-6) are written, respectively, in the forms:

$$\begin{aligned}
 f_{p11}(\bar{z}) &= \bar{\rho}_{f1} \sum_{n=0}^{\infty} \left[f_{1zn}(\bar{z}) A_n + f_{2zn}(\bar{z}) B_n \right] \\
 f_{p12}(\bar{z}) &= \bar{\rho}_{f1} \sum_{n=0}^{\infty} \left[f_{3zn}(\bar{z}) A_n + f_{4zn}(\bar{z}) B_n \right] \\
 f_{p21}(\bar{z}) &= \bar{\rho}_{f2} \sum_{n=0}^{\infty} \left[f_{5z\hat{n}}(\bar{z}) C_{\hat{n}} + f_{5z\bar{n}}(\bar{z}) C_{\bar{n}} \right] \\
 f_{p22}(\bar{z}) &= \bar{\rho}_{f2} \sum_{n=0}^{\infty} \left[f_{6z\hat{n}}(\bar{z}) C_{\hat{n}} + f_{6z\bar{n}}(\bar{z}) C_{\bar{n}} \right]
 \end{aligned} \tag{A-15}$$

where



$$f_{1z0}(\bar{z}) = \frac{\bar{e}_{10}}{\sqrt{g_{11}}} \sin(\sqrt{g_{11}} \bar{z}) + \frac{1}{g_{11}} \cos(\sqrt{g_{11}} \bar{z}) + \frac{\bar{z}}{\sqrt{g_{11}}} \sin(\sqrt{g_{11}} \bar{z})$$

$$f_{1zn}(\bar{z}) = \left[\frac{\alpha_n \alpha_n}{\alpha_n^2 + g_{11}} \cosh(\alpha_n \bar{z}) \cos(\sqrt{g_{11}} \bar{z}) + \frac{\alpha_n \sqrt{g_{11}}}{\alpha_n^2 + g_{11}} \sinh(\alpha_n \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) \right. \\ \left. + \frac{\alpha_n}{\alpha_n^2 + g_{11}} \sinh(\alpha_n \bar{z}) \cos(\sqrt{g_{11}} \bar{z}) \right. \\ \left. + \frac{\sqrt{g_{11}}}{\alpha_n^2 + g_{11}} \cosh(\alpha_n \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) \right] J_0(\alpha_n) \quad (n \geq 1)$$

$$f_{2z0}(\bar{z}) = 0$$

$$f_{2zn}(\bar{z}) = \frac{I_0(\beta_n)}{\cos(\beta_n \bar{b})} \left\{ \frac{\sin[(\beta_n - \sqrt{g_{11}})\bar{z} - \beta_n \bar{b}]}{2(\beta_n - \sqrt{g_{11}})} \right. \\ \left. + \frac{\sin[(\beta_n + \sqrt{g_{11}})\bar{z} - \beta_n \bar{b}]}{2(\beta_n + \sqrt{g_{11}})} \right\} \quad (n \geq 1)$$

$$f_{3z0}(\bar{z}) = -\frac{\bar{e}_{10}}{\sqrt{g_{11}}} \cos(\sqrt{g_{11}} \bar{z}) + \frac{1}{g_{11}} \sin(\sqrt{g_{11}} \bar{z}) - \frac{\bar{z}}{\sqrt{g_{11}}} \cos(\sqrt{g_{11}} \bar{z})$$

$$f_{3zn}(\bar{z}) = \left[\frac{\alpha_n \alpha_n}{\alpha_n^2 + g_{11}} \cosh(\alpha_n \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) - \frac{\alpha_n \sqrt{g_{11}}}{\alpha_n^2 + g_{11}} \sinh(\alpha_n \bar{z}) \cos(\sqrt{g_{11}} \bar{z}) \right. \\ \left. + \frac{\alpha_n}{\alpha_n^2 + g_{11}} \sinh(\alpha_n \bar{z}) \sin(\sqrt{g_{11}} \bar{z}) - \frac{\sqrt{g_{11}}}{\alpha_n^2 + g_{11}} \cosh(\alpha_n \bar{z}) \cos(\sqrt{g_{11}} \bar{z}) \right] J_0(\alpha_n) \\ (n \geq 1)$$

$$f_{4z0}(\bar{z}) = 0$$

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$$f_{4zn}(\bar{z}) = - \frac{I_0(\beta_n)}{\cos(\beta_n \bar{b})} \left\{ \frac{\cos[(\beta_n - \sqrt{g_{11}})\bar{z} - \beta_n \bar{b}]}{2(\beta_n - \sqrt{g_{11}})} + \frac{\cos[(\beta_n + \sqrt{g_{11}})\bar{z} - \beta_n \bar{b}]}{2(\beta_n + \sqrt{g_{11}})} , (n \geq 1) \right\}$$

$$f_{5z\hat{n}}(\bar{z}) = \sum_{i=0}^5 \sum_{k=0}^N \frac{(-1)^k (2\hat{n}-2k)! H_{ci}}{2^{\hat{n}} k! (\hat{n}-k)! (\hat{n}-2k)!} \int_{-\infty}^{\mu} \frac{\mu^{i+\hat{n}-2k}}{(1-\mu^2)^{\frac{\hat{n}+i+1}{2}}} d\mu$$

$$f_{5z\bar{n}}(\bar{z}) = \sum_{i=0}^5 \sum_{k=0}^N \frac{(-1)^k (2\bar{n}-2k)! H_{ci}}{2^{\bar{n}} k! (\bar{n}-k)! (\bar{n}-2k)!} \int_{-\infty}^{\mu} \frac{\mu^{i+\bar{n}-2k}}{(1-\mu^2)^{\frac{\bar{n}+i+1}{2}}} d\mu$$

$$f_{6z\hat{n}}(\bar{z}) = \sum_{i=0}^5 \sum_{k=0}^N \frac{(-1)^k (2\hat{n}-2k)! H_{si}}{2^{\hat{n}} k' (\hat{n}-k)! (\hat{n}-2k)!} \int_{-\infty}^{\mu} \frac{\mu^{i+\hat{n}-2k}}{(1-\mu^2)^{\frac{\hat{n}+i+1}{2}}} d\mu$$

$$f_{6z\bar{n}}(\bar{z}) = \sum_{i=0}^5 \sum_{k=0}^N \frac{(-1)^k (2\bar{n}-2k)! H_{si}}{2^{\bar{n}} k! (\bar{n}-k)! (\bar{n}-2k)!} \int_{-\infty}^{\mu} \frac{\mu^{i+\bar{n}-2k}}{(1-\mu^2)^{\frac{\bar{n}+i+1}{2}}} d\mu \quad (A-16)$$

The substitution of the expressions for $F_2(\mu_i)$, $f_{p11}(\bar{z})$, $f_{p12}(\bar{z})$, $f_{p21}(\bar{z})$, and $f_{p22}(\bar{z})$ from Eqs. (A-5) and (A-15) into Eqs. (44), (49), (84), (120) and (124) yields, after the lengthy mathematical manipulation, the expressions for q_{4rzn} , q_{5rzn} , ---, $q_{33z\bar{n}}$, which are the functional coefficients of A_n , B_n , $C_{\hat{n}}$ and $C_{\bar{n}}$, as follows:



$$q_{4rz0} = \bar{e}_{10} + \bar{z}$$

$$q_{4rzn} = \left[a_n \sinh(\alpha_n \bar{z}) + \cosh(\alpha_n \bar{z}) \right] J_0(\alpha_n \bar{r}), \quad (n \geq 1) \quad (A-17)$$

$$q_{5rz0} = 0$$

$$q_{5rzn} = \frac{\cos \left[\beta_n (\bar{z} - \bar{b}) \right]}{\cos(\beta_n \bar{b})} I_0(\beta_n \bar{r}), \quad (n \geq 1) \quad (A-18)$$

$$q_{6s\mu\hat{n}} = \bar{s}^{\hat{n}} P_{\hat{n}}(\mu) \quad (A-19)$$

$$q_{6s\mu\bar{n}} = \bar{s}^{\bar{n}} P_{\bar{n}}(\mu) \quad (A-20)$$

$$\begin{aligned} q_{11zn} = & \bar{\rho}_{f1} \left[\frac{v \ell_{k0}}{1 - \kappa_1 \omega^2} \left(\ell_{k11} f_{h1n} + \ell_{k12} f_{h3n} + \ell_{k12} \ell_{11} f_{h1n} \right. \right. \\ & \left. \left. + \ell_{k12} f_{h3n} \right) + \frac{v g_{12}}{1 - \kappa_1 \omega^2} \left(\ell_{13} f_{1zn} + \ell_{14} f_{3zn} \right) \right. \end{aligned}$$

$$\left. + \frac{1 + v g_{12}}{1 - \kappa_1 \omega^2} f_{p1zn} \right] \quad (A-21)$$

where



$$\ell_{k0} = -\ell_{c2}\ell_{13} + \ell_{c1}\ell_{14} - \ell_{c5}\ell_{14}$$

$$\ell_{k11} = \frac{-\ell_{c4}\ell_{13} + \ell_{c3}\ell_{14}}{\ell_{k0}}, \quad \ell_{k12} = \frac{\ell_{c2}\ell_{13} - \ell_{c1}\ell_{14}}{\ell_{k0}}$$

$$\ell_{k21} = \frac{\ell_{13} - \ell_{c6}\ell_{14}}{\ell_{k0}}, \quad \ell_{k22} = \frac{\ell_{c5}\ell_{14}}{\ell_{k0}}$$

$$\ell_{c1} = \frac{\ell_{11}}{\ell_{11}\ell_{14} - \ell_{12}\ell_{13}}, \quad \ell_{c2} = \frac{\ell_{12}}{\ell_{11}\ell_{14} - \ell_{12}\ell_{13}}, \quad \ell_{c5} = \frac{1}{\ell_{24}}$$

$$\ell_{c3} = \frac{\ell_{13}}{\ell_{11}\ell_{14} - \ell_{12}\ell_{13}}, \quad \ell_{c4} = \frac{\ell_{14}}{\ell_{11}\ell_{14} - \ell_{12}\ell_{13}}, \quad \ell_{c6} = \frac{\ell_{23}}{\ell_{24}}$$

$$f_{h1n} = - \left(1 - \frac{v^2}{1-\kappa_1\omega^2} \right) g_{12} \sqrt{g_{11}} \left[f_{1Hn} \sin(\sqrt{g_{11}} H) - f_{3Hn} \cos(\sqrt{g_{11}} H) \right]$$

$$+ \left\{ \left(1 - \frac{v^2}{1-\kappa_1\omega^2} \right) g_{12} - \frac{v}{1-\kappa_1\omega^2} \right\} f_{p1Hn}$$

$$f_{h2n} = - \left(1 - \frac{v^2}{1-\kappa_1\omega^2} \right) g_{12} \sqrt{g_{11}} \left[f_{2Hn} \sin(\sqrt{g_{11}} H) - f_{4Hn} \cos(\sqrt{g_{11}} H) \right]$$

$$+ \left\{ \left(1 - \frac{v^2}{1-\kappa_1\omega^2} \right) g_{12} - \frac{v}{1-\kappa_1\omega^2} \right\} f_{p2Hn}$$

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$$f_{h3n} = -g_{12}\sqrt{g_{11}} \left[f_{1bn} \sin(\sqrt{g_{11}} \bar{b}) - f_{3bn} \cos(\sqrt{g_{11}} \bar{b}) \right] + g_{12} f_{p1bn}$$

$$f_{h4n} = -g_{12}\sqrt{g_{11}} \left[f_{2bn} \sin(\sqrt{g_{11}} \bar{b}) - f_{4bn} \cos(\sqrt{g_{11}} \bar{b}) \right] + g_{12} f_{p2bn}$$

$$f_{h5n} = g_{12} f_{50n}$$

$$f_{h5\bar{n}} = g_{12} f_{50\bar{n}}$$

$$f_{h6n} = -g_{12}\sqrt{g_{11}} \left[f_{5bn} \sin(\sqrt{g_{11}} \bar{b}) - f_{6bn} \cos(\sqrt{g_{11}} \bar{b}) \right] + g_{12} f_{p3bn}$$

$$f_{h6\bar{n}} = -g_{12}\sqrt{g_{11}} \left[f_{5b\bar{n}} \sin(\sqrt{g_{11}} \bar{b}) - f_{6b\bar{n}} \cos(\sqrt{g_{11}} \bar{b}) \right] + g_{12} f_{p3b\bar{n}}$$

$$f_{p1z0} = \bar{e}_{10} + \bar{z}$$

$$f_{p1zn} = \left[a_n \sinh(\alpha_n \bar{z}) + \cosh(\alpha_n \bar{z}) \right] J_0(\alpha_n) \quad (n \geq 1)$$

$$f_{p2z0} = 0$$

$$f_{p2zn} = \frac{\cos[\beta_n(\bar{z}-\bar{b})]}{\cos(\beta_n \bar{b})} I_0(\beta_n) \quad (n \geq 1)$$

$$f_{p3z\hat{n}} = \left[1 + (\bar{b} - \bar{z})^2 \right]^{\frac{n}{2}} P_{\hat{n}}(\mu_{\bar{z}})$$

$$f_{p3z\bar{n}} = \left[1 + (\bar{b} - \bar{z})^2 \right]^{\frac{n}{2}} P_{\bar{n}}(\mu_{\bar{z}})$$



$$f_{J_i \hat{n}} = f_{J_i n} \quad \text{for } n = 0, 2, 4, \dots$$

$$f_{J_i \bar{n}} = f_{J_i n} \quad \text{for } n = 1, 3, 5, \dots$$

$$\mu_\alpha = \cos\alpha$$

$$\ell_{z13} = -\sqrt{g_{11}} \sin(\sqrt{g_{11}} \bar{z})$$

$$\ell_{z14} = \sqrt{g_{11}} \cos(\sqrt{g_{11}} \bar{z})$$

$$q_{12zn} = \bar{\rho}_{f1} \left[\frac{v \ell_{k0}}{1-\kappa_1 \omega^2} (\ell_{k11} f_{h2n} + \ell_{k12} f_{h4n} + \ell_{k12} \ell_{k11} f_{h2n} + \ell_{k12}^2 f_{h4n}) \right. \\ \left. + \frac{v g_{12}}{1-\kappa_1 \omega^2} (\ell_{13} f_{2zn} + \ell_{14} f_{4zn}) + \frac{1 + v g_{12}}{1-\kappa_1 \omega^2} f_{p2zn} \right] \quad (A-22)$$

$$q_{13z\hat{n}} = \bar{\rho}_{f1} \frac{v \ell_{k0}}{1-\kappa_1 \omega^2} \bar{\rho}_{fc} \ell_{k12} (\ell_{k21} f_{h5\hat{n}} + \ell_{k22} f_{h6\hat{n}}) \quad (A-23)$$

$$q_{13\bar{n}} = \bar{\rho}_{f1} \frac{v \ell_{k0}}{1-\kappa_1 \omega^2} \bar{\rho}_{fc} \ell_{k12} (\ell_{k21} f_{h5\bar{n}} + \ell_{k22} f_{h6\bar{n}}) \quad (A-24)$$

$$q_{21s\hat{n}} = \sum_i \left\{ \frac{2 \tan\alpha}{\ell^{-2}} \bar{\rho}_{f2} \frac{\mu_i^2}{(\sigma^2 + \mu_i^2)(\lambda^2 - \mu_i^2)} \frac{v \bar{s} \mu_i J_1(\mu_i \bar{s}) - J_0(\mu_i \bar{s})}{[J_0(\mu_i \bar{s})]^2} \right\} f_{J_i \hat{n}} \\ - \bar{\rho}_{f2} \bar{s}^{\hat{n}+2} P_{\hat{n}}(\mu_\alpha) \tan^2 \alpha \quad (A-25)$$



$$q_{21zn} = \sum_i \left\{ \frac{\frac{2\tan\alpha}{\ell^2}}{\bar{\rho}_{f2}} \frac{\mu_i^2}{(\sigma^2 + \mu_i^2)(\lambda^2 - \mu_i^2)} \frac{\frac{v\bar{s}\mu_i J_1(\mu_i \bar{s}) - J_0(\mu_i \bar{s})}{[J_0(\mu_i \ell)]^2}}{f_{Jin}} \right\} - \bar{\rho}_{f2} \bar{s}^{n+2} P_n(\mu_\alpha) \tan^2 \alpha \quad (A-26)$$

$$q_{31zn} = - \frac{v}{1-\kappa_1 \omega^2} \bar{\rho}_{f1} \ell_{c5} \ell_{z14} (\ell_{k11} f_{h1n} + \ell_{k12} f_{h3n}) \quad (A-27)$$

$$q_{32zn} = - \frac{v}{1-\kappa_1 \omega^2} \bar{\rho}_{f1} \ell_{c5} \ell_{z14} (\ell_{k11} f_{h2n} + \ell_{k12} f_{h4n}) \quad (A-28)$$

$$\begin{aligned} q_{33zn} = & - \frac{v}{1-\kappa_1 \omega^2} \bar{\rho}_{f2} \left[(\ell_{z13} - \ell_{c6} \ell_{z14}) f_{h5n} + \ell_{c5} \ell_{z14} f_{h6n} + \right. \\ & \left. + \ell_{c5} \ell_{z14} (\ell_{k21} f_{h5n} + \ell_{k22} f_{h6n}) \right] + \frac{v g_{12}}{1-\kappa_1 \omega^2} \bar{\rho}_{f2} (\ell_{z13} f_{5zn} \\ & + \ell_{z14} f_{6zn}) + \frac{1 + v g_{12}}{1-\kappa_1 \omega^2} \bar{\rho}_{f2} f_{p3zn} \end{aligned} \quad (A-29)$$

$$\begin{aligned} q_{33zn} = & - \frac{v}{1-\kappa_1 \omega^2} \bar{\rho}_{f2} \left[(\ell_{z13} - \ell_{c6} \ell_{z14}) f_{h5n} + \ell_{c5} \ell_{z14} f_{h6n} \right. \\ & \left. + \ell_{c5} \ell_{z14} (\ell_{k21} f_{h5n} + \ell_{k22} f_{h6n}) \right] + \frac{v g_{12}}{1-\kappa_1 \omega^2} \bar{\rho}_{f2} (\ell_{z13} f_{5zn} \end{aligned} \quad (A-30)$$

$$+ \ell_{z14} f_{6zn}) + \frac{1 + v g_{12}}{1-\kappa_1 \omega^2} \bar{\rho}_{f2} f_{p3zn} \quad (A-30)$$

B. Functional Errors $\bar{\epsilon}_i$ ($i = 1, 2, \dots, 5$)

The functional errors, $\bar{\epsilon}_i$, ($i = 1, 2, \dots, 5$) defined by Eqs. (139) through (143) are derived in a series form in this appendix. The element, I_{ijmn} , of the matrix eigenvalue problem is then calculated in terms of the functional coefficient which is derived in a series form of each functional error, $\bar{\epsilon}_i$.

From the expression of $\bar{\phi}_1(r, z)$ and $\bar{\phi}_2(s,)$ given in Eqs. (A-1), the derivatives of these two variables are obtained as

$$\left. \begin{aligned} \frac{\partial \bar{\phi}_1}{\partial \bar{r}} &= \sum_{n=0}^{\infty} (\gamma_{1rzn} A_n + \gamma_{2rzn} B_n) \\ \frac{\partial \bar{\phi}_1}{\partial \bar{z}} &= \sum_{n=0}^{\infty} (\gamma_{3rzn} A_n + \gamma_{4rzn} B_n) \\ \frac{\partial \bar{\phi}_2}{\partial \bar{s}} &= \sum_{n=0}^{\infty} (\gamma_{5s\mu n} C_n + \gamma_{5s\mu n}^{-} C_n^{-}) \\ \frac{1}{\bar{s}} \frac{\partial \bar{\phi}}{\partial \theta} &= \sum_{n=0}^{\infty} (\gamma_{6s\mu n} C_n + \gamma_{6s\mu n}^{-} C_n^{-}) \end{aligned} \right\} \quad (B-1)$$

where

$$\gamma_{1rzn} = -\alpha_n [a_n \sinh(\alpha_n \bar{z}) + \cosh(\alpha_n \bar{z})] J_1(\alpha_n \bar{r}) \quad (n \geq 1)$$

$$\gamma_{2rzn} = \beta_n \frac{\cos[\beta_n(z-b)]}{\cos(\beta_n b)} I_1(\beta_n \bar{r}) \quad (n \geq 1)$$



$$\gamma_{3rz_n} = \alpha_n \left[a_n \cosh(\alpha_n \bar{z}) + \sinh(\alpha_n \bar{z}) \right] J_0(\alpha_n \bar{r}) \quad (n \geq 1)$$

$$\gamma_{4rz_n} = -\beta_n \frac{\sin[\beta_n(\bar{z}-\bar{b})]}{\cos(\beta_n \bar{b})} I_0(\beta_n \bar{r}) \quad (n \geq 1)$$

$$\gamma_{5s\mu\hat{n}} = \bar{s}^{\hat{n}-1} P_{\hat{n}}(\mu)$$

$$\gamma_{5s\mu\bar{n}} = \bar{s}^{\bar{n}-1} P_{\bar{n}}(\mu)$$

$$\gamma_{6s\mu\hat{n}} = -\frac{1}{\sqrt{1-\mu^2}} \bar{s}^{\hat{n}-1} [P_{\hat{n}-1}(\mu) - \mu P_{\hat{n}}(\mu)]$$

$$\gamma_{6s\mu\bar{n}} = -\frac{1}{\sqrt{1-\mu^2}} \bar{s}^{\bar{n}-1} [P_{\bar{n}-1}(\mu) - \mu P_{\bar{n}}(\mu)]$$

$$\gamma_{1rz0} \stackrel{\Delta}{=} 0, \quad \gamma_{2rz0} \stackrel{\Delta}{=} 0$$

$$\gamma_{3rz0} \stackrel{\Delta}{=} 1, \quad \gamma_{4rz0} \stackrel{\Delta}{=} 0$$

Substituting Eqs. (A-1) and (B-1) into the functional errors, Eqs. (139) through (143) gives

$$\left. \begin{aligned}
 \bar{\varepsilon}_1 &= \sum_{n=0}^{\infty} (e_{11zn} A_n + e_{12zn} B_n + e_{13z\hat{n}} C_{\hat{n}} + e_{13z\bar{n}} C_{\bar{n}}) \\
 \bar{\varepsilon}_2 &= \sum_{n=0}^{\infty} (e_{21rn} A_n + e_{224n} B_n + e_{23r\bar{n}} C_{\bar{n}}) \\
 \bar{\varepsilon}_3 &= \sum_{n=0}^{\infty} (e_{31rn} A_n + e_{32rn} B_n + e_{33r\hat{n}} C_{\hat{n}}) \\
 \bar{\varepsilon}_4 &= \sum_{n=0}^{\infty} (e_{43s\hat{n}} C_{\hat{n}} + e_{43s\bar{n}} C_{\bar{n}}) \\
 \bar{\varepsilon}_5 &= \sum_{n=0}^{\infty} (e_{51zn} A_n + e_{52zn} B_n + e_{53z\hat{n}} C_{\hat{n}} + e_{53z\bar{n}} C_{\bar{n}})
 \end{aligned} \right\} \quad (B-2)$$



where

$$e_{11zn} = \gamma_{11zn} - \beta_{K1} q_{11zn}$$

$$e_{12zn} = \gamma_{21zn} - \beta_{K1} q_{12zn}$$

$$e_{13z\hat{n}} = -\beta_{K1} q_{13z\hat{n}}$$

$$e_{13z\bar{n}} = -\beta_{K1} q_{13z\bar{n}}$$

$$e_{21rn} = \gamma_{3rbn}$$

$$e_{22rn} = \gamma_{4rbn}$$

$$e_{23s\bar{n}} = e_{23r\bar{n}} = -\gamma_{6s0\bar{n}}$$

$$e_{31rn} = q_{4rbn}$$

$$e_{32rn} = q_{5rbn}$$

$$e_{33s\hat{n}} = e_{33r\hat{n}} = -q_{6s0\hat{n}}$$

$$e_{43s\hat{n}} = \gamma_{6s0\hat{n}} - \beta_{K2} q_{21s\hat{n}}$$

$$e_{43s\bar{n}} = \gamma_{6s0\bar{n}} - \beta_{K2} q_{21s\bar{n}}$$

$$e_{51zn} = -\beta_{K1} q_{31zn}$$

$$e_{52zn} = -\beta_{K1} q_{32zn}$$



$$e_{53z\hat{n}} = K_{sz} \gamma_{5z\hat{n}} + K_{\theta z} \gamma_{6z\hat{n}} - \beta_{K1} q_{33z\hat{n}}$$

$$e_{53z\bar{n}} = K_{sz} \gamma_{5z\bar{n}} + K_{\theta z} \gamma_{6z\bar{n}} - \beta_{K1} q_{33z\bar{n}}$$

The conditions for the minimum of the total integrated squared error are obtained by substituting Eqs. (B-2) into Eqs. (147) and then Eqs. (148). It follows four simultaneous equations:

$$\int_{\bar{b}}^{\bar{H}} \bar{\epsilon}_1 \frac{\partial \bar{\epsilon}_1}{\partial A_m} d\bar{z} + \int_0^1 \bar{\epsilon}_2 \frac{\partial \bar{\epsilon}_2}{\partial A_m} \bar{r} d\bar{r} + \int_0^1 \bar{\epsilon}_3 \frac{\partial \bar{\epsilon}_3}{\partial A_m} \bar{r} d\bar{r} + \int_0^{\bar{b}} \bar{\epsilon}_5 \frac{\partial \bar{\epsilon}_5}{\partial A_m} d\bar{z} = 0$$

$$\int_{\bar{b}}^{\bar{H}} \bar{\epsilon}_1 \frac{\partial \bar{\epsilon}_1}{\partial B_m} d\bar{z} + \int_0^1 \bar{\epsilon}_2 \frac{\partial \bar{\epsilon}_2}{\partial B_m} \bar{r} d\bar{r} + \int_0^1 \bar{\epsilon}_3 \frac{\partial \bar{\epsilon}_3}{\partial B_m} \bar{r} d\bar{r} + \int_0^{\bar{b}} \bar{\epsilon}_5 \frac{\partial \bar{\epsilon}_5}{\partial B_m} d\bar{z} = 0$$

$$\int_{\bar{b}}^{\bar{H}} \bar{\epsilon}_1 \frac{\partial \bar{\epsilon}_1}{\partial C_{\hat{m}}} d\bar{z} + \int_0^1 \bar{\epsilon}_3 \frac{\partial \bar{\epsilon}_3}{\partial C_{\hat{m}}} \bar{r} d\bar{r} + \int_0^1 \bar{\epsilon}_4 \frac{\partial \bar{\epsilon}_4}{\partial C_{\hat{m}}} \bar{s} \sin \alpha d\bar{s} + \int_0^{\bar{b}} \bar{\epsilon}_5 \frac{\partial \bar{\epsilon}_5}{\partial C_{\hat{m}}} d\bar{z} = 0$$

$$\int_{\bar{b}}^{\bar{H}} \bar{\epsilon}_1 \frac{\partial \bar{\epsilon}_1}{\partial C_{\bar{m}}} d\bar{z} + \int_0^1 \bar{\epsilon}_2 \frac{\partial \bar{\epsilon}_2}{\partial C_{\bar{m}}} \bar{r} d\bar{r} + \int_0^{\bar{l}} \bar{\epsilon}_4 \frac{\partial \bar{\epsilon}_4}{\partial C_{\bar{m}}} \bar{s} \sin \alpha d\bar{s} + \int_0^{\bar{b}} \bar{\epsilon}_5 \frac{\partial \bar{\epsilon}_5}{\partial C_{\bar{m}}} d\bar{z} = 0$$

(B-3)

The substitution of Eqs. (B-2) into Eqs. (B-3) yields a desired linear algebraic system of $4 \times (N+1)$ homogeneous simultaneous equations for $4 \times (N+1)$ unknowns A_n , B_n , $C_{\hat{n}}$ and $C_{\bar{n}}$ ($n = 0, 1, 2, \dots, N$) as follows:



$$\left. \begin{aligned}
 \sum_{n=0}^{\infty} (I_{11mn} A_n + I_{12mn} B_n + I_{13m\hat{n}} C_{\hat{n}} + I_{13m\bar{n}} C_{\bar{n}}) &= 0 \\
 \sum_{n=0}^{\infty} (I_{21mn} A_n + I_{22mn} B_n + I_{23m\hat{n}} C_{\hat{n}} + I_{23m\bar{n}} C_{\bar{n}}) &= 0 \\
 \sum_{n=0}^{\infty} (I_{31m\hat{n}} A_n + I_{32m\hat{n}} B_n + I_{33m\hat{n}} C_{\hat{n}} + I_{33m\bar{n}} C_{\bar{n}}) &= 0 \\
 \sum_{n=0}^{\infty} (I_{31m\bar{n}} A_n + I_{32m\bar{n}} B_n + I_{33m\hat{n}} C_{\hat{n}} + I_{33m\bar{n}} C_{\bar{n}}) &= 0
 \end{aligned} \right\} \quad (B-4)$$

where

$$\begin{aligned}
 I_{11mn} &= \int_{\bar{b}}^{\bar{H}} e_{11zm} e_{11zn} d\bar{z} + \int_0^1 e_{21rm} e_{21rn} \bar{r} d\bar{r} \\
 &\quad + \int_0^1 e_{31rm} e_{31rn} \bar{r} d\bar{r} + \int_0^{\bar{b}} e_{51zn} e_{51zn} d\bar{z} \\
 I_{12mn} &= \int_{\bar{b}}^{\bar{H}} e_{11zm} e_{12zn} d\bar{z} + \int_0^1 e_{21rm} e_{22rn} \bar{r} d\bar{r} \\
 &\quad + \int_0^1 e_{31rm} e_{32rn} \bar{r} d\bar{r} + \int_0^{\bar{b}} e_{51zm} e_{52zn} d\bar{z} \\
 I_{13m\hat{n}} &= \int_{\bar{b}}^{\bar{H}} e_{11zm} e_{13z\hat{n}} d\bar{z} + \int_0^1 e_{31rm} e_{33r\hat{n}} \bar{r} d\bar{r} \\
 &\quad + \int_0^{\bar{b}} e_{51zm} e_{53z\hat{n}} d\bar{z}
 \end{aligned}$$



$$I_{13mn} = \int_{\bar{b}}^{\bar{H}} e_{11zm} e_{13zn} d\bar{z} + \int_0^1 e_{21rm} e_{23rn} \bar{r} d\bar{r}$$

$$+ \int_0^{\bar{b}} e_{51zm} e_{53zn} d\bar{z}$$

$$I_{21mn} = \int_{\bar{b}}^{\bar{H}} e_{12zm} e_{11zn} d\bar{z} + \int_0^1 e_{22rm} e_{21rn} \bar{r} d\bar{r}$$

$$+ \int_0^1 e_{32rm} e_{31rn} \bar{r} d\bar{r} + \int_0^{\bar{b}} e_{52zm} e_{51zn} d\bar{z}$$

$$I_{22mn} = \int_{\bar{b}}^{\bar{H}} e_{12zm} e_{12zn} d\bar{z} + \int_0^1 e_{22rm} e_{22rn} \bar{r} d\bar{r}$$

$$+ \int_0^1 e_{32rm} e_{32rn} \bar{r} d\bar{r} + \int_0^{\bar{b}} e_{52zm} e_{52zn} d\bar{z}$$

$$I_{23mn} = \int_{\bar{b}}^{\bar{H}} e_{12zm} e_{13zn} d\bar{z} + \int_0^1 e_{32rm} e_{33rn} \bar{r} d\bar{r}$$

$$+ \int_0^{\bar{b}} e_{52zm} e_{53zn} d\bar{z}$$

$$I_{23mn} = \int_{\bar{b}}^{\bar{H}} e_{12zm} e_{13zn} d\bar{z} + \int_0^1 e_{22rm} e_{23rn} \bar{r} d\bar{r}$$

$$+ \int_0^{\bar{b}} e_{52zm} e_{53zn} d\bar{z}$$

$$I_{31mn} = \int_{\bar{b}}^{\bar{H}} e_{13zm} e_{11zn} d\bar{z} + \int_0^1 e_{33rm} e_{31rn} \bar{r} d\bar{r}$$

$$+ \int_0^{\bar{b}} e_{53zm} e_{51zn} d\bar{z}$$



$$I_{32\hat{m}\hat{n}} = \int_{\bar{b}}^{\bar{H}} e_{13z\hat{m}} e_{12zn} d\bar{z} + \int_0^1 e_{33r\hat{m}} e_{32rn} \bar{r} d\bar{r} \\ + \int_0^{\bar{b}} e_{53z\hat{m}} e_{52zn} d\bar{z}$$

$$I_{33\hat{m}\hat{n}} = \int_{\bar{b}}^{\bar{H}} e_{13z\hat{m}} e_{13z\hat{n}} d\bar{z} + \int_0^1 e_{33r\hat{m}} e_{33r\hat{n}} \bar{r} d\bar{r} \\ + \int_0^{\bar{\ell}} e_{43s\hat{m}} e_{43s\hat{n}} \bar{s} \sin\alpha d\bar{s} + \int_0^{\bar{b}} e_{53z\hat{m}} e_{53z\hat{n}} d\bar{z}$$

$$I_{33\hat{m}\hat{n}} = \int_{\bar{b}}^{\bar{H}} e_{13z\hat{m}} e_{13z\hat{n}} d\bar{z} + \int_0^{\bar{\ell}} e_{43s\hat{m}} e_{43s\hat{n}} \bar{s} \sin\alpha d\bar{s} \\ + \int_0^{\bar{b}} e_{53z\hat{m}} e_{53z\hat{n}} d\bar{z}$$

$$I_{31\bar{m}\bar{n}} = \int_{\bar{b}}^{\bar{H}} e_{13z\bar{m}} e_{11zn} d\bar{z} + \int_0^1 e_{23r\bar{m}} e_{21rn} \bar{r} d\bar{r} \\ + \int_0^{\bar{b}} e_{53z\bar{m}} e_{51zn} d\bar{z}$$



$$I_{32\bar{m}\bar{n}} = \int_{\bar{b}}^{\bar{H}} e_{13z\bar{m}} e_{12zn} d\bar{z} + \int_0^1 e_{23r\bar{m}} e_{22rn} \bar{r} d\bar{r}$$

$$+ \int_0^{\bar{b}} e_{53z\bar{m}} e_{52zn} d\bar{z}$$

$$I_{33\bar{m}\hat{n}} = \int_{\bar{b}}^{\bar{H}} e_{13z\bar{m}} e_{13z\hat{n}} d\bar{z} + \int_0^{\bar{\ell}} e_{43s\bar{m}} e_{43s\hat{n}} \bar{s} \sin\alpha ds$$

$$+ \int_0^{\bar{b}} e_{53z\bar{m}} e_{53z\hat{n}} d\bar{z}$$

$$I_{33\bar{m}\bar{n}} = \int_{\bar{b}}^{\bar{H}} e_{13z\bar{m}} e_{13z\bar{n}} d\bar{z} + \int_0^1 e_{23r\bar{m}} e_{23r\bar{n}} \bar{r} d\bar{r}$$

$$+ \int_0^{\bar{\ell}} e_{43s\bar{m}} e_{43s\bar{n}} \bar{s} \sin\alpha ds + \int_0^{\bar{b}} e_{53z\bar{m}} e_{53z\bar{n}} d\bar{z}$$

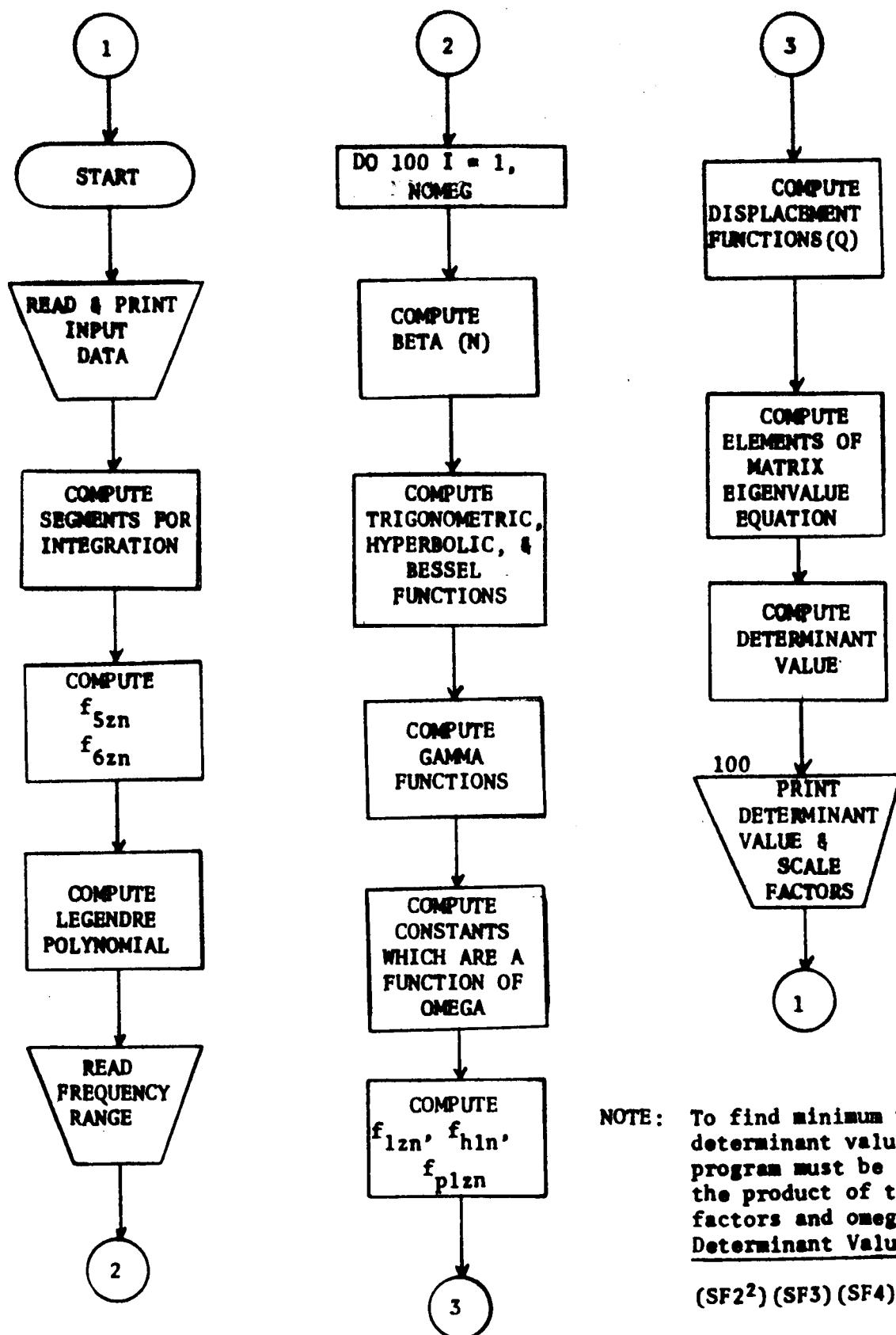


C. COMPUTER PROGRAM DESCRIPTION

A matrix eigenvalue problem of the longitudinal oscillation of a liquid-filled elastic cylindrical tank with a flexible inverted conical bulkhead is formulated in Chapter IV and a procedure for the numerical calculation is presented in Chapter V. The FORTRAN computer program for this matrix eigenvalue problem is described in this appendix.

The purpose of the program is to determine the natural frequencies and the corresponding mode shapes of the system described above. A numerical example is worked out with the use of the first three terms of the series of the velocity potentials. The input data for this numerical example are shown in Table 1. The flow chart of the program is also shown in Figure 5.

Figure 5. Frequencies and Modes Program Logic

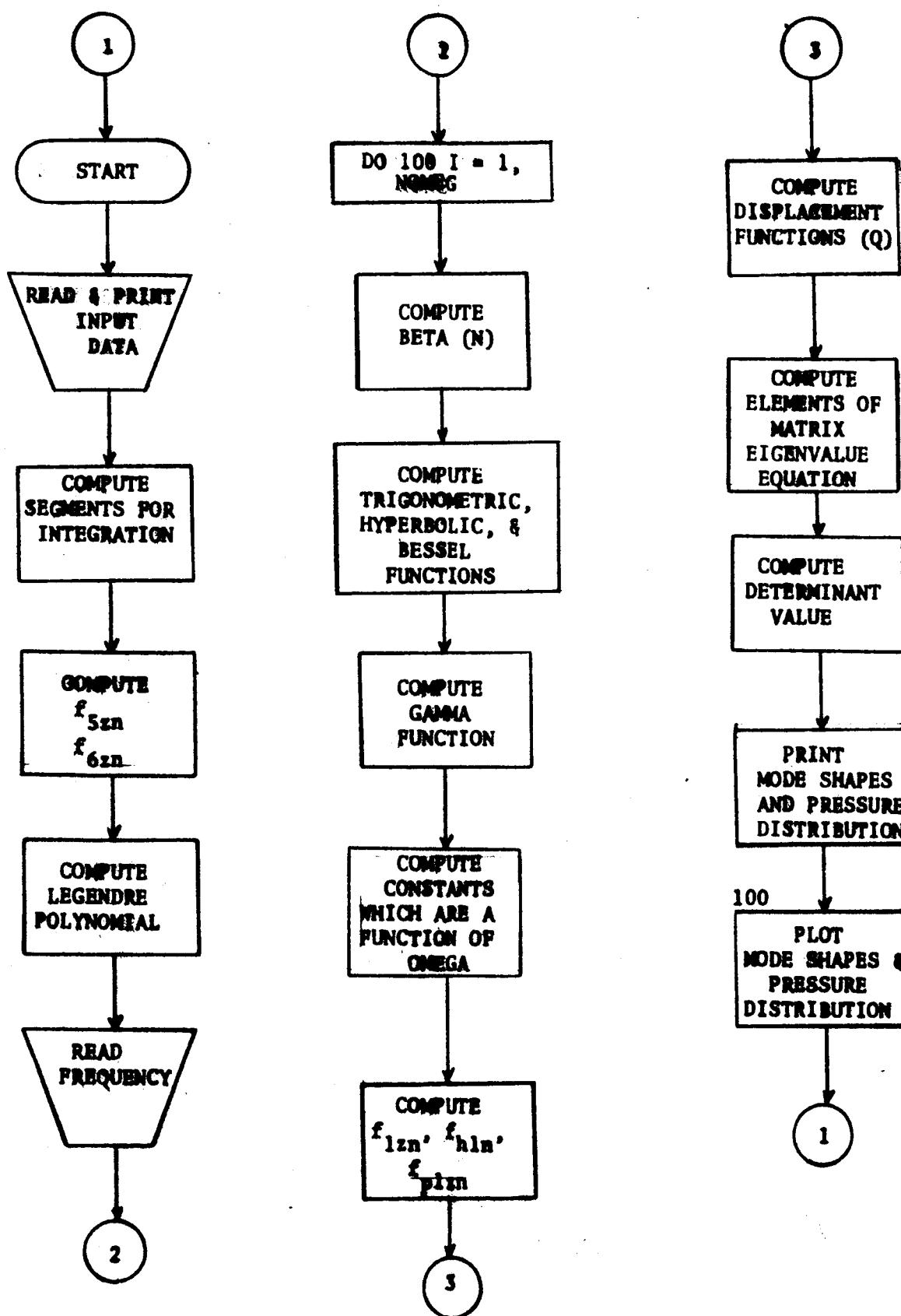


NOTE: To find minimum values, the determinant values from the program must be divided by the product of the scale factors and ω^{10} , i.e.
Determinant Value

$$(SF_2^2)(SF_3)(SF_4)(\omega^{10})$$

(a) Flow Diagram for Evaluation of Frequencies

Figure 5. Frequencies and Modes Program Logic



(b) C

(b) Flow Diagram for Evaluation of Mode Shapes

```

$16FTC OSCI
COMMON/DAT2/ NITER,N,NSEG
COMMON/DMAT/ CMAT(14,14)
COMMON/DAT3/ ALPHAN(5), B(5), U(5)
COMMON/Q/ Q1(12,11), Q2(12,11), Q3(12,11), Q4(6,11)
COMMON/SH/ BMAT(14,14)
DIMENSION AMAT(14,14), OMEGA(50)
LOGICAL ERR
1 CALL INPUT
1 CALL CONST
CALL F5ZN
CALL POLY
READ(5,10) NOMEGL(OMEGA(I)), I=1,NOMEGL
10 FORMAT(1I12/(6E12.8))
DO 100 I=1,NOMEGL
F = 1.
OMEG = OMEGA(I)
CALL BETAN(OMEG,ERR)
IF(ERR) GO TO 90
CALL TRANSF
CALL GAMAN
CALL COEF(OMEG)
CALL F1ZN
CALL FH1N(OMEG)
CALL FP1ZN
CALL QFUNCT
CALL EPSLON
CALL CMATX
CALL CMATX
C *** DELETE ROW-4 AND COL.-4 ***
DO 15 J=1,3
DO 15 K=1,14
15 AMAT(K,J) = CMAT(K,J)
DO 20 K=1,14
DO 20 J=5,14
20 AMAT(K,J-1) = CMAT(K,J)
DO 30 K=5,14
DO 30 J=1,13
30 AMAT(K-1,J) = AMAT(K,J)
WRITE(6,40)
40 FORMAT(1H1,45X,18HDETERMINANT MATRIX)
DO 45 K=1,13
45 WRITE(6,50) K, (AMAT(K,J), J=1,13)
50 FORMAT(1H0,1X,4HROW=, 12, 2X,1P6E16.4/(10X,1P6E16.4))

```

```

CW = (5./OMEG1)**6 OS390372
XX = CW**4 OSC10373
SF3 = 1.0E-6 *CW OSC10374
SF2 = 1.0E-4 * CW OSC10376
SF4 = 1.0E-3 * CW OSC10378
DO 52 J=1,13 OSC10380
AMAT(2,J) = SF2 * AMAT(2,J) OSC10382
AMAT(3,J) = SF3 * AMAT(3,J) OSC10384
AMAT(4,J) = SF4 * AMAT(4,J) OSC10386
52 AMAT(5,J) = SF2 * AMAT(5,J) OSC10388
DO 54 K = 1,10 OSC10389
DO 54 J = 1,13 OSC10390
54 BMAT(K,J) = AMAT(K,J) OSC10391
DO 55 K=T2,T3 OSC10392
DO 55 J = 1,13 OSC10393
55 BMAT(K-1,J) = AMAT(K,J) OSC10394
KK = IDETRM(14,13,AMAT,F) OSC10397
IF(KK.EQ.2) GO TO 110
WRITE(6,60) OMEG,F OSC10410
60 FORMAT(1H-,TOX,BHOMEKA = ,1PE13.6, 10X,20HDETERMINANT VALUE = ,1P0SC10420
1E13.6) OSC10430
WRITE(6,65) SF2,SF3, SF4, XX OSC10434
65 FORMAT(1H-,2X,6HSF2 = ,1PE13.6,5X,6HSF3 = ,1PE13.6, 5X, 6HSF4 = ,0SC10436
1 1PE13.6,3X,8HCW**4 = ,1PE13.6) OSC10438
GO TO 100 OSC10440
90 WRTTET6,95) OSC10490
95 FORMAT(1H-,10X,32H*** NO SOLUTIONS FOR BETA(I) ***)
100 CONTINUE OSC10500
GO TO 1 OSC10520
110 WRITE(6,120) OSC10525
120 FORMAT(1H-,5X,40HOVERFLOW OCCURED CALCULATING DETERMINANT) OSC10527
END OSC10530
5* OSC10540

```

```

$IBFTC OSCI
COMMON/DAT2/ NITER,N,NSFG
COMMON/DMAT/ CMAT(14,14)
COMMON/DAT3/ ALPHAN(5), B(5),U(5)
COMMON/Q/ Q1(12,11), Q2(12,11),Q3(12,11), Q4(6,11)
COMMON/SH/ BMAT(14,14)
DIMENSION AMAT(14,14), OMEGA(50)
LOGICAL ERR
1 CALL INPUT
CALL CONST
CALL F5ZN
CALL PULY
READ(5,10) NOMEGL(OMEGA(I), I=1,NOMEGL)
10 FORMAT(1I2/(6E12.8))
DC 100 I=1,NOMEGL
F = 1.
CMEG = OMEGA(I)
CALL BETAN(OMEG,ERR)
IF(ERR) GO TO 90
CALL TRANSF
CALL GAMAN
CALL COEF(OMEG)
CALL F1ZN
CALL FH1N(OMEG)
CALL FP1ZN
CALL QFUNCT
CALL EPSLON
CALL CMATX
C *** DELETE ROW-4 AND COL.-4 ***
DO 15 J=1,3
DO 15 K=1,14
15 AMAT(K,J) = CMAT(K,J)
DO 20 K=1,14
DO 20 J=5,14
20 AMAT(K,J-1) = CMAT(K,J)
DO 30 K=5,14
DO 30 J=1,13
30 AMAT(K-1,J) = AMAT(K,J)
WRITE(6,40)
40 FORMAT(1H1,45X,18HDETERMINANT MATRIX)
DO 45 K=1,13
45 WRITE(6,50) K, (AMAT(K,J),J=1,13)
50 FORMAT(1H0,1X,4HROW=, 12, 2X,1P6E15.4/(1.0X,1P6E16.4))

```

```

CW = (5./OMEG)*6          05390372
XX = CW**4                OSC10373
SF3 = 1.0E-6 *CW           OSC10374
SF2 = 1.0E-4 * CW          OSC10376
SF4 = 1.0E-3 * CW          OSC10378
DO 52 J=1,13               OSC10380
      AMAT(2,J) = SF2 * AMAT(2,J)    OSC10382
      AMAT(3,J) = SF3 * AMAT(3,J)    OSC10384
      AMAT(4,J) = SF4 * AMAT(4,J)    OSC10386
52   AMAT(5,J) = SF2 * AMAT(5,J)    OSC10388
      DO 54 K = 1,10             OSC10389
      DO 54 J = 1,13             OSC10390
      DO 55 K=12,13             OSC10391
      DO 55 J = 1,13             OSC10392
      DO 55 K=12,13             OSC10393
55   BMAT(K,J) = AMAT(K,J)       OSC10394
      KK = IDETRM(14,13,AMAT,F)     OSCJ0397
      IF(KK.EQ.2) GO TO 110
      WRITE(6,60) OMEG,F
      60 FORMAT(1H-,10X,B'OMEGA = ,IPET3.6,10X,Z'DETERMINANT VALUE = ,IPOSET420
      1E13.6)
      WRITE(6,65) SF2, SF3, SF4, X
      65 FORMAT(1H-,2X,6HSF2 = ,1PE13.6,5X,6HSF3 = ,1PE13.6, 5X, 6HSF4 = ,0SC10436
      1 1PE13.6,3X,8HCW**4 = ,1PE13.6)
      CALL SHAPE(OMEG)
      GO TO 100
      WRITE(6,95)
      95 FORMAT(1H-,10X,32H*** NO SOLUTIONS FOR BETA(1) ***)
      100 CONTINUE
      GO TO 1
      110 WRITE(6,120)
      120 FORMAT(1H-,5X,40HOVERFLOW OCCURED CALCULATING DETERMINANT)    OSC10526
      GO TO 1
      END
      OSC10527
      OSC10530
      OSC10540
$*

```

```

*IBFTC QFUNC• SDD
C *** COMPUTE Q FUNCTION ***
SUBROUTINE QFUNC
COMMON/DAT1/RADIUS,HTCON,AL,YOUNG,HTCY,RHO,RHOPRO,POTSON,ALPHA,D,
1 THICK
COMMON/TRIG/SINZ(5,11),CSINZ(5,11),SINHZ(5,11),CSINHZ(5,11),
1 CSINBN(5),CSINHR(5),SINHR(5)
COMMON/RESSFL/JIAN(5),JOAN(5),TUBN(5),TJOANR(5,11)
COMMON/FL/AL11,AL12,AL13,AL14,ALK11,ALK12,ALK21,ALK22,WK1,ALK0,
1ALK5,ALK6
COMMON/PAR/HBAR,BBAR,CAP,G11,E10,G12,RHOBAR
COMMON/DAT2/NITER,N,NSEG
COMMON/F1/F12(24,11)
COMMON/H1/FH1(8,6)
COMMON/FPZ/FP12(24,11)
COMMON/F56Z/F5Z(6,11),F6Z(6,11)
COMMON/DAT3/ALPHAN(5),B(5),U(5)
REAL JOANR
COMMON/Q/Q(12,11),Q2(12,11),Q3(12,11),Q4(6,11)
COMMON/PLDR/PO(6),PA(6),PZ(6,11)
COMMON/DELTA/DZRH(11),DZS(11),DSL(11),DR(11),DZB(11),DS(11)
DIMENSION P(10),SRAR(11),BFSI(5,11)
IF(ALPHA.GT.89.) GO TO 2
BX = (SIND(ALPHA) / COSD(ALPHA))**2
GO TO 4
2 BX = 0.
4 N2 = N+1
NSEG2 = NSEG + 1
DENOM = 1. - WK1
A = POISON * ALK0 / DENOM
B = POISON * G12 / DENOM
C = (1. + POISON * G12) / DENOM
CX = .POISON * RHORAR / DENOM
DFLZ = RBAR / FLOAT(NSEG)
DZ = 0.
DO 25 J=1,NSEG2
RAD = G11 * DZ
Q1(1,J) = -G11 * SIN(RAD)
Q2(2,J) = G11 * COS(RAD)
25 DZ = DZ + DELZ
K1 = 1
DO 50 I=1,N2

```

```

QFUNC000
QFUNC005
QFUNC010
QFUNC020
QFUNC030
QFUNC032
QFUNC034
QFUNC036
QFUNC040
QFUNC042
QFUNC048
QFUNC050
QFUNC060
QFUNC070
QFUNC080
QFUNC090
QFUNC100
QFUNC105
QFUNC117
QFUNC118
QFUNC119
QFUNC120
QFUNC124
QFUNC126
QFUNC130
QFUNC140
QFUNC150
QFUNC160
QFUNC170
QFUNC171
QFUNC172
QFUNC174
QFUNC176
QFUNC178
QFUNC180
QFUNC182
QFUNC184
QFUNC188
QFUNC189
QFUNC190

```

```

X = ALK11 * FH1(1,I)
Y = ALK12 * FH1(3,I)
Q11 = A * ( X + Y + ALK12 * X + ALK12 * Y )
X2 = ALK11 * FH1(2,I)
Y2 = ALK12 * FH1(4,I)
Q12 = A * ( X2 + Y2 + ALK12 * X2 + ALK12 * Y2 )
Q15 = RHOBAR * A * ALK12
Q13 = ALK21 * FH1(5,I) + ALK22 * FH1(7,I)
Q14 = ALK21 * FH1(6,I) + ALK22 * FH1(8,I)
DO 40 J =1,NSFG2
Q1(K,J) = RHOBAR * (Q11 + B * (AL13 * F12(K,J) + AL14 * F12(K+2,J)) + C * F12(K,J))
Q1(K+1,J) = RHOBAR * (Q12 + B * (AL13 * F12(K+1,J) + AL14 * F12(K+2,J)) + C * F12(K+1,J))
Q1(K+2,J) = Q13 * Q15
Q1(K+3,J) = Q14 * Q15
X5 = ALC5 * Q3(2,J)
X6 = ALC6 * Q3(2,J)
Q2(K,J) = -QX * X5 * (X + Y)
Q2(K+1,J) = -QX * X5 * (XZ + YZ)
Q2(K+2,J) = -QX * ((Q3(1,J) - X6) * FH1(5,I) + X5 * FH1(7,I) + X5 * QFUNC400
1 Q13) + B * RHOBAR * (Q3(1,J) * F5Z(K1,J) + Q3(2,J)*F6Z(K1,J)*G11) + QFUNC410
2 C * RHOBAR * FP12(K+2,J)
40 Q2(K+3,J) = -QX * ((Q3(1,J) - X6) * FH1(5,I) + X5 * FH1(8,I) + X5 * QFUNC430
1* Q14) + B * RHOBAR * (Q3(1,J)*F5Z(K1+1,J) + Q3(2,J)*F6Z(K1+1,J)*G1QFUNC440
21) + C*RHOBAR * FP12(K+3,J)
K1 = K1 + 2
50 K = K + 4
DO 90 I=1,NSEG2
X = B(I) * DR(J)
90 BESI(I,J) = BFSFL(X,0,2)
DO 95 J=1,NSEG2
Q3(1,J) = E10 + BBAR
Q3(2,J) = 0.
Q3(3,J) = 1.
P(1) = -5
P(2) = .375
95 Q3(4,J) = 0.
K = 5
DO 100 I=1,N
NCAP = 2 * I
DO 97 J=1,NSEG2

```

```

Q3(K,J) = (SINHR(I) + CSINHR(I)) * J(OANR(I,J))
Q2(K+1,J) = RFSI(I,J) / CSINRN(I)
Q2(K+2,J) = DR(J)*NCAP * P(I)
Q3(K+3,J) = U.

100 K = K + 4
      RX = RHOBAR * RX
      DO 102 J=1,NSFG2
      Q4(1,J) = -RBX * PA(1) * DSL(J)**2
      Q4(2,J) = -RBX * PA(2) * DSL(J)**3
102   K = 2
      DO 110 I=1,N
      NCAP = 2 * I
      NBAR = 2 * I + 1
      DO 105 J=1,NSFG2
      Q4(K,J) = -RBX * PA(NBAR) * DSL(J)**(NCAP+2)
105   Q4(K+1,J) = -RX * PA(NBAR+1) * DSL(J)**(NBAR+2)
110   K = K + 2
      RETURN
      END

```

\$*

\$IBFTC COEF.

SDD

COEF0001

SUBROUTINE COEFTOMEGA

COEF0002

COMMON/DAT1/ RADIUS,HTCON,AL,YOUNG,HTCY,RHO,RHOPRO,POISON,ALPHA,D,COEF0005

COEF0010

1 THICK

COEF0015

COMMON/EL/ AL11,AL12,AL13,AL14,ALK11,ALK12,ALK21,ALK22,WK1,ALK0,

COEF0016

1 ALC5• ALC6

COEF0017

COMMON/BAR/ HBAR,BBAR,CAP,G11,E10,G12, RHOBAR

COEF0018

OMEG2= OMEGA**2

COEF0019

E10 = (386.08970MEG2 - HTCY) / RADIUS

COEF0020

BK1 = RHO * OMEG2 * RADIUS**3 / D

COEF0022

WK1 = CAP * OMEG2

COEF0025

DENOM = 1. - POISON**2 - WK1

COEF0030

G11 = SQRT(WK1 - WK1**2) / DENOM

COEF0035

G12 = POISON / DENOM

COFF0040

X = G11 * HBAR

COFF0045

X2 = (1. - POISON**2 / (1. - WK1)) * G11

C056005

AL11 = -X2 * SIN(X)

COFF0055

AL12 = X2 * COS(X)

COFF0060

X = G11 * BBBAR

COFF0065

AL13 = -G11 * SIN(X)

COFF0070

AL14 = G11 * COS(X)

COEF0075

DENOM = AL11 * AL14 - AL12 * AL13

COFF0080

ALC1 = AL11 / DENOM

COEF0085

ALC2 = AL12 / DENOM

COFF0090

ALC3 = AL13 / DENOM

COFF0095

ALC4 = AL14 / DF NOM

COFF0100

ALC5 = 1. / AL14

COFF0105

ALC6 = AL13 / AL14

COFF0110

ALK0 = 1.

COFF0120

X1 = COS(G11 * BBBAR)

COFF0122

X2 = COS(G11 * HBAR)

COFF0124

X = G11 * (HBAR-BBAR)

COEF0126

X3 = COS(X)

COEF0128

X4 = SIN(X)

COFF0130

DN = 1. - POISON**2 / (1. - WK1)

COEF0132

ALK11 = X1 / (DN * X2)

COEF0134

ALK12 = -X1 * X3 / X2

COEF0136

ALK21 = G11 * X4 / X2

COEF0138

ALK22 = SIN(G11*BBAR) * X4 / X2

COEF0140

RETURN

COEF0150

FND

.

```

$!RFTC FP1ZN. SDN
      SURROUNTF FP1ZN
COMMON/RAR/ HBAR, RBAR, CAP, G11, F10, G12, RHOBAR
COMMON/DAT2/ NITER, N, NSEG
COMMON/FPZ/ FP1Z(24,11)
      REAL JIAN, JOAN, IOBN, I1BN, JOANR
COMMON/TRIG/ SINZ(5,11), CSINZ(5,11), SINHZ(5,11), CSINHZ(5,11),
1 CSINRN(5), CSTNRH(5), SINHB(5)           FP1ZN005
COMMON/BESSEL/ JIAN(5), JOAN(5), IUBN(5), I1BN(5), JUANR(5,11)   FP1ZN010
COMMON/PLDR/ PC(6), PA(6), PZ(6,11)          FP1ZN015
COMMON/DELTA/DZBH(11), DZS(11), DSL(11), DR(11), DZB(11), DS(11)   FP1ZN020
COMMON/LIMIT/N2, NSEG2                      FP1ZN024
DO 10 I=1,NSEG2                            FP1ZN025
      FP1Z(1,I) = E10 + DZRH(I)             FP1ZN026
      FP1Z(2,I) = U*                         FP1ZN027
      FP1Z(3,I) = PZ(1,I) * DZS(I)**2       FP1ZN030
10   FP1Z(4,I) = PZ(2,I) * DZS(I)**2       FP1ZN035
      K = 1                                     FP1ZN040
DO 20 I=1,N
      NCAP = 2 * I                           FP1ZN082
      NBAR = 2 * I + 1                        FP1ZN086
      K = K + 4                               FP1ZN086
      A = IOBN(I) / CSINBN(I)                 FP1ZN088
DO 20 J=1,NSEG2
      FP1Z(K,J) = JOAN(I) * (SINHZ(I,J) + CSINHZ(I,J))   FP1ZN090
      FP1Z(K+1,J) = A * CSINZ(I,J)            FP1ZN090
      XX = DZS(J)**2                         FP1ZN100
      FP1Z(K+2,J) = PZ(I+1,J) * XX***NCAP    FP1ZN100
20   FP1Z(K+3,J) = PZ(I+2,J) * XX***NBAR    FP1ZN110
      RETURN
      END

```

\$*

SIBFTC EPSLN SDD

EPSLN000

EPSLN005

EPSLN010

EPSLN015

EPSLN020

EPSLN025

EPSLN030

EPSLN035

EPSLN040

EPSLN045

EPSLN050

EPSLN055

EPSLN060

EPSLN065

EPSLN070

EPSLN075

EPSLN080

EPSLN090

EPSLN100

EPSLN110

EPSLN120

EPSLN130

EPSLN140

EPSLN145

EPSLN150

EPSLN160

EPSLN170

EPSLN180

EPSLN190

EPSLN200

EPSLN210

EPSLN215

EPSLN220

EPSLN230

EPSLN240

EPSLN250

EPSLN260

EPSLN270

EPSLN280

EPSLN290

EPSLN300

EPSLN310

EPSLN320

EPSLN330

EPSLN340

EPSLN350

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COMMON/DAT2/ NITER, N, NSEG
COMMON/GAMA/ GAMA1(9,11), GAMA2(12,11), GAMA3(9,11)
COMMON/Q/ Q1(12,11), Q2(12,11), Q3(12,11), Q4(6,11)
COMMON/BK/ BK1
COMMON/BAR/ HBAR, BBAR,CAP, G11, F10, G12,RHOBAR
COMMON/EPS/ E11(5,11), E12(5,11), E13(5,11), E21(5,11), E23(5,11), E31(5,11), E32(5,11), E33(5,11), E43(5,11), E43B(5,11)
1 E23(5,11), E31(5,11), E32(5,11), E33(5,11), E43(5,11), E43B(5,11)
2 E51(5,11), E52(5,11), E53(5,11), E53B(5,11)
COMMON/SK/SKZ(11), OKZ(11)
NSEG2 = NSEG + 1
DELZ = BBAR / FLOAT(NSEG)
DZ = 0.
DO 10 I=1,NSEG2
X = BBAR - DZ
SKZ(I) = 1. / SQRT(1. + X**2)
OKZ(I) = X * SKZ(I)
10 DZ = DZ + DELZ
      K2 = 1
      K3 = 1
      N2 = N + 1
      DO 30 I=1,N2
      DO 20 J=1,NSEG2
E11(I,J) = GAMA1(K,J) - BK1 * Q1(K2,J)
E12(I,J) = GAMA1(K+1,J) - BK1 * Q1(K2+1,J)
E13(I,J) = -BK1 * Q1(K2+2,J)
E21(I,J) = GAMA1(K+2,J)
E13B(I,J) = -BK1 * Q1(K2+3,J)
E23(I,J) = -GAMA3(K+2,J)
E31(I,J)= Q3TR2,J
E32(I,J) = Q3(K2+1,J)
E33(I,J) = -Q3(K2+2,J)
E51(I,J) = -BK1 * Q2(K2,J)
E52(I,J) = -BK1 * Q2(K2+1,J)
E53(I,J) = SKZ(J) * GAMA2(K2,J) + OKZ(J) * GAMA2(K2+2,J) - BK1 *
      1 Q2(K2+2,J)
E53B(I,J) = SKZ(J) * GAMA2(K2+1,J) + OKZ(J) * GAMA2(K2+3,J) - BK1
      1 * Q2(K2+3,J)
E43(I,J) = GAMA3(K,J) - BK1 * Q4(K3,J)
20 E43B(I,J) = GAMA3(K+1,J) - BK1 * Q4(K3+1,J)
      K = K+3

```

K2 = K2+4
30 K3 = K3+2
RETURN
END

EPSLN360
EPSLN370
EPSLN400
EPSLN500

```

$IBFTC SHAPE.
C*** ROUTINE CALCULATE AND PLOTS MODESHAPES ***
SUBROUTINE SHAPE(OMEG)
COMMON/DMAT/ CMAT(14,14)
COMMON/SK/ SKZ(11),OKZ(11)
COMMON/GAMA/ GAMA1(9,11), GAMA2(12,11), GAMA3(9,11)
COMMON/RK/ RK1
COMMON/LIMIT/ N2,NSEG2
COMMON/DAT3/ ALPHAN(5),B(5),U(5)
COMMON/PLDR/ PO(6),PA(6),PZ(6,11)
COMMON/DELTA/ DZBH(11),DZS(11),DSL(11),DR(11),DZB(11),DS(11)
COMMON/TRIG/ SINZ(5,11),CSINZ(5,11),SINHZ(5,11),CSINHZ(5,11),
1 CSINBN(5),CSINHB(5),SINHB(5)
COMMON/BESSEL/ JIAN(5),JOAN(5),TOBN(5),TIBN(5),JOANR(5,11)
COMMON/RAR/ HRAR,RRAR,CAP,G11,F10,G12,RHOBAR
COMMON/DAT1/RADIUS,HTCON,AL,YOUNG,HTCY,RHO,RHOPRO,POTSON,ALPHA,SHAPE048
1 D, THICK
COMMON/Q/ Q1(12,11), Q2(12,11), Q3(12,11), Q4(6,11)
COMMON/SH/ BMAT(14,14)
CALL SETMIV(48,48,48,48)
DIMENSION Y(30),X(30),DPHI(30)
DIMENSION C(12),XA(12),A(12,12)
REAL JIAN, JOAN, IOBN, JOANR
CPS = OMEG / 6.283185
WRITE(6,5) OMEG, CPS
5 FORMAT(1H1,10X,12HFREQUENCY = ,1PE13.6,10H RAD.7SEC.,10X,12HFREQUESHAPE066
1NCY = ,1PE13.6,7H C.P.S.)
NN = 4 * N2
DO 10 I=1,NN
C(I) = -BMAT(I,1)
DO 10 J=1,NN
10 A(I,J) = BMAT(I,J+1)

C COMPUTE A(N),B(N),C(N)
K = ISIIMEQ(12,NN,1,A,C,O,XA)
IF(K-2) 20,100,120
C*** STORE A(N), B(N), C(NCAP), C(NBAR), IN CMAT BY COLUMNS ***
20 CMAT(1,1) = 1.
CMAT(1,2) = 0.
N = N2 - 1
DO 25 I=1,N
CMAT(I+1,1) = A(I,1)
K = N + I

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25 CMAT(I+1,2) = A(K,1)           SHAPE220
K = ? * N                         SHAPE230
K1 = K + N2                        SHAPE240
D O 30 I=1,N2                      SHAPE250
J = K + I                          SHAPE260
CMAT(I,3) = A(J,1)                SHAPE270
J = K1 + I                         SHAPE280
C*** STORE MODE SHAPES OF W11 IN FIRST ROW OF MATRIX A ***          SHAPE290
C*** STORE MODE SHAPES OF W11 IN THE SECOND ROW OF MATRIX A ***          SHAPE300
C*** STORE MODE SHAPES OF W11 IN THE THIRD ROW OF MATRIX A ***          SHAPE305
C*** STORE MODE SHAPES OF W11 IN THE FOURTH ROW OF MATRIX A ***          SHAPE307
DO 50 J=1,NSEG2                     SHAPE310
K = 1                               SHAPE310
SUM = 0.                            SHAPE320
SUM2 = 0.                           SHAPE330
DO 40 I=1,N2                       SHAPE340
SUM = Q2(K,J) * CMAT(I,1) + Q2(K+1,J) * CMAT(I,2) + Q2(K+2,J) * CMAT(I,3) + Q2(K+3,J) * CMAT(I,4) + SUM
1 AT(I,3) + Q2(K+3,J) * CMAT(I,4) + SUM
SUM2 = Q1(K,J) * CMAT(I,1) + Q1(K+1,J) * CMAT(I,2) + Q1(K+2,J) * CMAT(I,3) + Q1(K+3,J) * CMAT(I,4) + SUM2
1 MAT(I,3) + Q1(K+3,J) * CMAT(I,4) + SUM2
40 K = K + 4
A(1,J) = SUM
50 A(2,J) = SUM2
C*** STORE MODE SHAPES OF W2 IN THE THIRD ROW OF MATRIX A ***          SHAPE380
DO 70 J=1,NSEG2                     SHAPE390
K = 1                               SHAPE400
SUM = 0.                            SHAPE410
SUM2 = 0.                           SHAPE420
DO 60 I=1,N2                       SHAPE430
SUM = Q4(K,J) * CMAT(I,1) + Q4(K+1,J) * CMAT(I,2) + Q4(K+2,J) * CMAT(I,3) + Q4(K+3,J) * CMAT(I,4) + SUM
60 K = K+2
70 A(3,J) = SUM
WRITE(6,75) (I,CMAT(I,1),I,CMAT(I,2),I,CMAT(I,3),I,CMAT(I,4),I,I,NSHAPE510
12)
75 FORMAT(1H-/ (1H0,10X,2HA(,I1*4H) = ,1PF13.6/ 1H0,2HR(,I1*4H) = SHAPE530
1,1PF13.6/ 1H0,10X,2HC(,I1,4H) = ,1PE13.6/ 1H0,10X,5HCBAR(,I1,4H) SHAPE540
2 = ,1PE13.6)
WRITE(6,80) (A(1,J),J=1,NSEG2)           SHAPE550
80 FORMAT(1H1,45X,14HW12 MODE SHAPE/ (1H0,3X,1P6E17.5))
WRITE(6,85) (A(2,J),J=1,NSEG2)           SHAPE560
85 FORMAT(1H-,45X,14HW11 MODE SHAPE/ (1H0,3X,1P6E17.5))
WRITE(6,90) (A(3,J),J=1,NSEG2)           SHAPE570
90 FORMAT(1H-,45X,13HW2 MODE SHAPE/ (1H0,3X,1P6E17.5))           SHAPE580

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DO 97 J= 1,NSEG2
C(J) = A(3,J)
X(J) = DZB(J)
DPHI(J)=((GAMA2(3,J) * CMAT(1,3) + GAMA2(4,J) * CMAT(1,4) + GAMA2(5,J) * CMAT(1,5) + GAMA2(6,J) * CMAT(1,6)) / BK1
17,J) * CMAT(2,3) + GAMA2(8,J) * CMAT(2,4) + GAMA2(11,J) * CMAT(3,3) SHAPE605
2) + GAMA2(12,J) * CMAT(3,4) * OKZ(J) + (GAMA2(I,J) * CMAT(I,3)) / SHAPE606
3+ GAMA2(2,J) * CMAT(1,4) + GAMA2(5,J) * CMAT(2,3) + GAMA2(6,J) * CMAT(3,4) * CSHAPE608
4MAT(2,4) + GAMA2(9,J) * CMAT(3,3) + GAMA2(10,J) * CMAT(3,4)) * SKZT SHAPE609
5J)) / BK1
SHAPE610
92 Y(J) = A(1,J)
SHAPE611
K = NSEG2
DO 93 J=2,NSEG2
K = K + 1
X(K) = DZBH(J)
DPHI(K) =(GAMA1(1,J) + GAMA1(4,J) * CMAT(2,1) + GAMA1(7,J) * CMAT(3,1) + GAMA1(8,J) * CMAT(3,2)) / RK1
13,1) + GAMA1(5,J) * CMAT(2,2) + GAMA1(9,J) * CMAT(3,2)) / RK1
SHAPE622
SHAPE623
93 Y(K) = A(2,J)
WRITE(6,94) (DPHI(I),I=1,21)
SHAPE624
94 FORMAT(1H1,4UX,5HPhi/R/(1H0,3X,6E17.5))
CALL GRAPH(3,42,-21,X,Y,1H,1H)
SHAPE625
DO 99 J=1,NSEG2
99 DPHI(J) =(GAMA3(1,J) * CMAT(1,3) + GAMA3(2,J) * CMAT(1,4) + GAMA3(3,J) * CMAT(1,5) + GAMA3(4,J) * CMAT(2,3) + GAMA3(5,J) * CMAT(2,4) + GAMA3(7,J) * CMAT(3,3) SHAPE655
2) + GAMA3(8,J) * CMAT(3,4)) / BK1
CALL GRAPH(3,42,-11,DSL,C,1H,1H)
SHAPE660
C *** COMPUTE PRESSURE DISTRIBUTION AND PLOT ON CRT ***
CALL GRAPH(3,42,-11,DSL,C,1H,1H)
SHAPE665
DO 95 I=1,11
X1 = CMAT(2,1) * (SINHB(1) + CSINHB(1))
X2 = CMAT(3,1) * (SINHB(2) + CSINHB(2))
GAW = 386.089 / (RADIUS * OMEG**2) - HBAR + BBAR
SHAPE710
DO 95 I=1,11
Y(I) = X1 * JOANR(1,I) + X2 * J0ANR(2,I) + CMAT(2,2) * BESEL(R(I)*SHAPE828
1DR(I),0,2) / CSINRN(1) + CMAT(3,2) * RESELT(BT2)*DR(I),0,2) / CSTNSHAPE829
2BN(2) + GAW
SHAPE830
95 C(I) = CMAT(1,3) + CMAT(2,3) * PO(3) * DR(I)**2
CALL GRAPH(3,42,-11,DR,C,1H,1H)
SHAPE831
WRITE(6,96) (Y(I),I=1,11)
SHAPE839
96 FORMAT(1H1,37X,30HTHETA(1) PRESSURE DISTRIBUTION/ (1H0,3X,1P6E17.5SHAPE841
1))
SHAPE840
WRITE(6,97) (C(I),I=1,11)
SHAPE842
97 FORMAT(1H-,37X,30HTHETA(2) PRESSURE DISTRIBUTION/ (1H0,3X,1P6E17.5SHAPE846
1))
SHAPE848
RFTURN
SHAPE850
100 WRITE(6,110)
SHAPE852

```

```
110 FORMAT(1H-,10X,53H*** ILL CONDITIONED MATRIX - EXECUTION TERMINATESHAPE854  
1D ***)  
STOP  
120 WRITE(6,125)  
125 FORMAT(1H-, 10X,46H** MATRIX IS SINGULAR - EXECUTION TERMINATED**) SHAPE880  
STOP  
SHAPE900  
END
```

SIBFTC ROOTD* SDD

ZEROS OF AN ARBITRARY FUNCTION

C ROOT SEARCHES FOR ZEROS BY SAMPLING F(X) AND FITTING L(X) TO

C THE SAMPLED VALUES. L(X) IS USED TO PREDICT F(X). IF THE
 C PREDICTION IS BAD (SEE REL), DX IS DECREASED. IF THE PREDICTION
 C IS VERY GOOD, DX IS INCREASED. WHEN A PREDICTION IS GOOD, X IS
 C INCREASED BY DX. WHEN F(X) CHANGES SIGN, A ROOT HAS BEEN
 C ISOLATED, AND X IS INTERPOLATED TO AN ACCURACY OF TOL.

C INPUT -- X = START OF SEARCH.

C DX = INCREMENT IN X (+ OR -).

C F = NAME OF THE ARBITRARY FUNCTION.

C TOL = RELATIVE ERROR IN THE ANSWER (X).

C REL = MAXIMUM RELATIVE ERROR IN $|F(X)-L(X)|$.

C LIMIT = MAXIMUM NUMBER OF ITERATIONS TRIED.

C OUTPUT -- NTRY = 1, EXACT ROOT FOUND.

C = 2, APPROXIMATE ROOT FOUND.

C = 3, LOCAL /MINIMUM/ FOUND.

C = 4, DISCONTINUITY FOUND.

C = 5, LIMIT REACHED (NO ROOT FOUND).

C X = ROOT (OR CURRENT VALUE).

C DX = ADJUSTED (CURRENT) VALUE.

C Y = F(X).

C*****SUBROUTINE ROOT (NTRY,X,DX,Y,F,TOL,REL,LIMIT)

EXTERNAL F
 INTEGER NTRY,LIMIT
 REAL REL,TOL,DX

C-----
 LOGICAL SLOW, MED, FAST, BAD, FIRST
 C-----
 C LAGRANGIAN POLYNOMIAL

L(X) = $(X-X1)(X-X2)\dots(X-XN)/(X1-X0)(X2-X0)\dots(XN-X0)$ *Y1
 C*****
 C INITIALIZATION
 OLDDX = 0
 RATIO = 2.0

```

FIRST = •TRUE•
ITER = 1
X0 = X
Y0 = F(X0)
IF (Y0 •EQ. 0.0) GO TO 60
C-----  

9 IF ((NOT•FIRST)) GO TO 20
10 IF (X0+DX •EQ. XC) GO TO 80
IF (ITER •GE. LIMIT) GO TO 90
ITER = ITER+1
X1 = X0+DX
Y1 = F(X1)
IF (X1 •EQ. X0) GO TO 80
DX = X1-X0
IF (SIGN(1.0,Y0) •EQ. SIGN(1.0,Y1) •AND. Y1 •NE. 0.0) GO TO 20
DX = DX/4.0
GO TO 10
C-----  

C SEARCH
20 IF ((X1+DX •EQ. X1) GO TO 80
IF (ITER •GE. LIMIT) GO TO 90
ITER = ITER+1
X = X1+DX
Y = F(X)
IF (X •EQ. X1) GO TO 80
DX = X-X1
ERROR = ABS(Y-L(X))
30 IF (ERROR •NE. 0.0) GO TO 31
DX = DX*2.0
GO TO 32
31 D = (X0-X1)/2.0
R = ABS(Y)*(X-X0)*(X-X1)/ERROR
DX = AMAX1(DX/4.000,AMIN1(DX*2.000,D+ SGRT(D*D+R*REL/2.000)))
IF (ERROR •GT. REL* ABS(Y)) GO TO 9
32 IF (Y •EQ. 0.0) GO TO 60
IF (SIGN(1.0,Y) •NE. SIGN(1.0,Y1)) GO TO 46
IF ( ABS(Y1) •GE. ABS(Y) •OR. ABS(Y) •LT. ABS(Y1)) GO TO 35
C-----  

IF (OLDDX •EQ. 0.0 ) OLDDX = X-X1
IF ( ABS(X-X1) •LT. TOL* ABS(X) ) GO TO 70
DX = (X-X1)/4.0
GO TO 10

```

```

C----- 35 X0 = X1          00000318
      Y0 = Y1          00000320
      X1 = X           00000330
      Y1 = Y           00000340
      FIRST = •FALSE•  00000350
      GO TO 20         00000360
C----- 40 SLOW = RATIO •LT• 2•0  00000370
      RATIO = -SNGL(YY/Y)  00000371
      MED = •NOT• SLOW •AND• RATIO •LE• 4•0  00000372
      FAST = •NOT• SLOW •AND• RATIO •GT• 4•0  00000373
      IF (FAST) X = XX+(X0-X1)*(Y1+Y0)/(Y1-Y0) 00000381
      IF (MED) X = XX+(X0-X1)*(YY/(Y1-Y0)) 00000382
      BAD = X •LE• AMIN1(X0,X1) •OR• X •GE• AMAX1(X0,X1) 00000383
      IF (SLOW •OR• BAD) X = (X0+X1)/2•0 00000390
      Y = F(X)          00000391
      IF (Y •EQ• 0•0) GO TO 60 00000392
      IF (SIGN(1•0,Y) •EQ• SIGN(1•0,Y0)) GO TO 45 00000393
C----- RATIO = SNGL(Y1/Y) 00000394
      X1 = X           00000400
      Y1 = Y           00000410
      XX = X0          00000411
      YY = Y0          00000420
      GO TO 50         00000430
C----- 45 RATIO = SNGL(YY/YY) 00000440
      46 X0 = X         00000450
      Y0 = Y           00000460
      XX = X1          00000470
      YY = Y1          00000480
      GO TO 50         00000490
C----- 50 IF ( ABS(X0-X1) •GT• TOL* ABS(X)) GO TO 40 00000500
      NTRY = 2          00000510
      RETURN            00000515
      60 NTRY = 1          00000520
      RETURN            00000530
      70 DX = OLDDX    00000540
      NTRY = 3          00000550
      RETURN            00000551
      NTRY = 3          00000552

```

RETURN
NTRY = 4
80 RETURN
NTRY = 5
90 RETURN
END

00000553
00000560
00000570
00000580
00000590
00000600

\$IBFTC CMAT.

SUBROUTINE CMATX
COMMON/Q/ Q1(12,11), Q2(12,11), Q3(12,11), Q4(6,11)

COMMON/DMAT/ CMAT(14,14)
COMMON/DAT1/ RADIUS, HTCON, AL,YOUNG, HTCY, RHOU, RHOPRO, POISON,

1 ALPHA, D, THICK
COMMON/DAT2/ NITER,N, NSEG

COMMON/FPS/ E11(5,11), E12(5,11), E13(5,11), E14(5,11), E21(5,11), E22(5,11), E23(5,11), E31(5,11), E32(5,11), E33(5,11), E43(5,11), E44(5,11), E51(5,11), E52(5,11), E53(5,11), E54(5,11), E55(5,11)

COMMON/BAR/ HBAR,BBAR,CAP,G11,E12,RHOBAR
COMMON/LIMIT/ N2,NSEG2
COMMON/DELTA/ DZBH(11), DZS(11), DSL(11),DR(11),DZB(11),DS(11)

COMMON/DEL/A,B,C,X
COMMON/TERM/ T(4,11), ISEG,KSEG, NN

KSEG = NSEG2
ISEG = NSEG

NN = NSEG2 - 3
C *** COMPUTE I11 MATRIX ***

DO 20 I=1,N2
DO 20 J=1,N2

DO 15 K=1,NSEG2
T(1,K) = E11(I,K) * E11(J,K) * DR(K)

T(2,K) = E21(I,K) * E21(J,K) * DR(K)

T(3,K) = E31(I,K) * E31(J,K) * DR(K)

15 T(4,K) = E51(I,K) * E51(J,K)

20 CMAT(I,J) = SRULE(4,A,X,X,B)

C *** COMPUTE I12 MATRIX ***

DO 30 I=1,N2
DO 30 J=1,N2

L = N2 + J
DO 25 K=1,NSEG2
T(1,K) = E11(I,K) * E12(J,K)

T(2,K) = E21(I,K) * E32(J,K) * DR(K)

25 T(3,K) = E51(I,K) * E52(J,K)

30 CMAT(I,L) = SRULE(3,A,X,B)

C *** COMPUTE I13(CAP N) MATRIX ***

N4 = 2 * N2
DO 40 I=1,N2
DO 40 J=1,N2

L = N4 + J
DO 35 K=1,NSEG2

T(1,K) = E11(I,K) * E13(J,K)

CMAT0000

CMAT0005

CMAT0009

CMAT0010

CMAT0015

CMAT0020

CMAT0025

CMAT0030

CMAT0040

CMAT0050

CMAT0060

CMAT0070

CMAT0067

CMAT0072

CMAT0065

CMAT0100

CMAT0122

CMAT0170

CMAT0175

CMAT0180

CMAT0190

CMAT0200

CMAT0210

CMAT0220

CMAT0230

CMAT0240

CMAT0250

CMAT0255

CMAT0260

CMAT0270

CMAT0280

CMAT0290

CMAT0300

CMAT0310

CMAT0320

CMAT0330

CMAT0340

CMAT0350

CMAT0360

CMAT0370

CMAT0380

CMAT0390

CMAT0400

```

T(2,K) = E31(I,K) * E33(J,K) * DR(K)          CMAT0410
T(3,K) = E51(I,K) * E53(J,K)          CMAT0420
35 CMAT(I,L) = SRULE(3,A,X,B)          CMAT0430
40 COMPUTE I13(BAR N) MATRIX ***          CMAT0440
      N5 = N4 + N2          CMAT0450
DO 50 I=1,N2          CMAT0460
DO 50 J=1,N2          CMAT0470
L = N5 + J          CMAT0480
DO 45 K = 1,NSEG2          CMAT0490
T(1,K) = E11(I,K) * E13B(J,K)          CMAT0500
T(2,K) = E21(I,K) * E23(J,K) * DR(K)          CMAT0510
45 T(3,K) = E51(I,K) * E53B(J,K)          CMAT0520
50 CMAT(I,L) = SRULF(3,A,X,B)          CMAT0530
C *** COMPUTE I122 MATRIX ***          CMAT0540
DO 60 I=1,N2          CMAT0550
L1 = N2 + I          CMAT0560
DO 60 J=1,N2          CMAT0570
L2 = N2 + J          CMAT0580
DO 55 K=1,NSEG2          CMAT0590
T(1,K) = E12(I,K) * E12(J,K)          CMAT0600
T(2,K) = E32(I,K) * E32(J,K) * DR(K)          CMAT0610
55 T(3,K) = E52(I,K) * E52(J,K)          CMAT0620
60 CMAT(L1,L2) = SRULE(3,A,X,B)          CMAT0630
C *** COMPUTE I123(CAP N) MATRIX          CMAT0640
DO 70 I=1,N2          CMAT0650
L1 = N2 + I          CMAT0660
DO 70 J=1,N2          CMAT0670
L2 = N4 + J          CMAT0680
DO 65 K=1,NSEG2          CMAT0690
T(1,K) = E12(I,K) * E13(J,K)          CMAT0700
T(2,K) = E32(I,K) * E33(J,K) * DR(K)          CMAT0710
65 T(3,K) = E52(I,K) * E53(J,K)          CMAT0720
70 CMAT(L1,L2) = SRULE(3,A,X,B)          CMAT0730
C *** COMPUTE I123(BAR N) MATRIX ***          CMAT0740
DO 80 I=1,N2          CMAT0750
L1 = N2 + I          CMAT0760
DO 80 J=1,N2          CMAT0770
L2 = N5 + J          CMAT0780
DO 75 K = 1,NSFG2          CMAT0790
T(1,K) = E12(I,K) * E13B(J,K)          CMAT0800
T(2,K) = E52(I,K) * E53B(J,K)          CMAT0810
75 T(3,K) = 0.          CMAT0820
80 CMAT(L1,L2) = SRULE(3,A,X,B)          CMAT0830

```

```

C *** COMPUTE I33(CAP N) MATRIX ***
DO 90 I=1,N2
L1 = N4 + 1
DO 90 J=1,N2
L2 = N4 + J
DO 85 K=1,NSEG2
T(1,K) = E13(I,K) * E13(J,K)
T(2,K) = E33tt(K) * E33t(J,K) * DR(K)
T(3,K) = E43(I,K) * E43(J,K) * DS(K)
85 T(4,K) = E53(I,K) * E53(J,K)
90 CMAT(L1,L2) = SRULE(4,A,X,C,B)
C ***
COMPUTE I33(BAR N) MATRIX ***
DO 100 I=1,N2
L1 = N4 + I
DO 100 J=1,N2
L2 = N5 + J
DO 95 K=1,NSEG2
T(1,K) = E13(I,K) * E13R(J,K)
T(2,K) = E43(I,K) * E43B(J,K) * DS(K)
95 T(3,K) = E53(I,K) * E53B(J,K)
100 CMAT(L1,L2) = SRULE(3,A,C,B)
C ***
COMPUTE I33(BAR N) MATRIX ***
DO 110 I=1,N2
L1 = N5 + I
DO 110 J=1,N2
L2 = N5 + J
DO 105 K=1,NSEG2
T(1,K) = E13B(I,K) * E13B(J,K)
T(2,K) = E23(I,K) * E23(J,K) * DR(K)
T(3,K) = E43B(I,K) * E43B(J,K) * DS(K)
105 T(4,K) = E53B(I,K) * E53B(J,K)
110 CMAT(L1,L2) = SRULE(4,A,X,C,B)
NE = 4 * (N+1)
DO 120 I=1,NE
DO 120 J=I,NE
120 CMAT(J,I) = CMAT(I,J)
K = 0.
DO 53 IN=1,4
L = IN
DO 53 J=1,3
K = K + 1
CMAT(K,13) = Q1(L,11)
CMAT(K,14) = (Q1(L,1) - Q2(L,11))

```

```
CMAT(13,K) = CMAT(K,13)
CMAT(14,K) = CMAT(K,14)
53 L4 = L4 + 4
      CMAT(13,13) = 0•
      CMAT(13,14) = 0•
      CMAT(14,14) = 0•
      CMAT(14,13) = 0•
      RRETURN
      END
```

```
CMAT1242
CMAT1243
CMAT1244
CMAT1245
CMAT1246
CMAT1247
CMAT1248
CMAT2180
CMAT2190
```

SIBFTC INPUT.

```
      SUBROUTINE INPUT
COMMON/DAT1/ RADIUS,HTCON,AL,YOUNG,HTCY,RHO,RHOPRO,POISON,ALPHA,D,INPUT005
      1 THICK
COMMON/DAT2/ NITER,N,NSEG
COMMON/BAR/ HBAR,BBAR,CAP,G11,E10, G12, RHOBAR
COMMON/DAT3/ ALPHAN(5),B(5),U(5)
      READ5,TOT RADIUS,HTCY,THICK,HTCON,AL,YOUNG,RHO,RHOPRO,ALPHA,D,
      1POISON
      10 FORMAT(6E12.8)
      READ(5,15) N,NSEG,NITER
      READ(5,10) (ALPHAN(I),I=1,5)
      READ(5,10) (B(I),I=1,5)
      READ(5,TOT TUTTY),T=1,5)
      15 FORMAT(6I12)
      HBAR = HTCY / RADIUS
      BBAR = HTCON / RADIUS
      CAP = (RHO * THICK * RADIUS**2) / D
      RHOBAR = RHOPRO / RHO
      WRITE(6,20)
      20 FORMAT(1H1,24X,6OHLONGITUDINAL OSCILLATION OF A LIQUID-FILLED CYLINDER)
      1 INDRICAL TANK/38X,33HW WITH AN INVERTED CONICAL BULKHEAD)
      WRITE(6,30) RADIUS,THICK,HTCY,HTCON,AL, ALPHA
      30 FORMAT(1H0,4UX,21HRADIUS OF CYLINDER = ,F7.27 38X,24HTHICKNESS OF
      1CYLINDER = ,F7.2/ 41X,21HHEIGHT OF CYLINDER = ,F7.2/1HU,43X,18HHEI INPUT100
      2GHT OF CONIC = ,F7.2/58X,4HL = ,F7.2754X, 8HATPHA = ,F7.2)
      WRITE(6,40) RHO,RHOPRO,POISON,YOUNG
      40 FORMAT(1H0,51X,10HDENSITY = ,1PE13.6/ 41X,21HPROPELLANT DENSITY =
      1,1PE13.6/ 46X,16HPOISONS RATIO = ,1PE13.6/ 45X,17HYOUNGS MODULUS = INPUT130
      2 ,1PE13.6)
      WRITE(6,50) N,NSEG
      50 FORMAT(1H0,53X,8HDEGREE = ,13/-45X,17HNO. OF SEGMENTS = ,13)
      RETURN
      END
```

111-

SID 67-212-2

```
      INPUT003
      INPUT004
      INPUT005
      INPUT010
      INPUT015
      INPUT017
      INPUT020
      INPUT025
      INPUT030
      INPUT035
      INPUT040
      INPUT042
      INPUT043
      INPUT044
      INPUT045
      INPUT050
      INPUT055
      INPUT060
      INPUT062
      INPUT070
      INPUT080
      INPUT090
      INPUT095
      INPUT097
      INPUT100
      INPUT110
      INPUT115
      INPUT120
      INPUT130
      INPUT140
      INPUT150
      INPUT160
      INPUT170
      INPUT180
```

```

$IBFTC TRANS SDD
C *THIS ROUTINE COMPUTES SIN, COSIN, SINH, COSH, AND BESSSEL FUNCTIONS*TRANS001
SURROUNTE TRANSF
REAL JIAN, JOAN, IIBN, JOANR
COMMON/BAR/ HBAR,BBAR,CAP,G11,E10,G12, RHOBAR
COMMON/DAT2/ NITER,N,NSEG
COMMON/DAT3/ ALPHAN(5),B(5),U(5)
COMMON/TRIG/ SINZ(5,11),CSINZ(5,11),SINHZ(5,11),CSINHZ(5,11),
1CSINR(5),CSINHR(5),SINHB(5)
COMMON/BESSEL/ JIAN(5), JOAN(5), IIBN(5),JOANR(5,11)
DEL = (HBAR - BBAR) / FLOAT(NSEG)
N2 = NSEG + 1
DEL2 = 1. / FLOAT(NSEG)
DELTAR = 0.
DELTAZ = BBAR
DO 20 I=1,N
DO 10 J=1,N2
X2 = ALPHAN(I) * DELTAZ
X3 = B(I) * (DELTAZ - BBAR)
X4 = ALPHAN(I) * DELTAR
SINZ(I,J) = SIN(X3)
CSINZ(I,J) = COS(X3)
SINHZ(I,J) = SINH(X2)
CSINHZ(I,J) = COSH(X2)
JOANR(I,J) = BESEL(X4,0,1)
DELTAZ = DELTAZ + DEL
10 DELTAR = DELTAR + DEL2
DELTAZ = BBAR
20 DELTAR = 0.
DO 30 I=1,N
X = ALPHAN(I) * BBAR
SINHR(I) = SINH(X)
CSINHB(I) = COSH(X)
CSINBN(I) = COS(B(I)*BBAR)
JIAN(I) = BESEL(ALPHAN(I),1,1)
JOAN(I) = BESEL(ALPHAN(I),0,1)
IIBN(I) = BESEL(B(I),0,2)
30 IIBN(I) = BESEL(B(I),1,2)
RETURN
END

```

SIRFTC F5ZN.

SDD

SUBROUTINE F5ZN

COMMON/BAR / HBAR, BBAR, CAP, G11, E10, G12, RHOBAR

COMMON/DAT2/ NITER, N, NSEG

COMMON/F56Z/ F5Z(6,11), F6Z(6,11)

DEL = 0.

DELZ = BBAR / FLOAT(NSEG)

N2 = N + T

NSEG2 = NSEG + 1

A6 = 1./6.

A8 = 35./8.

A9. = 21./8.

A16 = 63./16.

A17 = 35.716.

A18 = 35./6.

A31 = 106./30.

A24 = 11./24.

A25 = 245./724.

A30 = 181./30.

A56 = 55.756.

DO 30 J=1,NSEG2

A = BBAR - DEL

IF (J.EQ.NSEG2) A = 0.

U = A / SQRT(1. + A**2)

EP = SQRT(1. - U**2)

UE = U / EP

UE3 = U**3 / EP**3

UE5 = U / EP**5

E1 = 1. / EP**2

E2 = 1. / EP**4

F5Z(1,J) = -UE

F5Z(2,J) = -UE

F5Z(3,J) = .5 * UE - UE3/3.

F5Z(4,J) = 1.25 * E1 - .25 * E2

F5Z(5,J) = 4. * UE - A24 * UE3 - A8 * UE5 + A25 * UE5 * U**2 - A30F5ZNU150

1 * UE3 * E1 * U**2

F5Z(6,J) = -A16 * E1 - A9 * E2 - A6 * E1 * E2

30 DEL = DEL + DELZ

ZBAR = 0.

DO 40 J=1,NSEG2

BZ = BBAR - ZBAR

IF (J.EQ.NSEG2) BZ = 0.

BZ2 = BZ**2

40 ZBAR = ZBAR + BZ2

F5ZN00000

F5ZN00005

F5ZN0010

F5ZN0015

F5ZN0020

F5ZN0025

F5ZN0030

F5ZN0035

F5ZN0040

F5ZN0045

F5ZN0050

F5ZN0055

F5ZN0060

F5ZN0062

F5ZN0063

F5ZN0064

F5ZN0065

F5ZN0070

F5ZN0075

F5ZN0076

F5ZN0085

F5ZN0090

F5ZN0092

F5ZN0095

F5ZN100

F5ZN0105

F5ZN0110

F5ZN0115

F5ZN0120

F5ZN0125

F5ZN0130

F5ZN0135

F5ZN0140

F5ZN0145

F5ZN0150

F5ZN0155

F5ZN0160

F5ZN0170

F5ZN0180

F5ZN0190

F5ZN0200

F5ZN0205

F5ZN0210

```

BZ3 = BZ2 * BZ          F5ZN0220
BZ5 = BZ3 * BZ2         F5ZN0230
B1Z = 1. + BZ2          F5ZN0240
B2Z = B1Z**2            F5ZN0250
B4Z = B2Z**2            F5ZN0260
F6Z(1,J) = -BBAR * BZ + .5 * B1Z   F5ZN0270
F6Z(2,J) = -BBAR * .5 * R2Z + BZ3/3.  F5ZN0280
F6Z(3,J) = BBAR * (.5 * BZ - BZ3/3.) - .75 * BTZ + .25 * BZZ  F5ZN0290
F6Z(4,J) = BBAR * (1.25 * B2Z - .25 * B4Z) - 2.5 * BZ + BZ3/3. + F5ZN0300
12.5 * BZ * B2Z - A18 * BZ3 * B1Z + A31 * BZ5    F5ZN0310
F6Z(5,J) = BBAR * (4.* BZ - A24 * BZ3 - A8 * BZ * B4Z + A25 * RZ3 * F5ZN0320
1 R2Z - A30 * BZ5) + A18 * B1Z - 1.25 * B2Z + A6 * B1Z * R2Z  F5ZN0330
F6Z(6,J) = BBAR * (-A16 * B2Z - A9 * B4Z - A6 * B2Z * B4Z) - 12.15 F5ZN0340
1* BZ + 7.875 * RZ * R1Z + 1.45 * BZ3 + 2.15 * BZ * BZZ - 18.375 * F5ZN0350
2RZ3 * B2Z + 11.3 * BZ5 - A56 * BZ2               F5ZN0360
40 ZRAR = ZBAR + DELZ
      RETURN
      END

```

\$IBFTC CONST.

```
SUBROUTINE CONST
COMMON/BAR/ HBAR, BBAR,CAP, G11, E10,G12,RHOBAR
COMMON/DAT?/ NITER,N, NSEG
COMMON/DEL/DEL1,DEL2,DEL3,DEL4
COMMON/DELT/A/ DZBH(11),DSL(11), DR(11), DZB(11),DS(11)
COMMON/LIMIT/ N2,NSEG2
COMMON/DAT1/RADIUS,HTCON,AL,YOUNG,HTCY,RHO, RHO PRO, POTSON,ALPHA, CONSTO27
      D, THICK
      N2 = N+1
      NSFG? = NSEG+1
      DFL1 = (HRAR-BBAR) / FLOAT(NSEG)
      DFL2 = BBAR / FLOAT(NSEG)
      BABL = SQRT(1.+BBAR**2)
      DFL3 = BABL / FLOAT(NSEG)
      DFL4 = 1. / FLOAT(NSEG)
      ZBAR = 0.
      DZBH(1) = BBAR
      DZS(1) = SQRT(1.+RRAR**2)
      DSL(1) = 0.
      DR(1) = 0.
      DB = 0.
      DO 10 I=2,NSEG2
      DZBH(I) = DZBH(I-1) + DEL1
      ZBAR = ZBAR + DEL2
      DZS(I) = SQRT(1.+(BBAR-ZBAR)**2)
      DSL(I) = DSL(I-1) + DEL3
      DR(I) = DR(I-1) + DEL4
10   DZB(I) = DZB(I-1) + DEL2
      A = SIND(ALPHA)
      DO 20 I=1,NSEG2
20   DS(I) = DSL(I) * A
      RETURN
      FND
```

\$*

*IBFTC F1ZN• SDD

SUBROUTINE F1ZN

C

COMMON/RAR/ HBAR,BBAR,CAP,G11,E10,G12,RHOBAR

COMMON/DAT2/ NITER,N,NSEG

COMMON/DAT3/ ALPHAN(5), B(5), U(5)

COMMON/F1/F1Z(24,11)

N2 = N + 1

N3 = NSEG + 1

DFLTA = (HRAR - RBAR) / FLOAT(NSEG)

DFL = BRAR

K2 = 1

DO 10 K=1,N3

F1Z(1,K) = 1. / G11**2

F1Z(2,K) = 0.

F1Z(3,K) = -E10 / G11

10 F1Z(4,K) = 0.

DO 40 I=1,N

K2 = K2 + 4

X = BESEL(B(I),0,2) / (B(I) * COS(B(I)*BEART))

DO 40 K = 1,N3

A = R(I) * DEL - BRAR

F1Z(K2,K) = 0.

F1Z(K2+1,K) = X * SIN(A)

F1Z(K2+2,K) = 0.

F1Z(K2+3,K) = X * COS(A)

40 DEL = DEL + DELTA

RETURN

END

\$*

F1ZN0000

F1ZN0002

F1ZN0005

F1ZN0010

F1ZN0015

F1ZN0017

F1ZN0020

F1ZN0025

F1ZN0030

F1ZN0035

F1ZN0036

F1ZN0038

F1ZN0050

F1ZN0055

F1ZN0060

F1ZN0070

F1ZN0080

F1ZN0090

F1ZN0100

F1ZN0105

F1ZN0110

F1ZN0115

F1ZN0120

F1ZN0125

F1ZN0130

F1ZN0140

F1ZN0150

F1ZN0160

F1ZN0170

FH1N0999

SIBFTC FH1N.

SDD

SUBROUTINE FH1N(OMEGA)
COMMON/BAR/HBAR,BBAR,CAP,G11,E10,G12, RHOBAR
COMMON/DAT2/ NITER,N,NSEG
COMMON/H1/ FH1(8,6)

COMMON/DAT1/RADIUS,HTCON,AL,YOUNG, HTCY,RHO,RHOPRO,POTSON,ALPHA,D,FH1N0025
1 THICK

```

COMMON/DAT3/ALPHAN(5),B(5),U(5)
X = G11 * HBAR
SIGN = SIN(X)
COSIGN = COS(X)
WK1 = 1. - CAP * OMEGA**2
WK2 = 1. - POISON**2 / WK1
WK3 = 1. - POISON / WK1
GG = G12 * G11
X2 = G11 * BBAR
SIGN2 = SIN(X2)
COSN2 = COS(X2)
FH1(1,1) = -WK2 * GG*(SIGN / G11)**2 + COSIGN * E10 / G11
FH1(2,1) = 0.
FH1(3,1) = -GG * (SIGN2 / G11)**2 + COSN2 * E10 / G11
FH1(4,1) = 0.
Y2 = HBAR - BBAR
DO 40 I=1,N
XA = ALPHAN(I) * HBAR
XX = EXP(XA)
HYSIN = (XX - 1. / XX) / 2.
HYCOS = (XX + 1. / XX) / 2.
BES = (BESEL(R(I),0,2)) / (B(I) * COS(B(I)*BBAR))
FH1(1,I+1) = (WK2 * G12 - POISON / WK1) * (HYSIN + HYCOS) * BESELFH1N0210
1(ALPHAN(I),0,1)
FH1(2,I+1) = -WK2 * GG * (SIGN * BES * TSINT(B(I)*Y2) - COSTN * BESFH1N0230
1 * COS(R(I)*Y2) + (WK2 * G12 - POISON / WK1) * COS(B(I)*Y2) / COS(FH1N0235
2B(I)*BBAR) * BESEL(B(I),0,2)
FH1(3,I+1) = G12 *(SINH(ALPHAN(I)*BBAR) +COSH(ALPHAN(I)*BBAR)) *
1BESEL(ALPHAN(I),0,1)
40 FH1(4,I+1) = GG * COSH2 * BES * COS(B(I)) + G12 * BESEL(B(I),0,2)FH1N0270
1 / COS(B(I)* BBAR)
U1 = BBAR / SQRT(1. + BBAR**2)
EP = SQRT(1. - U1**2)
UE = U1 / EP
UE3 = UE**3
FH1(5,1) = -UE

```

FH1N0000	FH1N0005
COMMON/BAR/HBAR,BBAR,CAP,G11,E10,G12, RHOBAR	FH1N0010
COMMON/DAT2/ NITER,N,NSEG	FH1N0015
COMMON/H1/ FH1(8,6)	FH1N0020
COMMON/DAT1/RADIUS,HTCON,AL,YOUNG, HTCY,RHO,RHOPRO,POTSON,ALPHA,D,FH1N0025	FH1N0030
1 THICK	FH1N0032
COMMON/DAT3/ALPHAN(5),B(5),U(5)	FH1N0035
X = G11 * HBAR	FH1N0040
SIGN = SIN(X)	FH1N0045
COSIGN = COS(X)	FH1N0050
WK1 = 1. - CAP * OMEGA**2	FH1N0055
WK2 = 1. - POISON**2 / WK1	FH1N0060
WK3 = 1. - POISON / WK1	FH1N0065
GG = G12 * G11	FH1N0070
X2 = G11 * BBAR	FH1N0080
SIGN2 = SIN(X2)	FH1N0090
COSN2 = COS(X2)	FH1N0100
FH1(1,1) = -WK2 * GG*(SIGN / G11)**2 + COSIGN * E10 / G11	FH1N0110
FH1(2,1) = 0.	FH1N0120
FH1(3,1) = -GG * (SIGN2 / G11)**2 + COSN2 * E10 / G11	FH1N0130
FH1(4,1) = 0.	FH1N0140
Y2 = HBAR - BBAR	FH1N0150
DO 40 I=1,N	FH1N0160
XA = ALPHAN(I) * HBAR	FH1N0170
XX = EXP(XA)	FH1N0180
HYSIN = (XX - 1. / XX) / 2.	FH1N0190
HYCOS = (XX + 1. / XX) / 2.	FH1N0210
BES = (BESEL(R(I),0,2)) / (B(I) * COS(B(I)*BBAR))	FH1N0220
FH1(1,I+1) = (WK2 * G12 - POISON / WK1) * (HYSIN + HYCOS) * BESELFH1N0230	FH1N0235
1(ALPHAN(I),0,1)	FH1N0240
FH1(2,I+1) = -WK2 * GG * (SIGN * BES * TSINT(B(I)*Y2) - COSTN * BESFH1N0230	FH1N0250
1 * COS(R(I)*Y2) + (WK2 * G12 - POISON / WK1) * COS(B(I)*Y2) / COS(FH1N0235	FH1N0260
2B(I)*BBAR) * BESEL(B(I),0,2)	FH1N0270
FH1(3,I+1) = G12 *(SINH(ALPHAN(I)*BBAR) +COSH(ALPHAN(I)*BBAR)) *	FH1N0280
1BESEL(ALPHAN(I),0,1)	FH1N0290
40 FH1(4,I+1) = GG * COSH2 * BES * COS(B(I)) + G12 * BESEL(B(I),0,2)FH1N0270	FH1N0300
1 / COS(B(I)* BBAR)	FH1N0310
U1 = BBAR / SQRT(1. + BBAR**2)	FH1N0320
EP = SQRT(1. - U1**2)	FH1N0330
UE = U1 / EP	
UE3 = UE**3	
FH1(5,1) = -UE	

```

FH1(6,1) = -•5 / EP**2
FH1(5,2) = -•5 * FH1(5,1) - UE3 / 3.
FH1(6,2) = 1•25 / EP**2 - •25 / EP**4
FH1(5,3) = 4• * UF - 11• * UE3 / 24. - 4•375 * UE7 EP**4 + 24• * FH1N0370
1 UE3 / (24. * EP**) - 181• * UE**5 / 30.
FH1(6,2) = -62. / (16.* FP**2) - 21. / (8. * EP**4) - 1. / 7(6. * EPFH1N0390
1**6)                                             FH1N0400
FH1(7,1) = -GG * (-G11/2. * COSN2) + G12
FH1(8,1) = -GG * (-•5 * SIGN2 - G11 * (-BBAR/2.))
FH1(7,2) = -GG * G11 * 0.5 * COSN2 - •5 * 612
FH1(8,2) = -GG * (SIGN2 - G11 * BBAR)
FH1(7,3) = -GG * (-53./48. * G11 * COSN2) + •375 * G12
FH1(8,3) = -GG * (-323./48. * SIGN2 + 323./48. * BBAR * G11)
RFTURN
FND

```

\$IBFTC POLY.

```
C *** ROUTINE COMPUTES THE FIRST SIX LEGENDRE POLYNOMIAL ***
C
C SUBROUTINE POLY
COMMON/BAR/ BBAR, CAP, G11,E10, G12, RHOBAR
COMMON/DAT2/ NITER,N,NSEG
COMMON/DATT/RADIUS,HTCON, AL, YOUNG, HTCY, RHO, RHOPRO, POTSON, POLY0000
      1 ALPHA, D, THICK
      COMMON/PLDR/ PO(6), PA(6), PZ(6,11)
      U = COSD(ALPHA)
      U2 = U**2
      U3 = U2 * U
      U4 = U2**2
      U5 = U2 * U3
      PA(1) = 1.
      PA(2) = U
      PA(3) = 0.5 * (3.*U2 - 1.)
      PA(4) = 0.5 * (5.*U3 - 3.*U)
      PA(5) = 0.125 * (35.*U4 - 30.*U2 + 3.)
      PA(6) = 0.125 * (63.*U5 - 70.*U3 + 15.*U)
      PO(1) = 1.
      PO(2) = 0.
      PO(3) = -5
      PO(4) = 0.
      PO(5) = .375
      PO(6) = 0.
      DEL = 0.
      DZ = BBAR / FLOAT(NSEG)
      NSEG2 = NSEG+1
      DO 10 J=1,NSEG2
      B1 = BBAR - DEL
      IF(J.EQ.NSEG2) B1 = 0.
      U = B1 / SQRT(1. + B1**2)
      U2 = U**2
      U3 = U**3
      PZ(1,J) = 1.
      PZ(2,J) = U
      PZ(3,J) = 0.5 * (3.*U2 - 1.)
      PZ(4,J) = 0.5 * (5.*U3 - 3.*U)
      PZ(5,J) = 0.125 * (35.*U2**2 - 30.*U2 + 3.)
      PZ(6,J) = 0.125 * (63.*U3*U2 - 70.*U3 + 15.*U)
      10 DEL = DEL + DZ
      POLY0005
      POLY0010
      POLY0015
      POLY0020
      POLY0025
      POLY0030
      POLY0035
      POLY0040
      POLY0045
      POLY0050
      POLY0055
      POLY0060
      POLY0065
      POLY0070
      POLY0075
      POLY0080
      POLY0085
      POLY0090
      POLY0095
      POLY0100
      POLY0105
      POLY0110
      POLY0115
      POLY0120
      POLY0125
      POLY0130
      POLY0135
      POLY0140
      POLY0145
      POLY0150
      POLY0155
      POLY0160
      POLY0162
      POLY0165
      POLY0170
      POLY0175
      POLY0180
      POLY0185
      POLY0190
      POLY0195
      POLY0200
      POLY0210
      POLY0215
```

RETURN
END

POLY 0225

-120-

SID 67-212-2

SIRFTC GAMA.

C***** COMPUTE GAMMA FUNCTION *****

SUBROUTINE GAMAN

COMMON/GAMA/ GAMA1(9,11),GAMA2(12,11),GAMA3(9,11)

COMMON/BAR/ HBAR, BBAR, CAP, G11,E10, G12, RHOBAR

COMMON/DAT1/ RADIUS,HTCON,AL,YOUNG,HTCY,RHO,RHOPRO,POTSON,ALPHA,D,GAMMA018

GAMMA019

1 THICK
COMMON/TRIG/ SINZ(5,11),CSINZ(5,11),CSINHZ(5,11),

1CSINBN(5),CSINHB(5),SINHB(5)

COMMON/BFSSEL/ JIAN(5),JOAN(5),IIBN(5),JOANR(5,11)

REAL JIAN, JUAN, IIBN,JOANR,BBAR

COMMON/DAT2/ NITER,N,NSEG

COMMON/DAT3/ ALPHAN(5),B(5),U(5)

LBAR = SQRT(1. + BBAR**2)

DEL1 = BBAR

N2 = N + 1

DEL2 = 0.

NSEG2 = NSEG + 1

DZ = (HBAR - BBAR) / FLOAT(NSEG)

DR = T. / FLOAT(NSEG)

K = 4

DO 20 I=1,N

A = B(I) * IIBN(I) / CSINBN(I)

A2 = ALPHAN(I) * (CSINHB(I) + SINHB(I))

DO 10 J=1,NSEG2

GAMA1(K,J) = 0.

GAMA1(K+1,J) = A * CSINZ(I,J)

10 GAMA1(K+2,J) = A2 * JOANR(I,J)

20 K = K + 3

DO 25 J=1,NSEG2

GAMA1(1,J) = 0.

GAMA1(2,J) = 0.

25 GAMA1(3,J) = 1.

C *** COMPUTE LEGENDRE POLYNOMIAL ***

DIMENSION P(6,11)

DEL = BBAR / FLOAT(NSEG)

DZ = 0.

DO 30 J=1,NSEG2

B1= BBAR - DZ

U1= B1/ SQR(1. + B1**2)

P(1,J) = 1.

P(2,J) = U1

U2 = U1**2

GAMMA000

GAMMA001

GAMMA002

GAMMA010

GAMMA015

GAMMA018

GAMMA019

GAMMA020

GAMMA025

GAMMA030

GAMMA035

GAMMA040

GAMMA050

GAMMA055

GAMMA060

GAMMA065

GAMMA070

GAMMA075

GAMMA080

GAMMA090

GAMMA100

GAMMA110

GAMMA120

GAMMA122

GAMMA130

GAMMA140

GAMMA150

GAMMA160

GAMMA170

GAMMA172

GAMMA174

GAMMA176

GAMMA178

GAMMA179

GAMMA180

GAMMA190

GAMMA200

GAMMA210

GAMMA220

GAMMA230

GAMMA250

GAMMA260

GAMMA270

SID 67-212-2

```

U3 = U2 * U1
P(3,J) = .5 * (3. * U2 - 1.)
P(4,J) = .5 * (5. * U3 - 3. * U1)
P(5,J) = .125 * (35. * U2**2 - 30. * U2 + 3.)
P(6,J) = .125 * (63. * U1**5 - 70. * U3 + 15. * U1)
30 D2 = DZ + DEL
D2 = 0.
DO 40 J=1,NSEG2
B1 = BBAR - DZ
GAMA3(1,J) = SQRT(1. + B1**2)
GAMA3(2,J) = B1 / GAMA3(1,J)
GAMA3(3,J) = 1. / SQRT(1. - GAMA3(2,J) * GAMA3(2,J))
40 D2 = DZ + DEL
K = 1
DO 55 J=1,NSEG2
GAMA2(1,J) = 0.
GAMA2(2,J) = P(2,J)
GAMA2(3,J) = 0.
55 GAMA2(4,J) = -GAMA3(3,J) * (P(1,J) - GAMA3(2,J) * P(2,J))
DO 70 I=1,N
NCAP = 2 * I
NRAR = 2 * I + 1
A = FLOAT(NCAP)
AX = FLOAT(NBAR)
DO 65 J=1,NSEG2
GAMA2(K,J) = A * GAMA3(1,J)**(NCAP-1) * P(NCAP+1,J)
GAMA2(K+1,J) = AX * GAMA3(1,J)**NCAP * P(NBAR+1,J)
GAMA2(K+2,J) = -GAMA3(3,J) * A * GAMA3(1,J)**(NCAP-1) * (P(NCAP,J)
1 - GAMA2(2,J) * P(NCAP+1,J))
65 GAMA2(K+3,J) = -GAMA3(3,J) * AX * GAMA3(1,J)**NCAP * (P(NBAR,J) -
1GAMA3(2,J) * P(NBAR+1,J))
70 K = K + 4
U1 = COSD(ALPHA)
P(1,1) = 1.
P(1,2) = U1
U2 = U1**2
P(1,3) = .5 * (3. * U2 - 1.)
U3 = U2 * U1
P(1,4) = .5 * (5. * U3 - 3. * U1)
P(1,5) = .125 * (35. * U2**2 - 30. * U2 + 3.)
P(1,6) = .125 * (63. * U3 * U2 - 70. * U3 + 15. * U1)
P(2,1) = 1.

```

```

P(2,2) = 0.
P(2,3) = -•5
P(2,4) = 0.
P(2,5) = •375
P(2,6) = 0.
DFL = 0.

A =-1. / SQRT(1. - U1**2)
DELL = LBAR / FLOAT(NSFG)
DEL2 = 0.
DFLR = 1. / FLOAT(NSEG)
DO 85 J=1,NSEG2
P(3,J) = DEL
P(4,J) = DEL2
DEL = DFL + DELL
85 DEL2 = DEL2 + DFLR
DO 90 J=1,NSEG2
GAMA3(1,J) = 0.
GAMA3(2,J) = A * (P(1,1) - U1 * P(1,2))
90 GAMA3(3,J) = -P(2,1)
GAMA3(2,1) = 0.
GAMA3(3,1) = 0.

K = 4
DO 100 I=1,N
NCAP = 2 * I
NBAR = NCAP + 1
NR = NBAR + 1
NC = NCAP - 1
DO 95 J=1,NSEG2
GAMA3(K,J) = A * FLOAT(NCAP) * P(3,J)**NC * (P(1,NCAP) - U1 * P(1,NCAP) - 1)
1NBAR)
GAMA3(K+1,J) = A * FLOAT(NBAR) * P(3,J)**NCAP * (P(1,NBAR) - U1 * GAMMA984
1P(1,NB))
95 GAMA3(K+2,J) = - FLOAT(NBAR) * P(4,J)**NCAP * (P(2,NBAR)) GAMMA990
100 K = K + 3
RETURN GAMMA995
END GAMMA999

```

```

$1RAFTC BETAN.
C*** SUBROUTINE TO CALCULATE BETAN(N) ***
      SUBROUTINE BETAN(OMEG,ERR)
COMMON/BAR/ HBAR, BRAR, CAP, G11, E10, G12, RHOBAR
COMMON/DATA/ ALPHAN(5),B(5),U(5)
COMMON/RET/ OMEG2,GA, BH
COMMON/DAT1/ RADIUS,HTCON, AL,YOUNG, HTCY,RHO, RHOPO, POISON,
1ALPHA,D, THICK
EXTERNAL BETA
LOGICAL FRR
ERR = .FALSE.
WRITE(6,5)
5 FORMAT(1H1)
GA = 386.089 / RADIUS
BH = HBAR - BRAR
TOL = .001
REL = 1.
NITER = 80
OMEGG2 = OMEG**2
K = 1
DFLRN = .5
RN = 1.58 / BH
11 BETN = GA * BN * TAN(RN * BH) + OMEG2
IF(BETN.LT.0.) GO TO 12
B(K) = BN
BN = BN + (3.14159 / BH)
IF(K.EQ.4) GO TO 50
K = K + 1
GO TO 11
12 CALL ROOT(NTRY,BN,DELBN,RESID,BETA,TOL,REL,NITER)
GO TO (15,15,30,30,35), NTRY
15 B(K) = BN
IF(K.EQ.4) GO TO 50
K = K + 1
BN = BN + (3.14159 / BH)
GO TO 11
20 WRITE(6,32) OMEG
22 FORMAT(1HO,10X,29H*** NO SOLUTIONS FOR OMEGA = ,1PE13.6)
ERR = .TRUE.
GO TO 50
35 WRITE(6,37) OMEG
37 FORMAT(1HO,10X,70H*** EXCEEDED MAXIMUM NO. OF ITERATIONS WITHOUT ABETAN330
1 SOLUTION FOR OMEGA = ,1PE13.6)

```

```
ERR = .TRUE.  
50 WRITE(6,55) OMEGA,T,R(I),I=1,K  
55 FORMAT(1H-,1UX,8HOMEGA = ,1PE13.6, 10H RAD./SEC./ (5X,2HB(,11,4H)BETAN370  
1 = ,1PE13.6)  
RETURN  
END
```

```
* I AFTC RFTA.  
FUNCTION RFTA(RN)  
COMMON/RFT/ OMEG2, GA, RH  
X = RN * RH  
RFTA = GA * BN * TAN(X) + OMEG2  
RETURN  
END
```

```
BETA0000  
BFTA0002  
BETA0004  
BETA0006  
BETA0008  
BETA0010  
BETA0012
```

```

$IBFTC SRULF.
      FUNCTION SRULE(N,A,B,C,D)
C *** INTEGRATION BY SIMPSON'S RULE ***
C   N = NO. OF INTEGRALS
COMMON/TERM/ T(4,11),ISEG,NSEG2,NN
SUM4 = 0.
SUM1 = T(1,1) + T(1,NSEG2) + 4. * T(1,ISEG)
SUM2 = T(2,1) + T(2,NSEG2) + 4. * T(2,ISEG)
SUM3 = T(3,1) + T(3,NSEG2) + 4. * T(3,ISEG)
IF (N.FQ.3) GO TO 10
SUM4 = T(4,1) + T(4,NSEG2) + 4. * T(4,ISEG)
DO. 5 I=2,NN,2
5 SUM4 = SUM4 + 4. * T(4,I) + 2. * T(4,I+1)
10 DO 20 T=2,NN,2
      SUM1 = SUM1 + 4. * T(1,I) + 2. * T(1,I+1)
      SUM2 = SUM2 + 4. * T(2,I) + 2. * T(2,I+1)
20 SUM3 = SUM3 + 4. * T(3,I) + 2. * T(3,I+1)
SRULE = A*SUM1/3. + B*SUM2/3. + C*SUM3/3. + D*SUM4/73.
RETURN
END

```

\$*